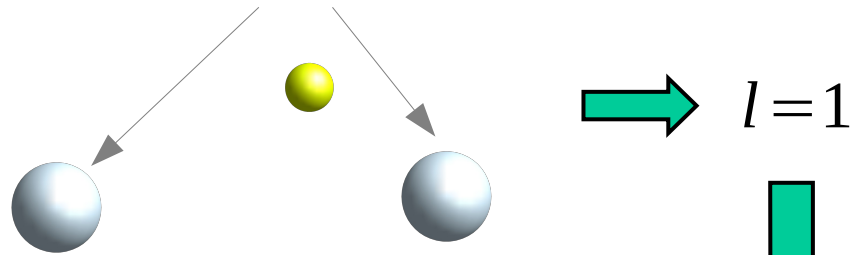


# Non-Efimovian (universal) clusters & 2D quantum anomaly

Dmitry Petrov

Laboratoire Physique Théorique et Modèles Statistiques (Orsay)

# Identical fermions



$$R \ll a \quad \longrightarrow \quad \tilde{U}_{\text{eff}}(R) \approx \frac{\hbar^2}{MR^2} \underbrace{\left( 2 - 0.16 \frac{M}{m} \right)}_{\beta}$$

$$\beta < -1/4, \quad M/m > 13.6$$



$$\chi(R) \propto \sqrt{R} \sin(\sqrt{-1/4 - \beta} \log R/r_3)$$



“Fall of a particle to the center in  $R^{-2}$  potential”. Infinite number of zeros of the wavefunction. Infinite number of trimer states. **Efimov effect**

Yb-Li, Dy-Li, Er-Li, ...?

$$\beta > -1/4, \quad M/m < 13.6$$

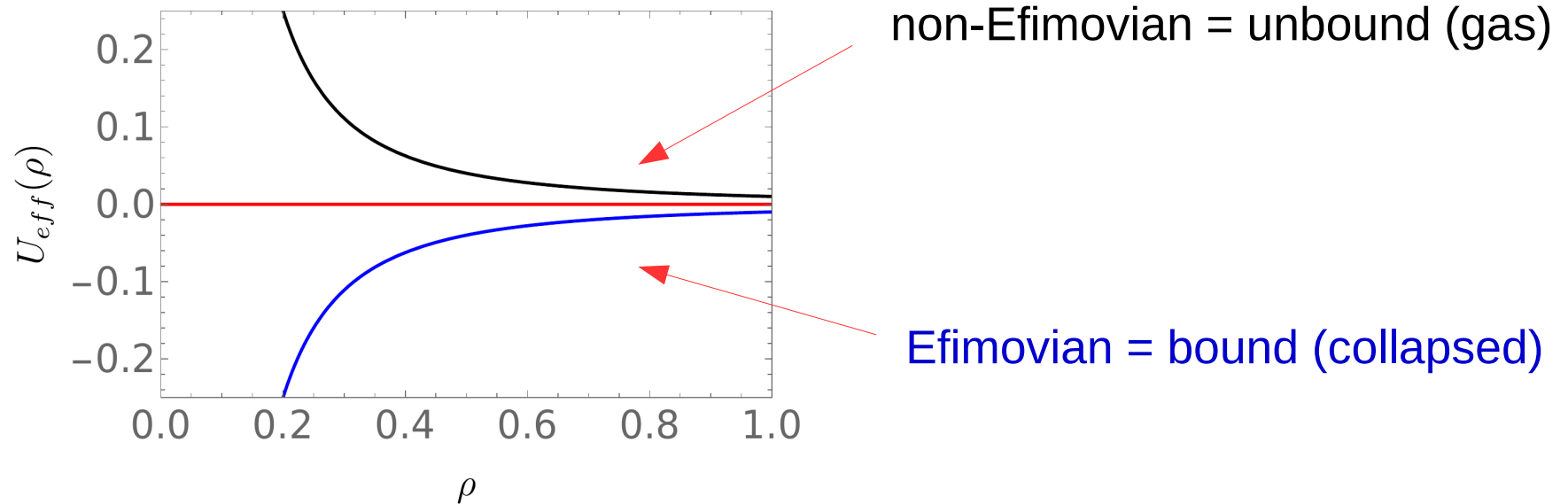


$$\chi(R) \propto R^{1/2 + \sqrt{\beta + 1/4}}$$



“Universal” regime in the sense that one needs no three-body parameter. **Fermi statistics wins over the induced attraction**

Homonuclear mixtures, K-Li, K-Dy, Li-Cr, K-Yb, K-Er, ...?



$$U_{eff}(\rho) \propto \frac{1}{\rho^2}$$

Valid only in the scale-invariant region  $\rho \ll |a|$

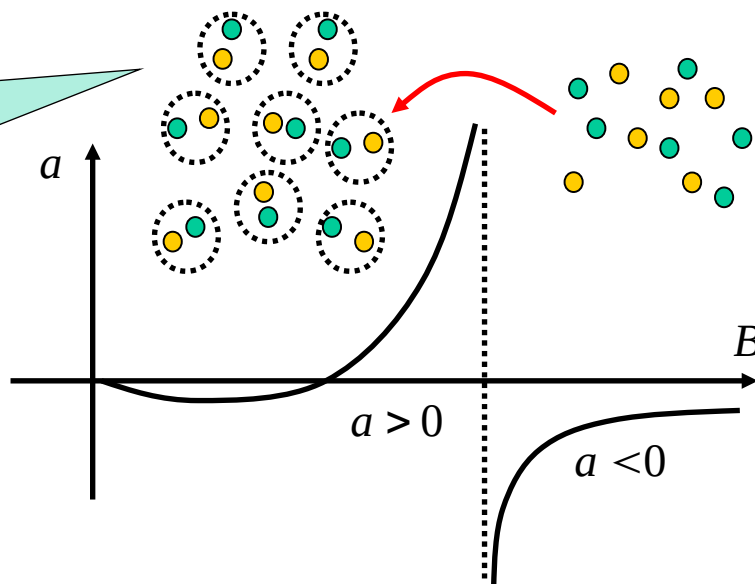
Consequence of the fact that for  $a = \infty$  the two-body interaction is scaling invariant

$$\frac{(r\psi(r))'}{r\psi(r)} = -\frac{1}{a} = 0 \quad \longleftrightarrow \quad \psi(r) \underset{r \rightarrow 0}{\rightarrow} C \left( \frac{1}{r} + o(r^0) \right)$$

Finite  $a$   $\rightarrow$  {  
 Efimovian  $\rightarrow$  loss resonances (« loss features »)  
 Non-Efimovian  $\rightarrow$  universality of few-body observables and universal bound states... and long lifetime

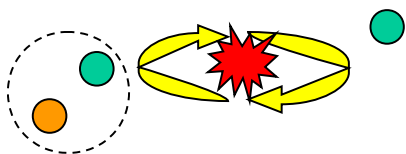
# 2003: BCS-BEC crossover, molecules

Bose gas of dimers  
("BEC side" of the resonance)



Two-component Fermi gas  
("BCS side" of the resonance)

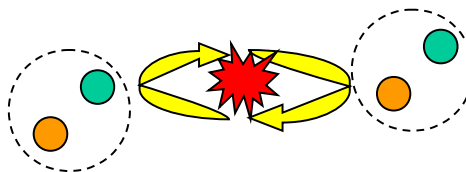
Few-body problem in this case is **non-efimovian**



$$a_{ad} = 1.2 a$$

Skorniakov,

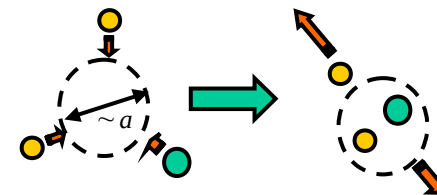
Ter-Martirosian (1957)



$$a_{dd} = 0.6 a$$

DSP, Salomon,

Shlyapnikov (2003)

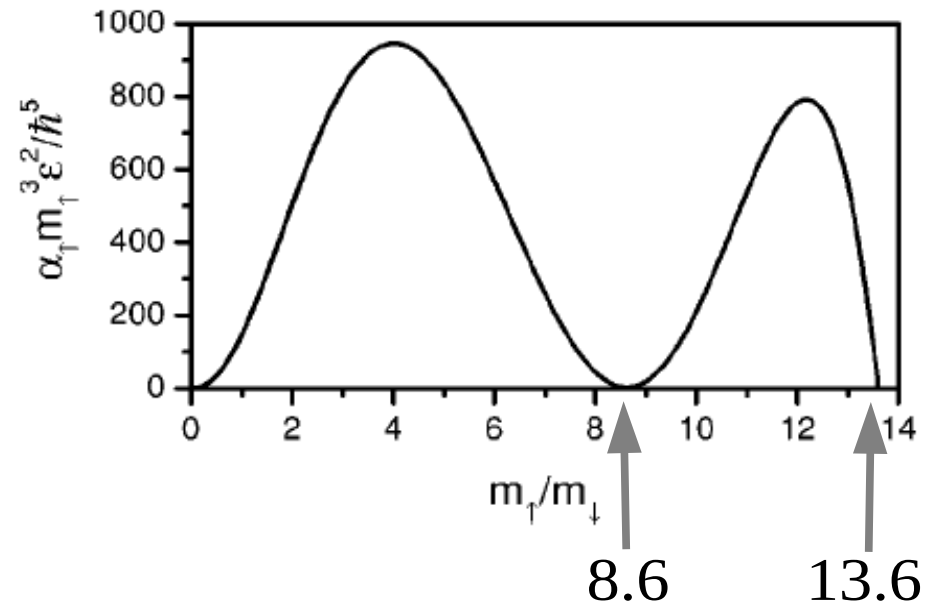
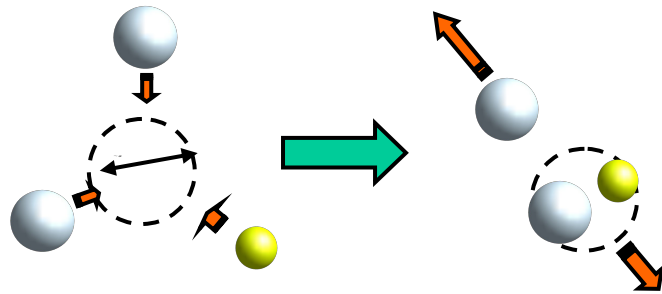


$$\alpha_{rec} = 148 \frac{\hbar a^4}{m} \cdot \frac{\bar{\epsilon}}{\epsilon_0}$$

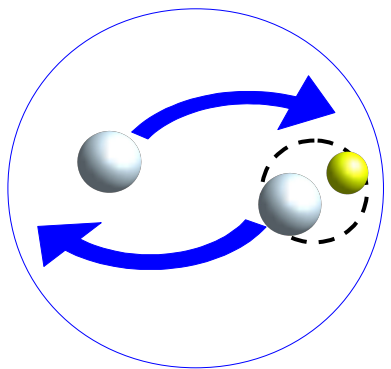
DSP (2003)

# Heavy-heavy-light problem, magic mass ratios

3-body recombination to a weakly bound level DSP'03

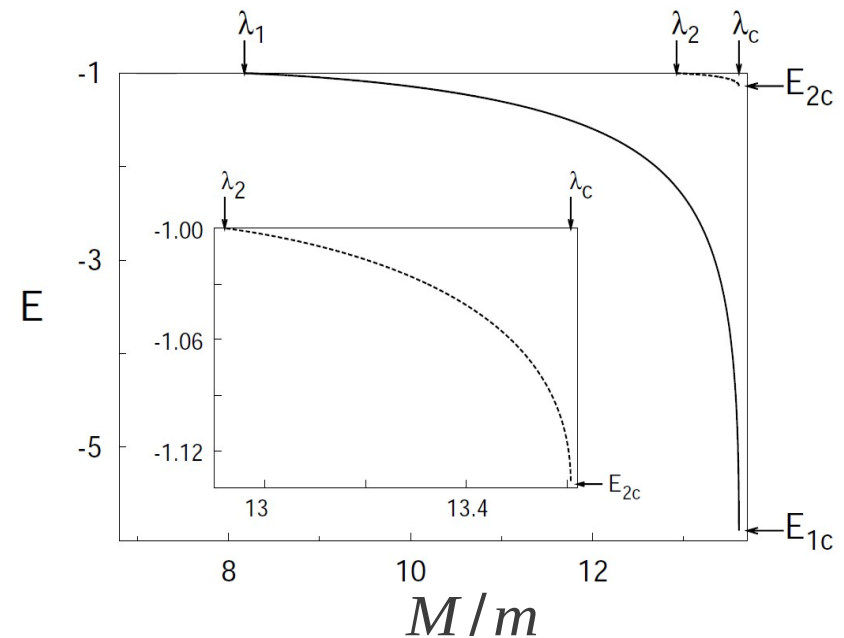


Emergence of a non-Efimovian trimer state for  $M/m > 8.2$  Kartavtsev&Malykh'06

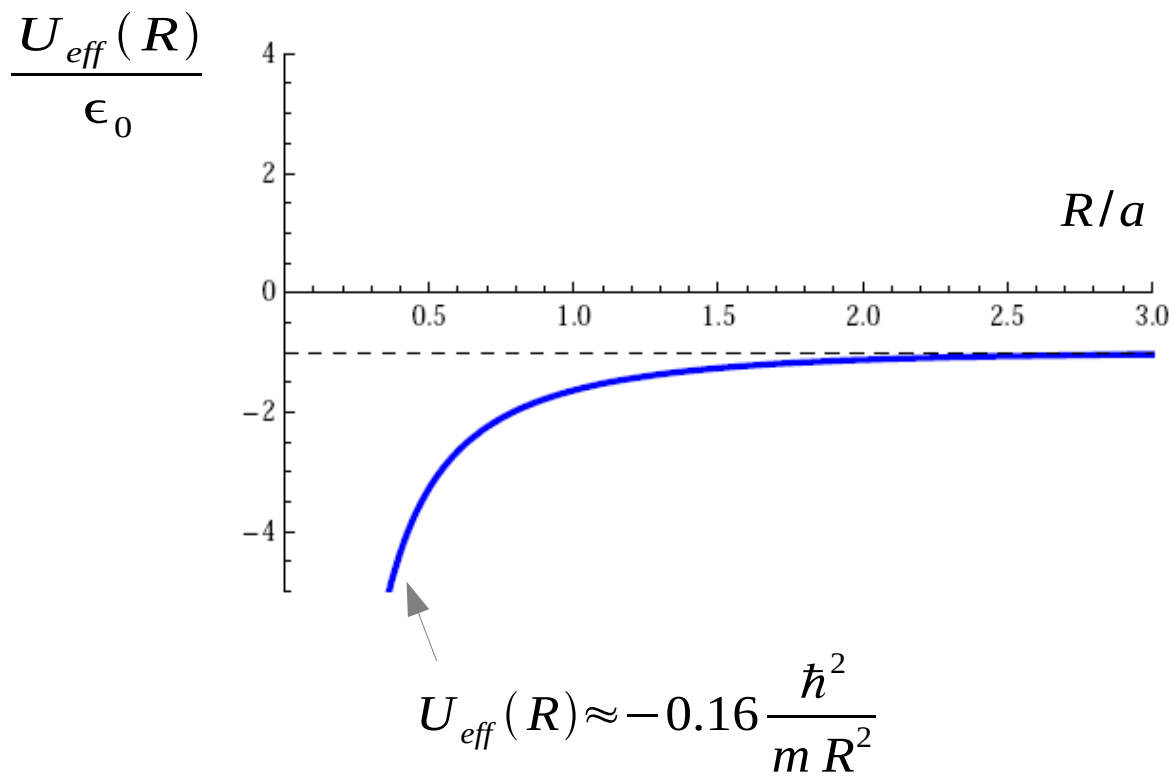
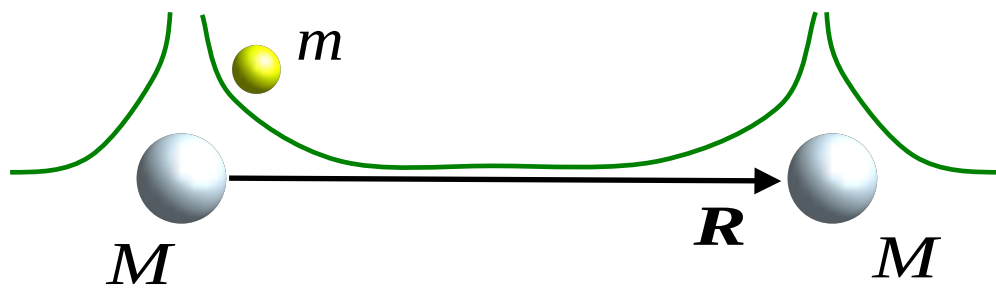


$M/m < 8.2$   $p$ -wave atom-dimer scattering resonance

$M/m > 8.2$  trimer state with  $l=1$



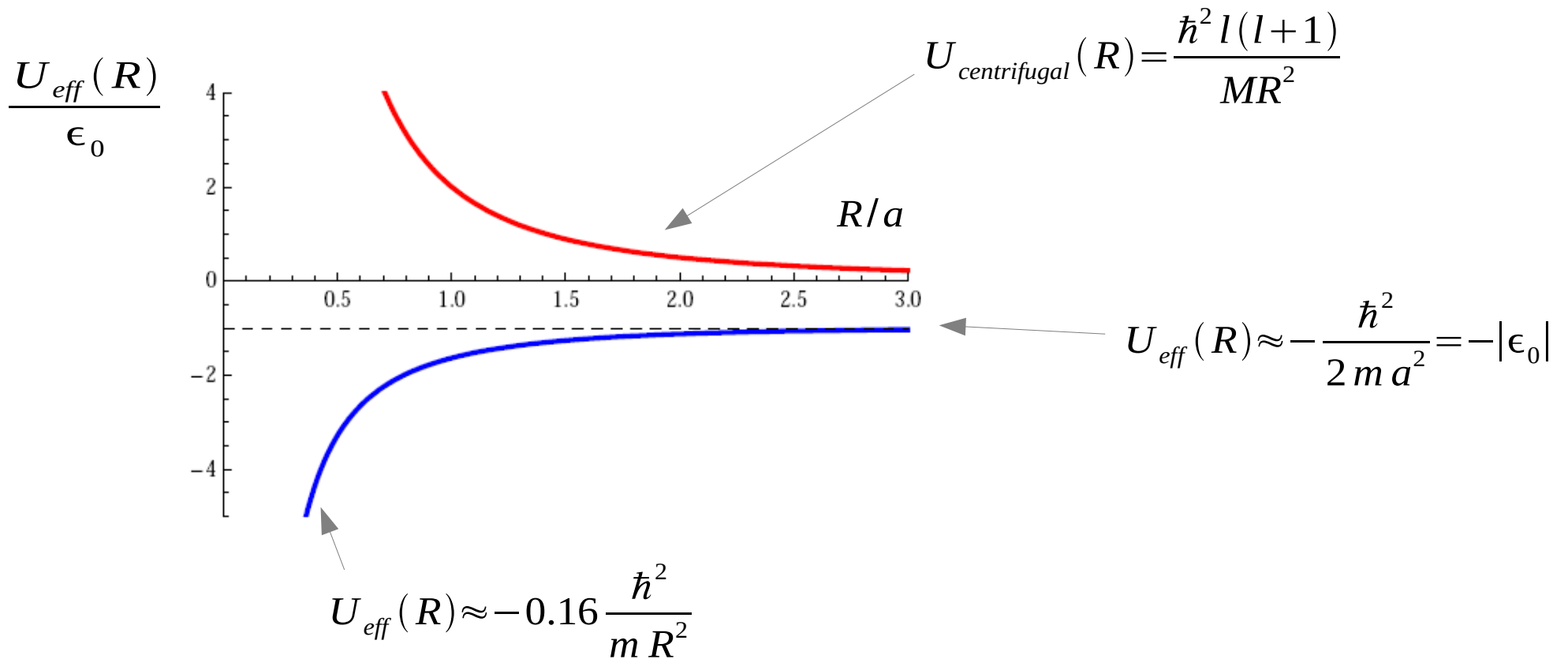
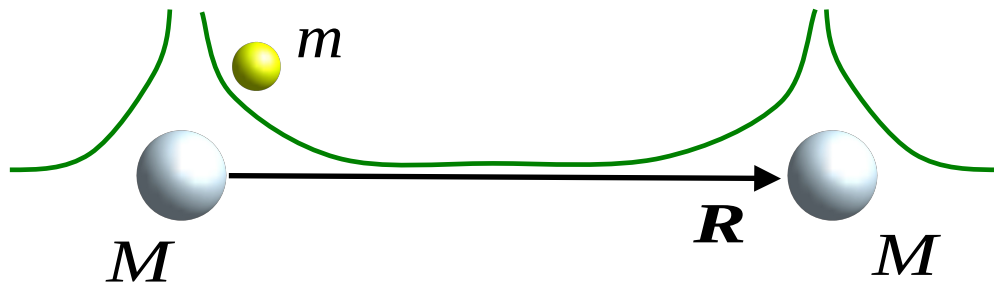
# Born-Oppenheimer picture



$$U_{eff}(R) \approx -\frac{\hbar^2}{2ma^2} = -|\epsilon_0|$$

$$U_{eff}(R) \approx -0.16 \frac{\hbar^2}{mR^2}$$

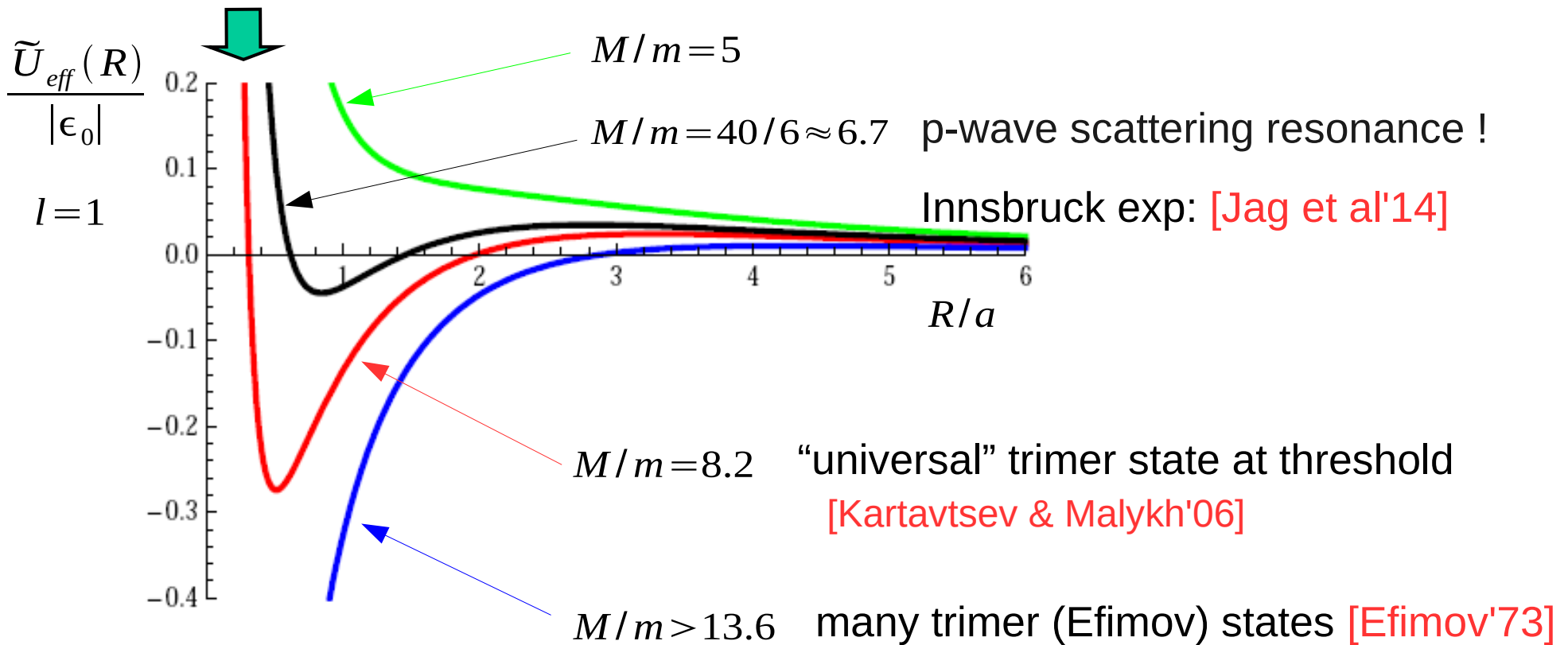
# Born-Oppenheimer picture



$$\left[ -\frac{\hbar^2}{M} \frac{\partial^2}{\partial R^2} + \tilde{U}_{eff}(R) \right] \chi(R) = E \chi(R)$$

$$\frac{\hbar^2}{MR^2} \left( l(l+1) - 0.16 \frac{M}{m} \right)$$

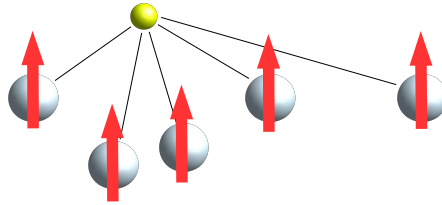
$$\tilde{U}_{eff}(R) = U_{eff}(R) + |\epsilon_0| + \frac{\hbar^2 l(l+1)}{MR^2}$$





# (N+1)-body problem

How many heavy fermions can be bound by a single light atom?



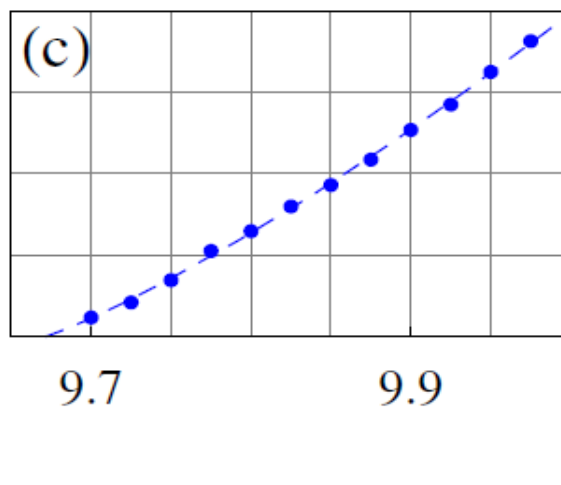
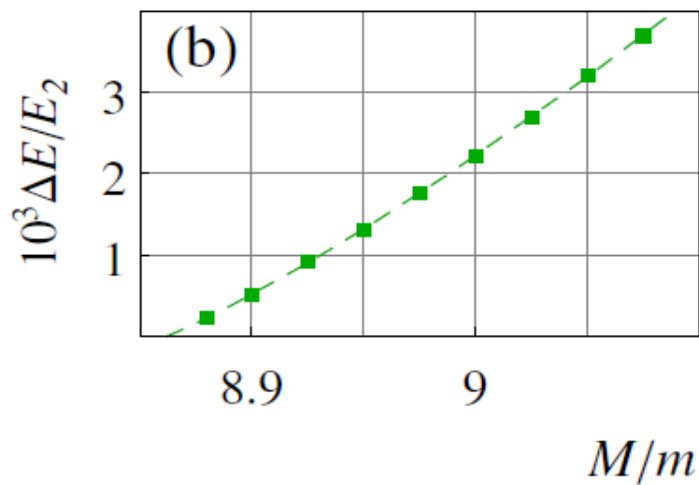
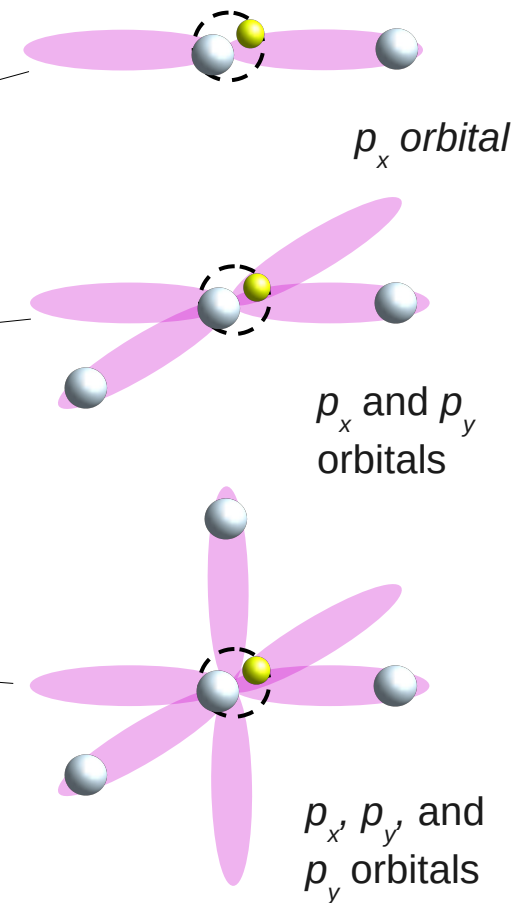
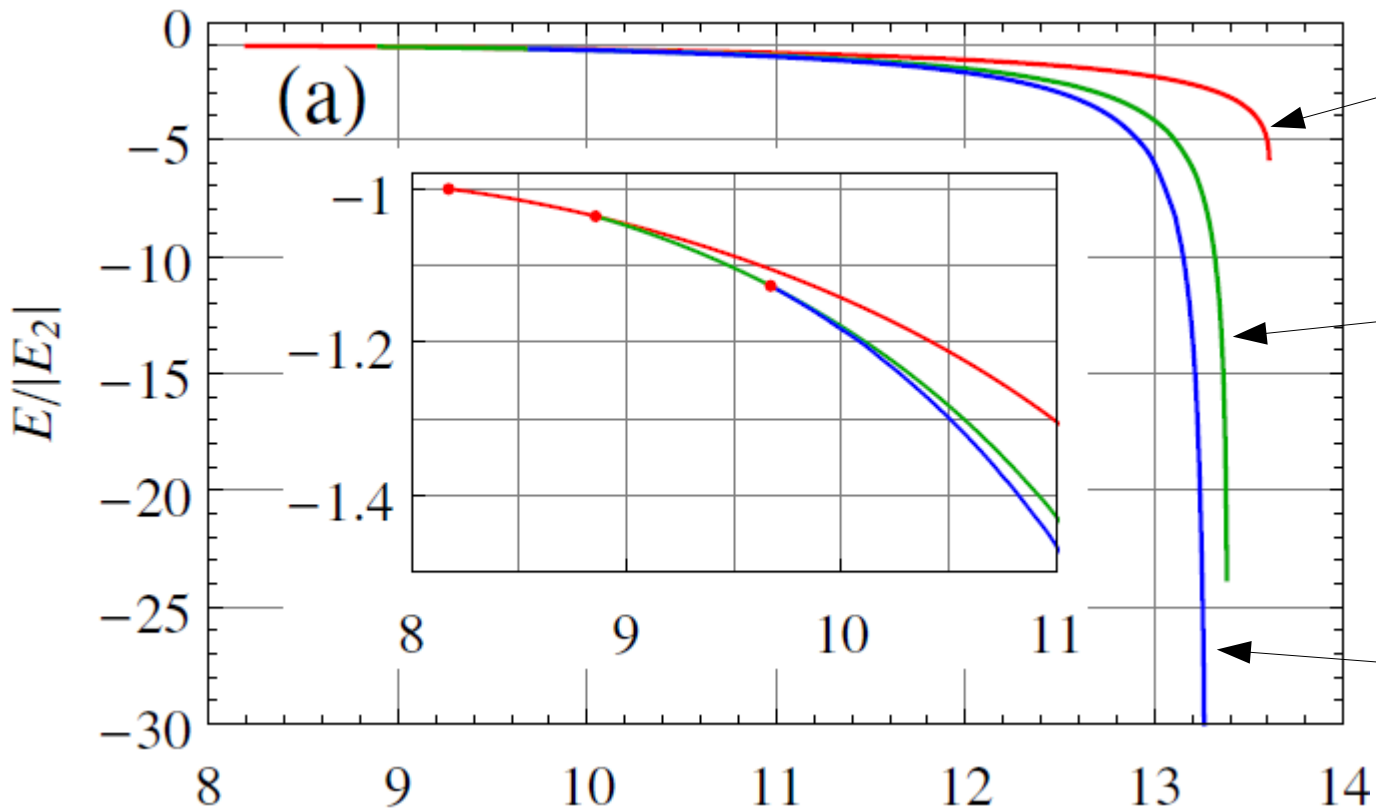
Kinetic energy of the heavy atoms  $\sim 1/M$

competes with

Attractive exchange potential of the light atom  $\sim 1/m$

Parameters of the **free-space zero-range** N+1-body problem:

- space dimension  $D$
- number of heavy atoms  $N$ 
  - mass ratio  $M/m$
- dimer size  $a$  (can be used as the length unit)



pentamer = closed  $p$ -shell



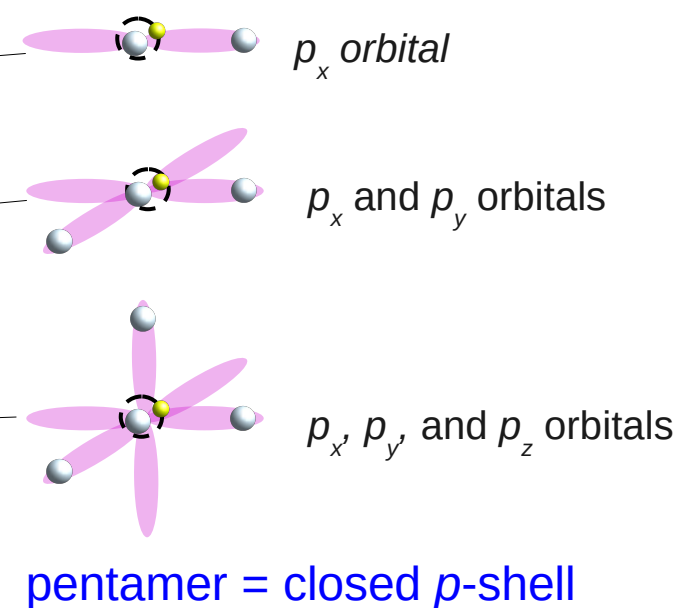
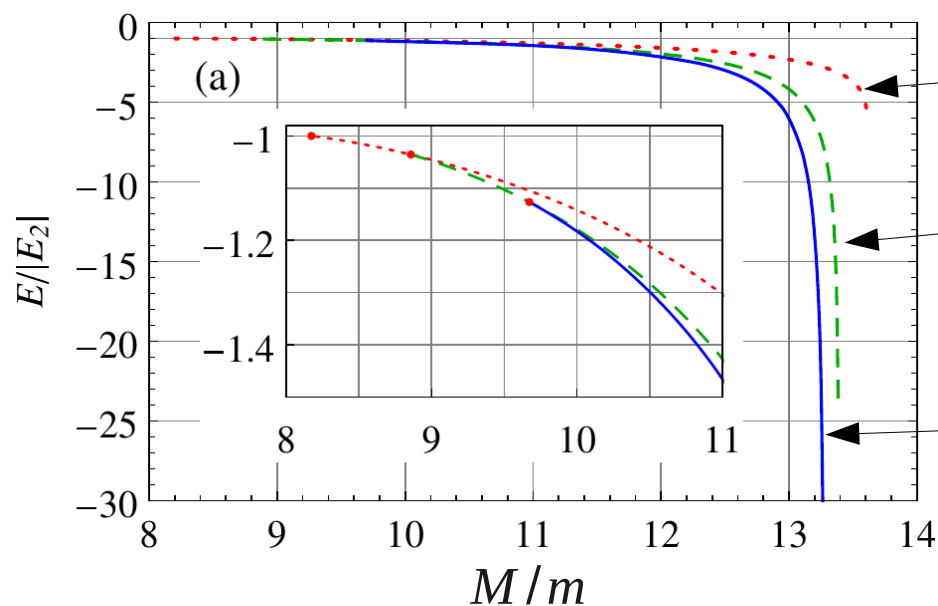
CONJECTURE:

No hexamer!

(requires justification)

# 3D trimer, tetramer, pentamer,...

	Symmetry $L^\pi$	appear at $M/m >$	Efimovian for $M/m >$
2+1 trimer	$1^-$	8.173 <small>Kartavtsev&amp;Malykh'06</small>	13.607 <small>Efimov'73</small>
3+1 tetramer	$1^+$	8.862(1) <small>Blume'12, Bazak&amp;DSP'17</small>	13.384 <small>Castin,Mora&amp;Pricoupenko'10</small>
4+1 pentamer	$0^-$	9.672(6) <small>Bazak&amp;DSP'17</small>	13.279(2) <small>Bazak&amp;DSP'17</small>
N+1-mer	?	?	?



# Physics at $a=\infty$ (& zero range)

Small-hyperradius behavior of the  $(N+1)$ -body wave function:

$$\left[ -\frac{\partial^2}{\partial R^2} - \frac{3N-1}{R^2} \frac{\partial}{\partial R} + \frac{s^2 - (3N/2 - 1)^2}{R^2} \right] \Psi(R) = 0$$



$$\Psi(R) \propto R^{-3N/2+1 \pm s}$$

$$s^2 > 0 \quad (s > 0)$$



$$\Psi(R) \propto R^{-3N/2+1+s}$$



“Universal” regime in the sense that one needs no three-body parameter

Non-Efimovian regime

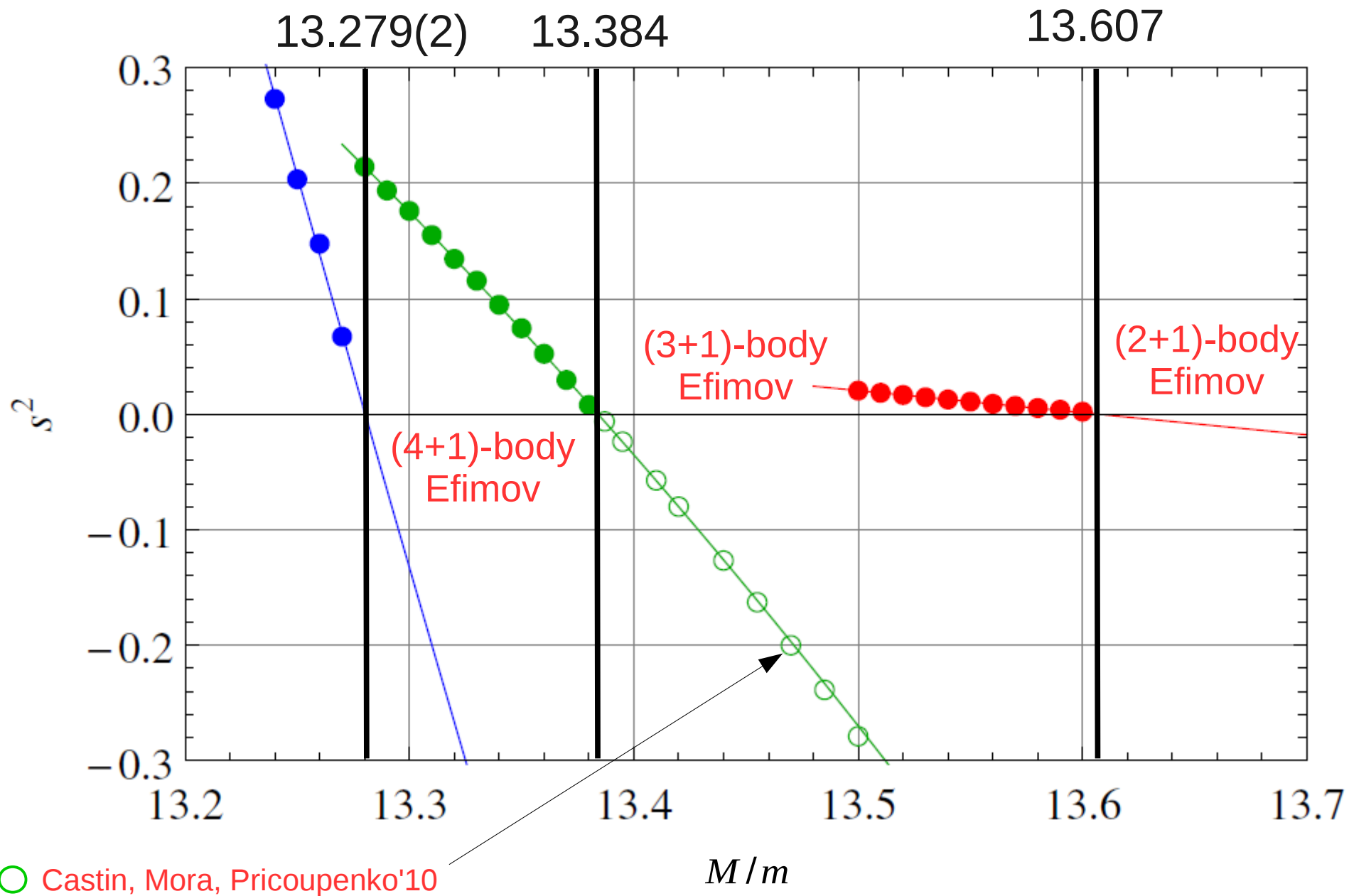
$$s^2 < 0 \quad (s = is_0)$$



$$\Psi(R) \propto R^{-3N/2+1} \sin(s_0 \ln R/R_0)$$



“Fall of a particle to the center in  $R^{-2}$  potential”. Infinite number of zeros of the wave function. Infinite number of trimer states. Efimov effect



A few words about low dimensions

$$\left[ -\frac{\hbar^2}{M} \frac{\partial^2}{\partial R^2} + \tilde{U}_{eff}(R) \right] \chi(R) = E \chi(R)$$

3D:  $\tilde{U}_{eff}(R) = U_{eff}(R) + |\epsilon_0| + \frac{\hbar^2 l(l+1)}{MR^2} \longrightarrow l=1 \rightarrow (M/m)_c = 8.2$

This is actually exact (not Born-Oppenheimer) number

different

2D:  $\tilde{U}_{eff}(R) = U_{eff}^{2D}(R) + |\epsilon_0| + \frac{\hbar^2 (l^2 - 1/4)}{MR^2} \longrightarrow$  Rough guess:

$$(M/m)_c^{2D} \approx \frac{l^2 - 1/4}{l(l+1)} (M/m)_c^{3D} = 3.1$$

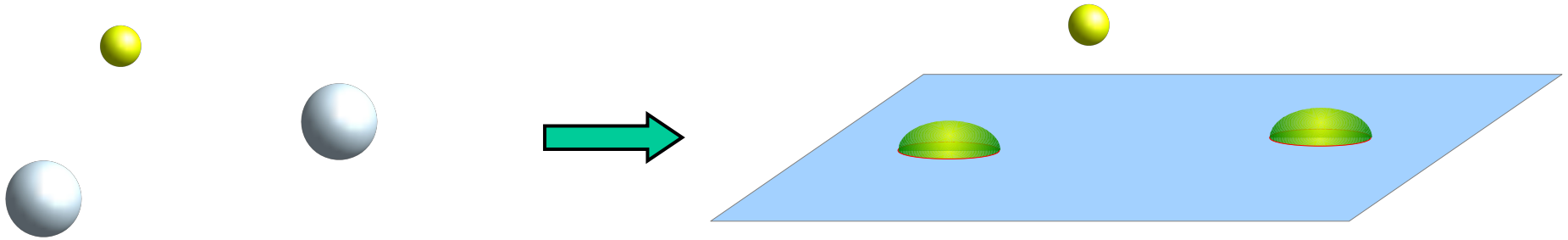
Exact ratio  $(M/m)_c^{2D} = 3.3$  [Pricoupenko & Pedri'10]

Centrifugal force weaker in 2D  $\rightarrow$  p-wave resonance for smaller mass ratio!

... and in 1D  $(M/m)_c^{1D} = 1$  exactly!

# Can we make a bound Li-K-K trimer state?

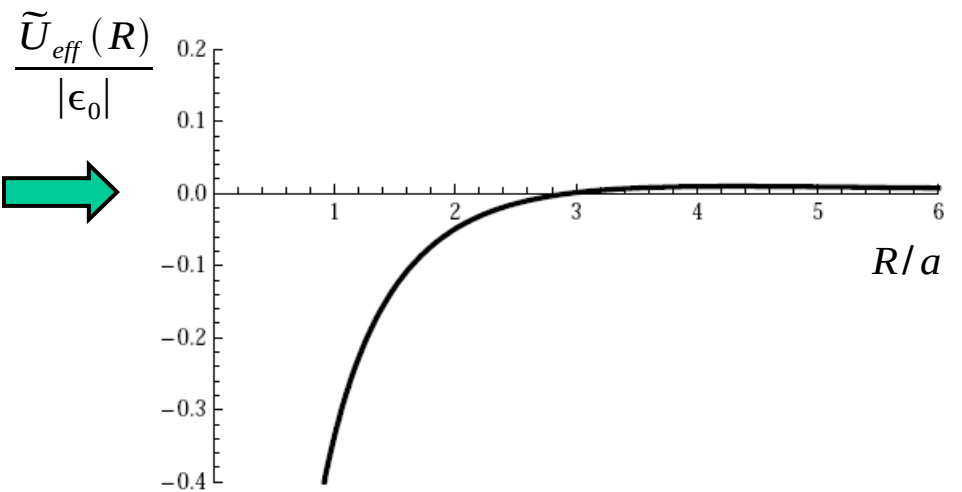
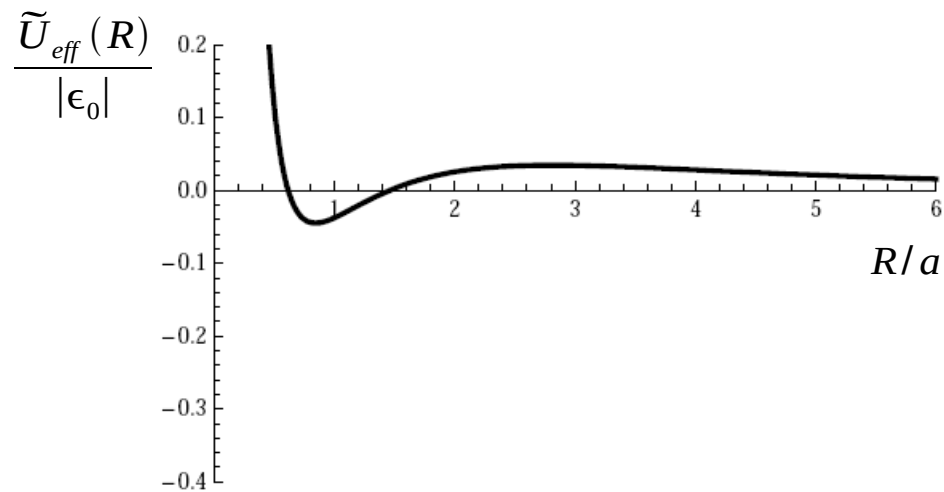
[Levinsen et al'09]



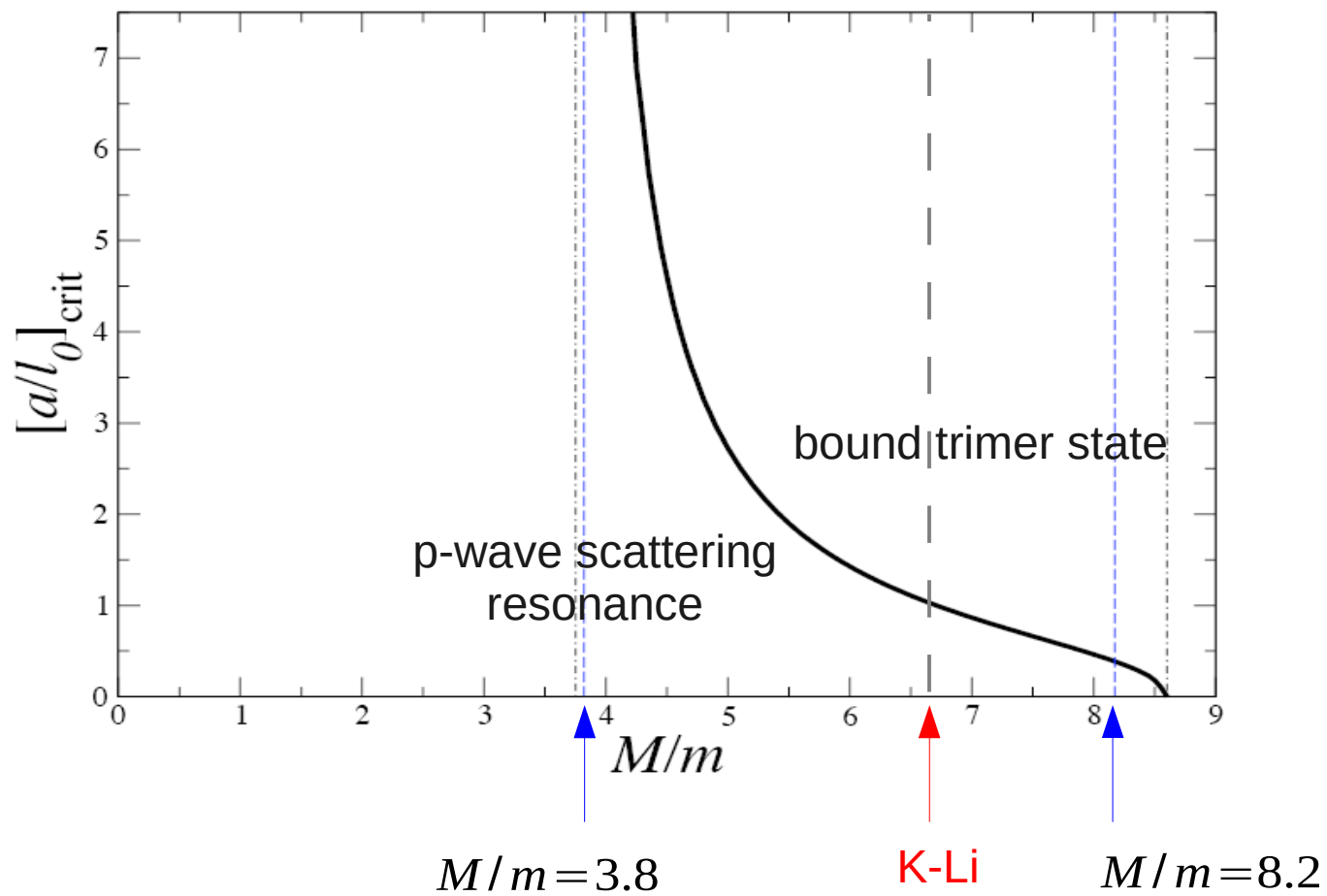
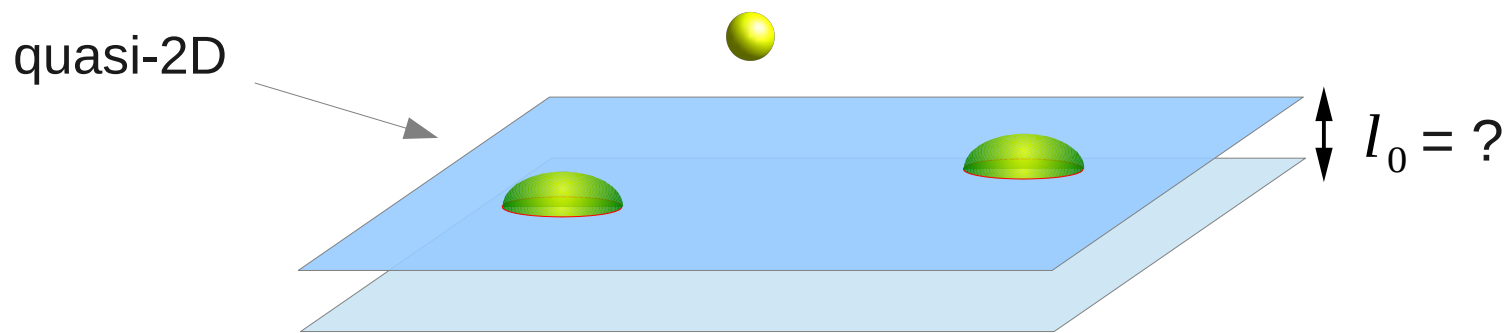
$$U_{\text{centrifugal}}(R) = \frac{\hbar^2 l(l+1)}{MR^2}$$



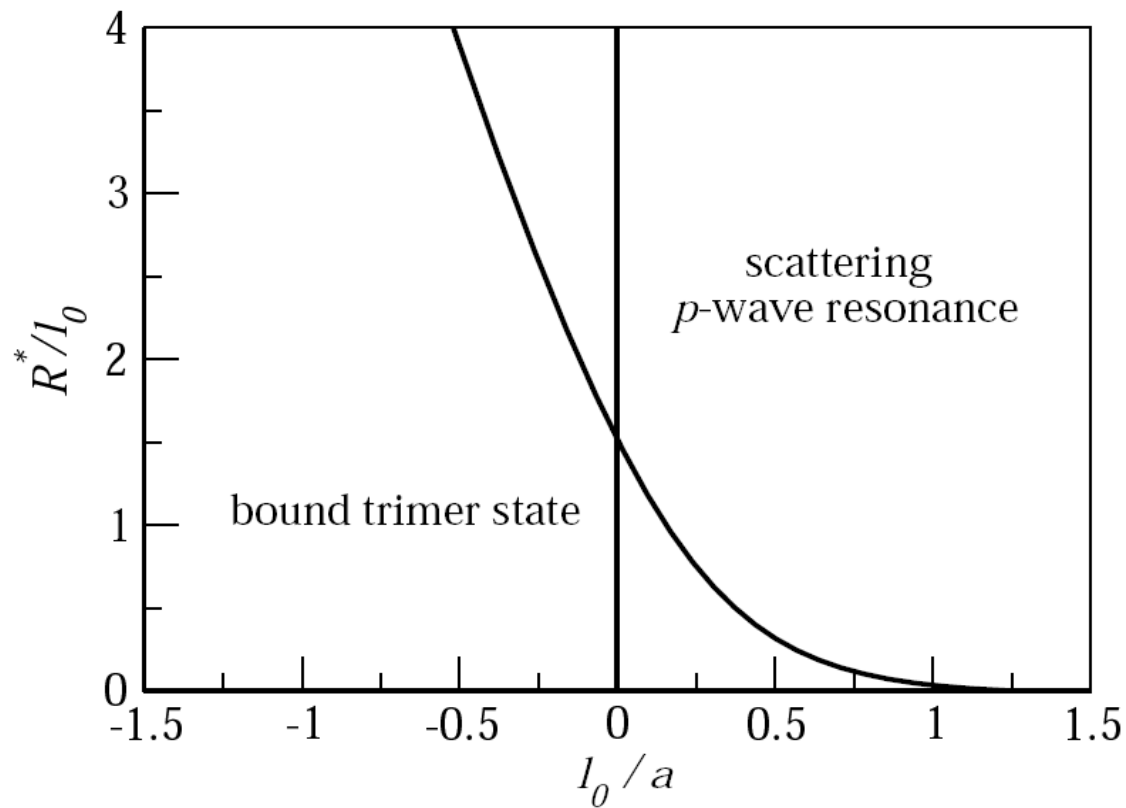
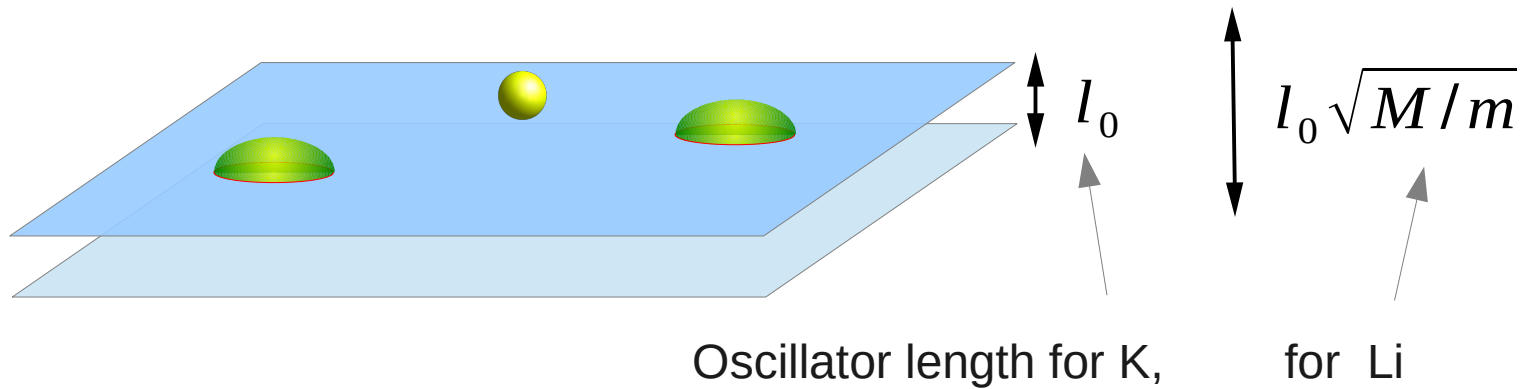
$$U_{\text{centrifugal}}(R) = \frac{\hbar^2 (l^2 - 1/4)}{MR^2}$$







# Quasi-2D – quasi-2D case $\omega_{Li} = \omega_K$



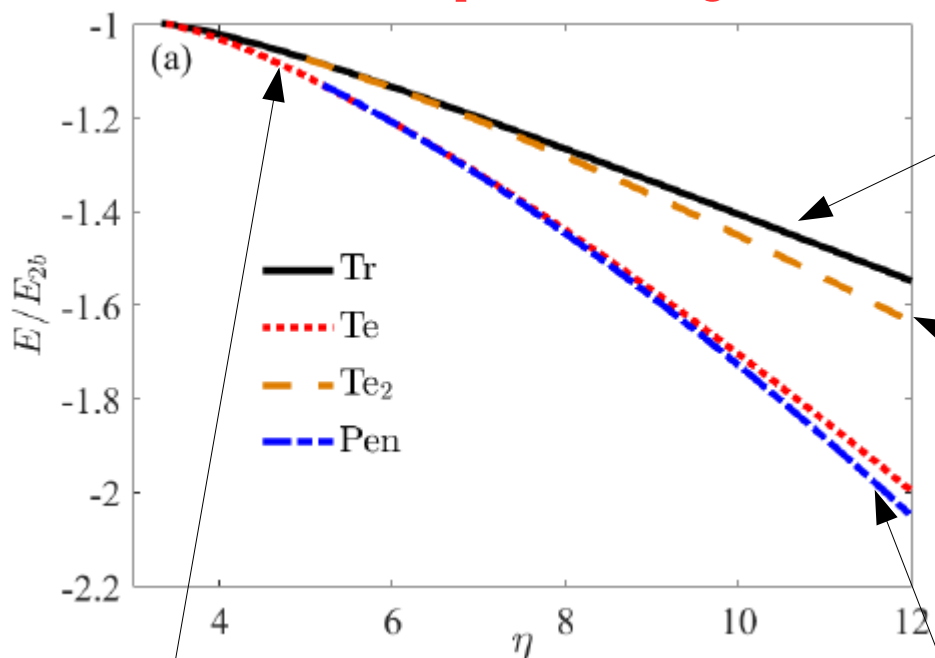
[Levinsen et al'09]

Expect similar effect for Li-Cr and K-Dy!

Exact ratio for the trimer formation in 2D  $(M/m)_c^{2D} = 3.3$  [Pricoupenko & Pedri'10]

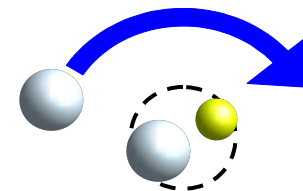
# 2D trimer, tetramer, pentamer...

[Liu & Peng & Cui'22]



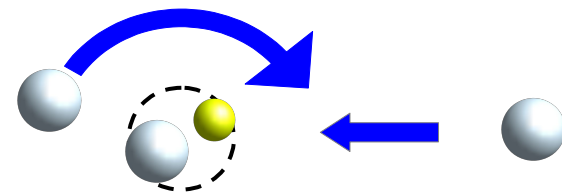
L=1 trimer  $(M/m)_c = 3.33$

[Pricoupenko & Pedri'10]



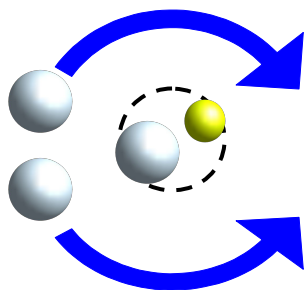
L=1 tetramer  $(M/m)_c^{2D} = 5.0$

[Levinsen & Parish'13]



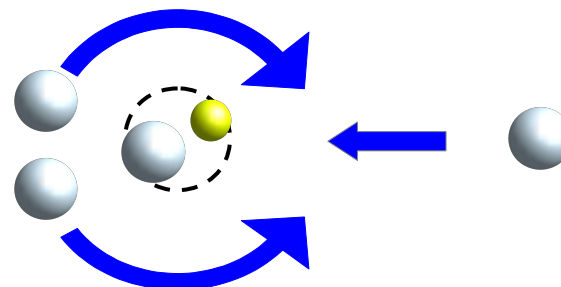
L=0 tetramer  $(M/m)_c = 3.38$

[Liu & Peng & Cui'22]



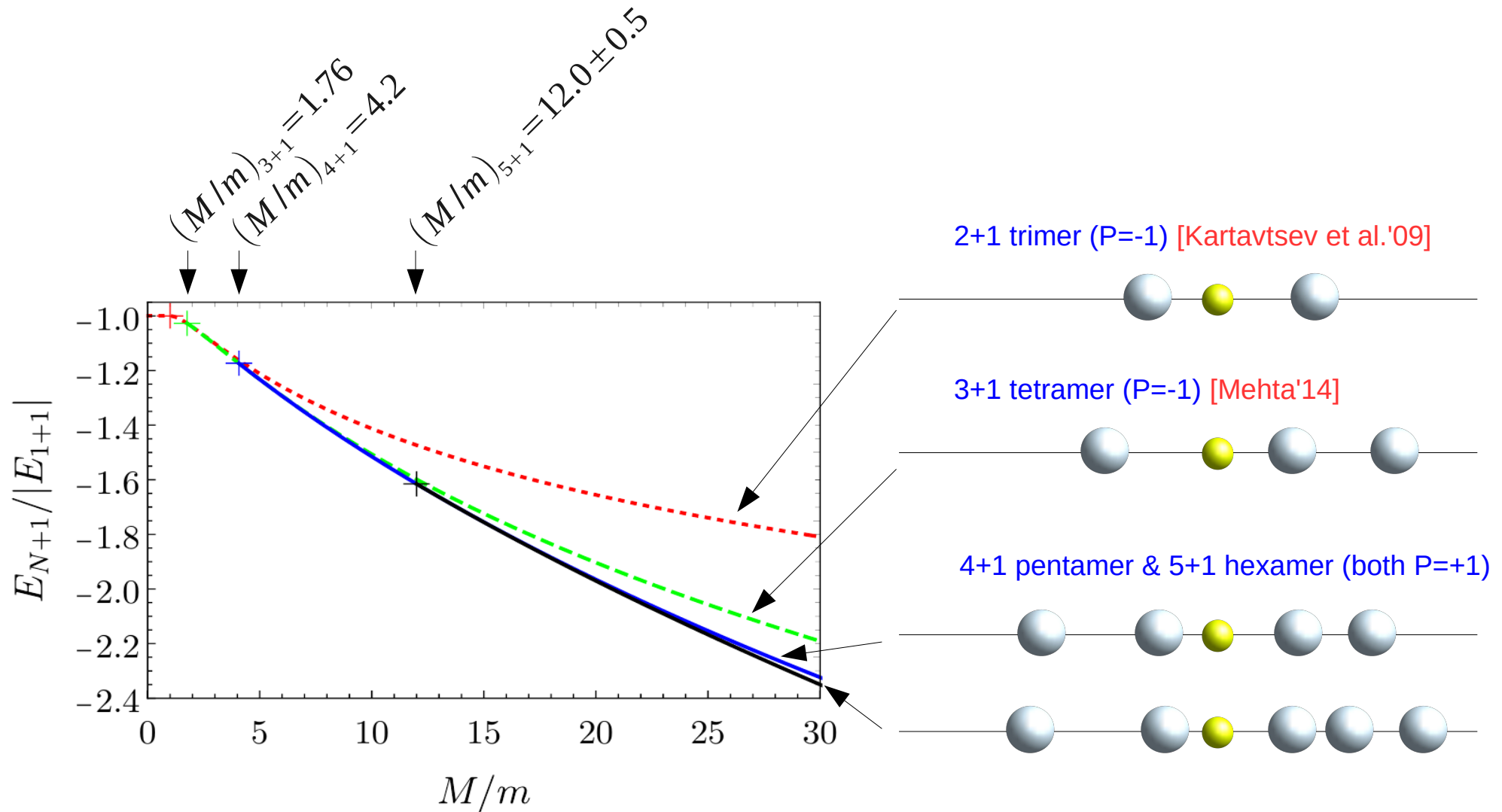
L=0 pentamer  $(M/m)_c = 5.14$

[Liu & Peng & Cui'22]



# 1D trimer, tetramer...(exact)

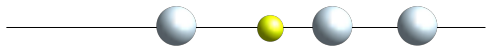
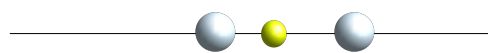
A. Tononi, J. Givois, DSP, Phys. Rev. A **106**, L011302 (2022)



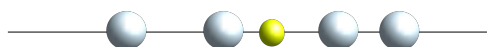
1D

2D

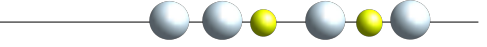
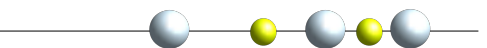
3D



Exact STM:



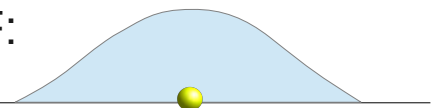
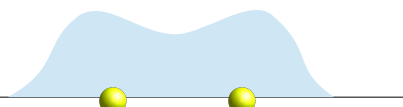
PRA 106, L011302 (2022)



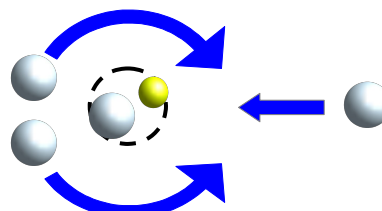
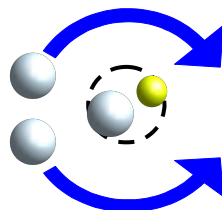
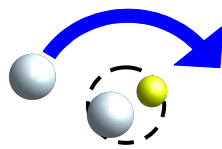
?

PRA 106, L011302 (2022)

MF:

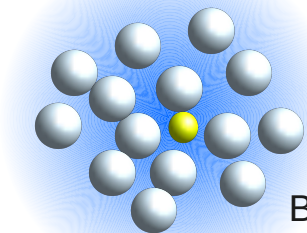
Binding for  $\frac{M}{m} > \frac{\pi^2}{36} N^3$ 

SciPost Phys. 14, 091 (2023)

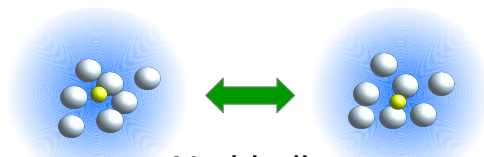


?

MF+BMF

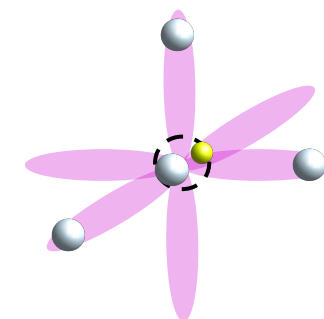
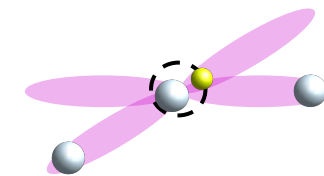
Binding for  $\frac{M}{m} > 1.074N^2$ 

SciPost Phys. 17, 050 (2024)



No binding

Few-Body Syst 65, 80 (2024)



?

 $M/m \approx 13$  Efimov threshold

?

MF or exact?

 $M/m$  or  $N$

Pitaevskii-Rosch scaling symmetry

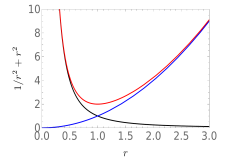
Classical (mean-field) Gross-Pitaevskii energy functional for 2D bosons (similar for 2D Fermi mixtures)

$$E_{\text{MF}}(\Psi, \Psi^*) = (1/2) \int d^2\rho [|\nabla_{\rho}\Psi(\rho,t)|^2 + g|\Psi(\rho,t)|^4]$$

$$V(\rho) = g\delta^2(\rho) \quad : \quad V(\lambda\rho) = \lambda^{-2}V(\rho)$$

dimensionless

$$V(r) = \beta/r^2 \quad - \text{other example (any dimension)}$$



$$E_{\text{MF}}(\Psi, \Psi^*) = (1/2) \int d^2\rho [|\nabla_{\rho}\Psi(\rho,t)|^2 + g|\Psi(\rho,t)|^4 + \omega^2\rho^2|\Psi(\rho,t)|^2] \longrightarrow \text{Undamped breathing mode with frequency } 2\omega \text{ independent of interaction}$$

[Pitaevskii'1996;Pitaevskii&Rosch'1997]

The model  $E_{\text{MF}}(\Psi, \Psi^*)$  (with  $V(\rho) = g\delta^2(\rho)$ ) does not survive quantization, but remains a good approximation in some cases (spoiler : too good)!

To « survive quantization » means to stay valid for the same quantum Lagrangian.

no survival → quantum anomaly (smoking gun = deviation from  $2\omega$ )

Quantum models featuring PR symmetry : 3D unitary (non-Efimovian) gases, 1D Tonks gas ( $V_{1D}(x) = \infty\delta(x)$ ),  $1/r^2$ -models in any dimension

Example of a Quantum Anomaly in the Physics of Ultracold Gases (2D bosons)

Maxim Olshanii,<sup>1,2</sup> H el ene Perrin,<sup>2</sup> and Vincent Lorent<sup>2</sup>

Quantum Anomaly, Universal Relations, and Breathing Mode of a Two-Dimensional Fermi Gas (2D BCS-BEC crossover)

Johannes Hofmann\*

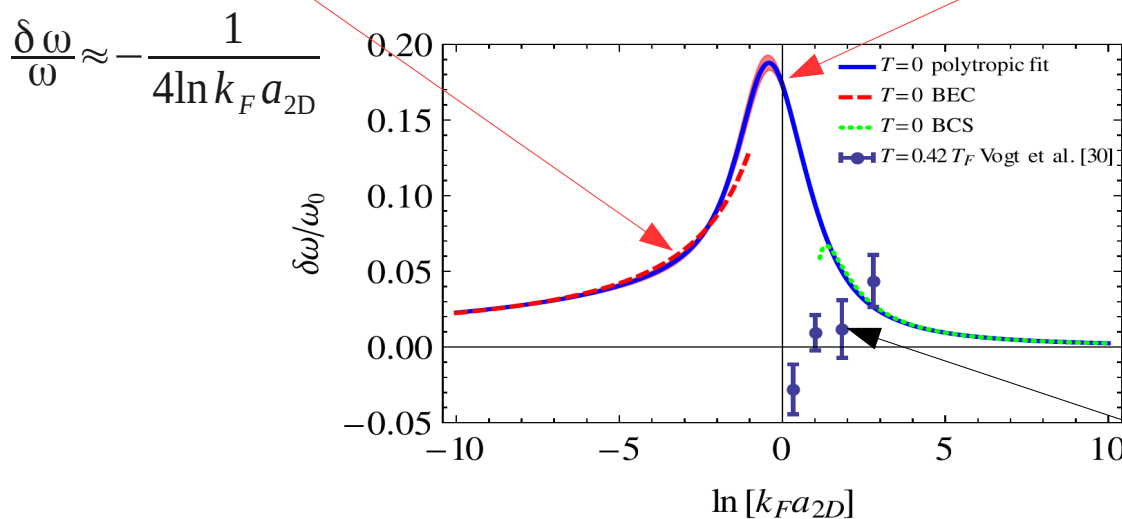
QM Hamiltonian with the Bethe-Peierls boundary condition  $\psi(\rho) \xrightarrow{\rho \rightarrow 0} C \ln\left(\frac{\rho}{a_{2D}}\right)$



Hydrodynamic equations with local-density equation of state

Perturbative for bosons [Schick'1971, Popov'1972, Mora&Castin'2009]

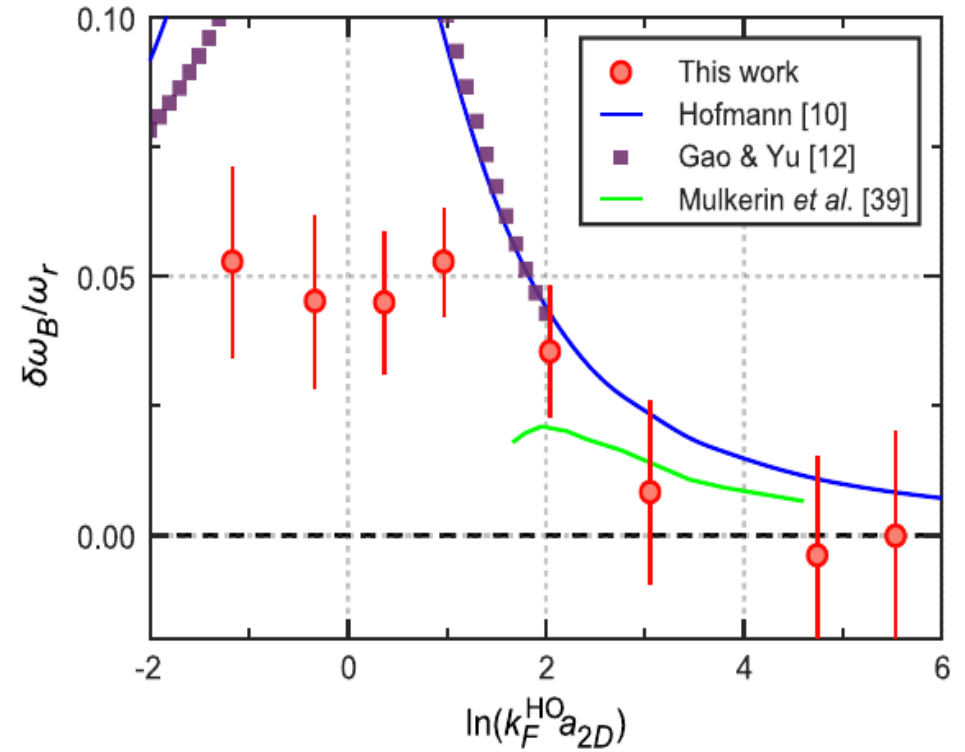
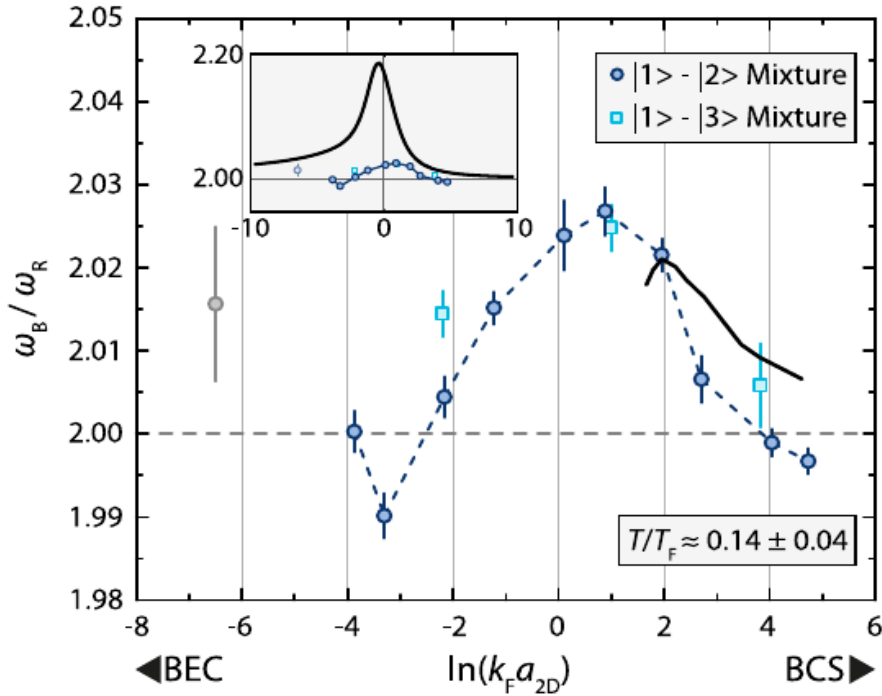
Monte-Carlo in the 2D BCS-BEC crossover [Bertaina&Giorgini'2011]



[Vogt et al.'2012] (Cambridge)



# FERMIONS

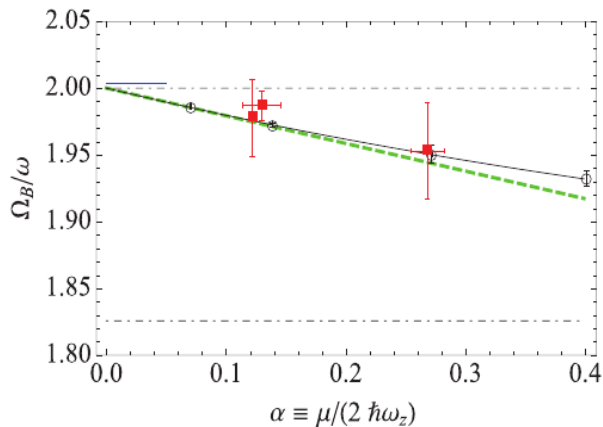


[Holten et al.'2018] (Heidelberg)

see references  
for theory works

[Peppler et al.'2018] (Melbourne)

# BOSONS



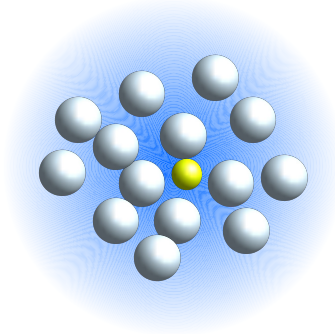
[Merloti et al.'2013] (Paris 13)

Problems :

- small relative shift
- finite-T
- third direction (2D-3D crossover)
- trap anisotropy and inharmonicity

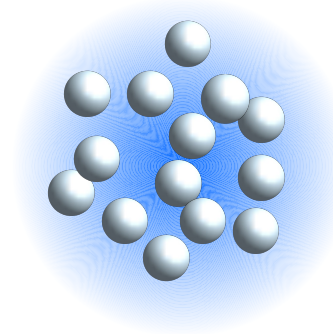
## Self-bound

2D N+1 fermionic clusters



## Self-bound

2D bosons



## If no trap

MF predicts  $2\omega=0$  for the breathing mode frequency

Self-evaporation, « self-cleaning »

relevant, not analyzed yet

No trap !

Problems :

- ~~small relative shift~~
- ~~finite-T~~
- third direction (2D-3D crossover)
- ~~trap anisotropy and inharmonicity~~

Next : 2D attractive bosons

