The (Lattice) QCD thermal transition as a function of the parameters

Owe Philipsen

- I. Chiral phase transition in the massless limit
- II. Evidence for chiral spin symmetric intermediate temperature regime
- III. Cold+dense, qualitatively: nuclear liquid gas transition + quarkyonic matter

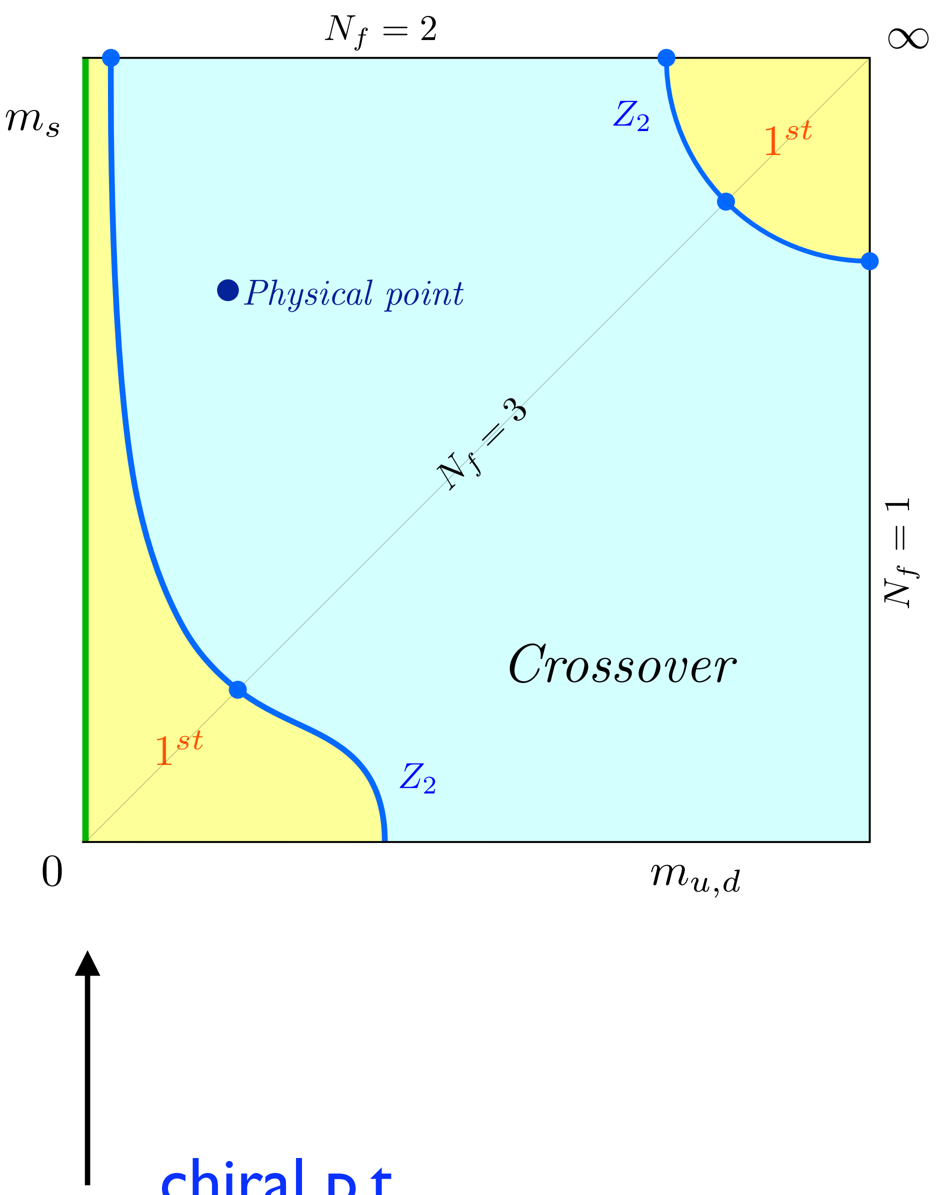
Tool: QCD away from physical point, study limits + parameter dependence: constraints!

CRC-TR 211 Strong-interaction matter under extreme conditions
EL EM EN TS
GOETHE UNIVERSITÄT FRANKFURT AM MAIN
NIC

I. Nature of the QCD thermal transition at zero density

$N_f = 2 + 1$

deconfinement p.t.:
breaking of global $Z(3)$ symmetry

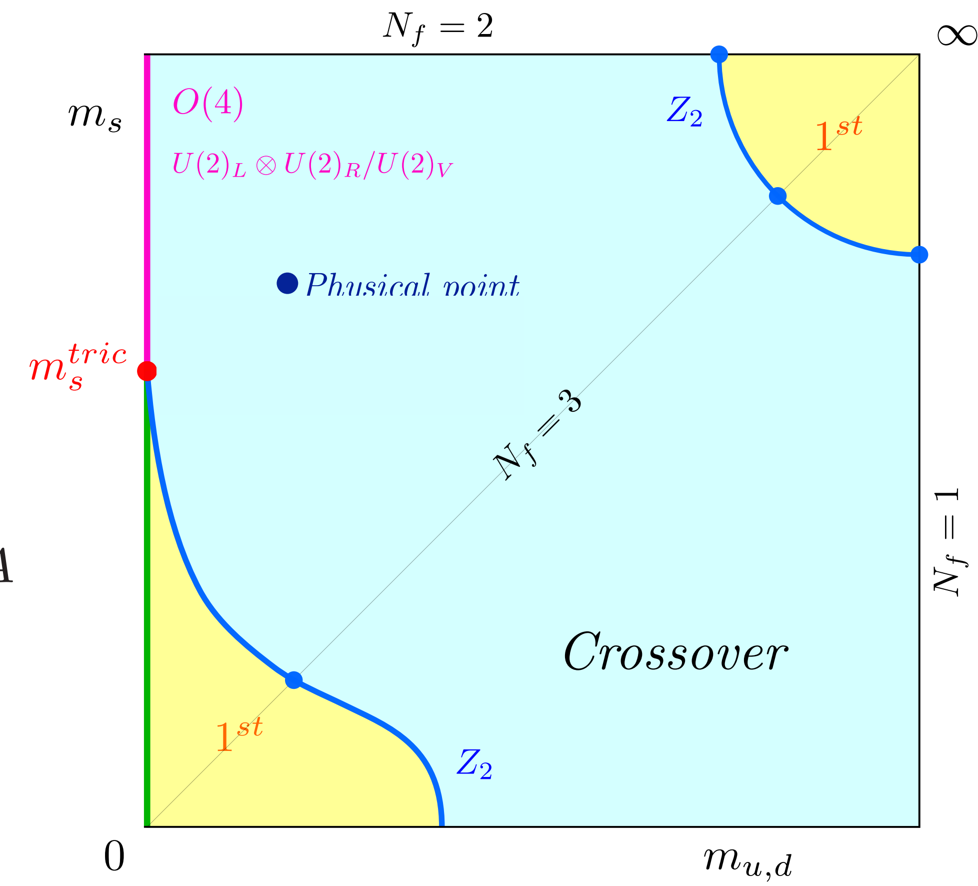


[Pisarski, Wilczek, PRD 84]:
(Linear sigma model in 3d)

$N_f = 2$ depends on $U(1)_A$

restored broken

$N_f \geq 3$ 1st order



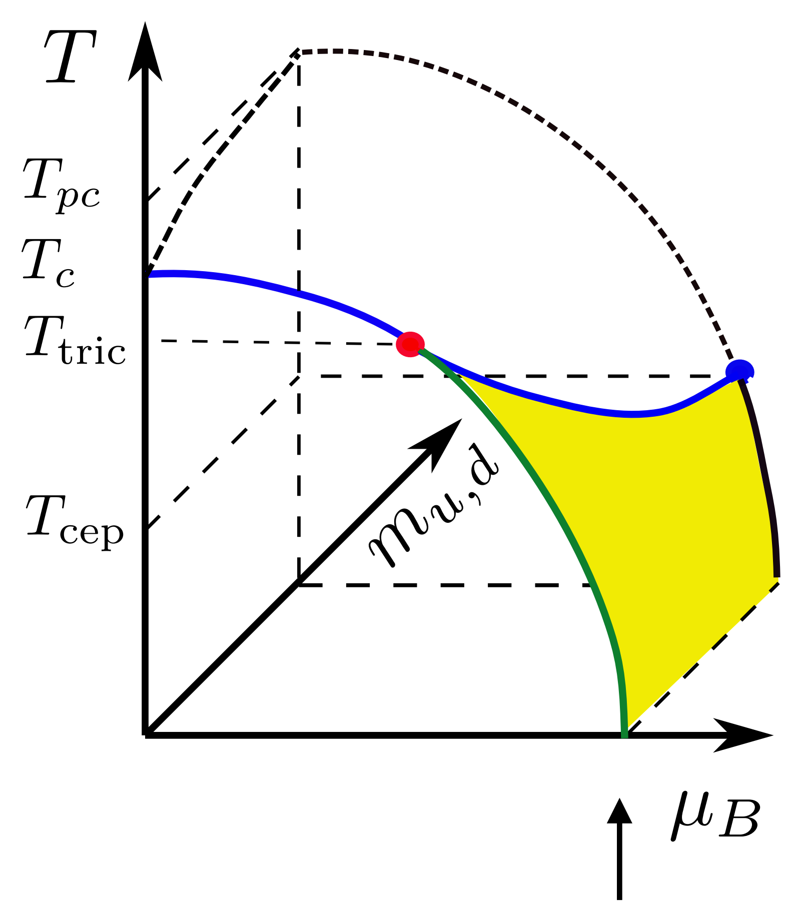
chiral p.t.
restoration of global symmetry in flavour space

$SU(2)_L \times SU(2)_R \times U(1)_A$

anomalous

Including chemical potential

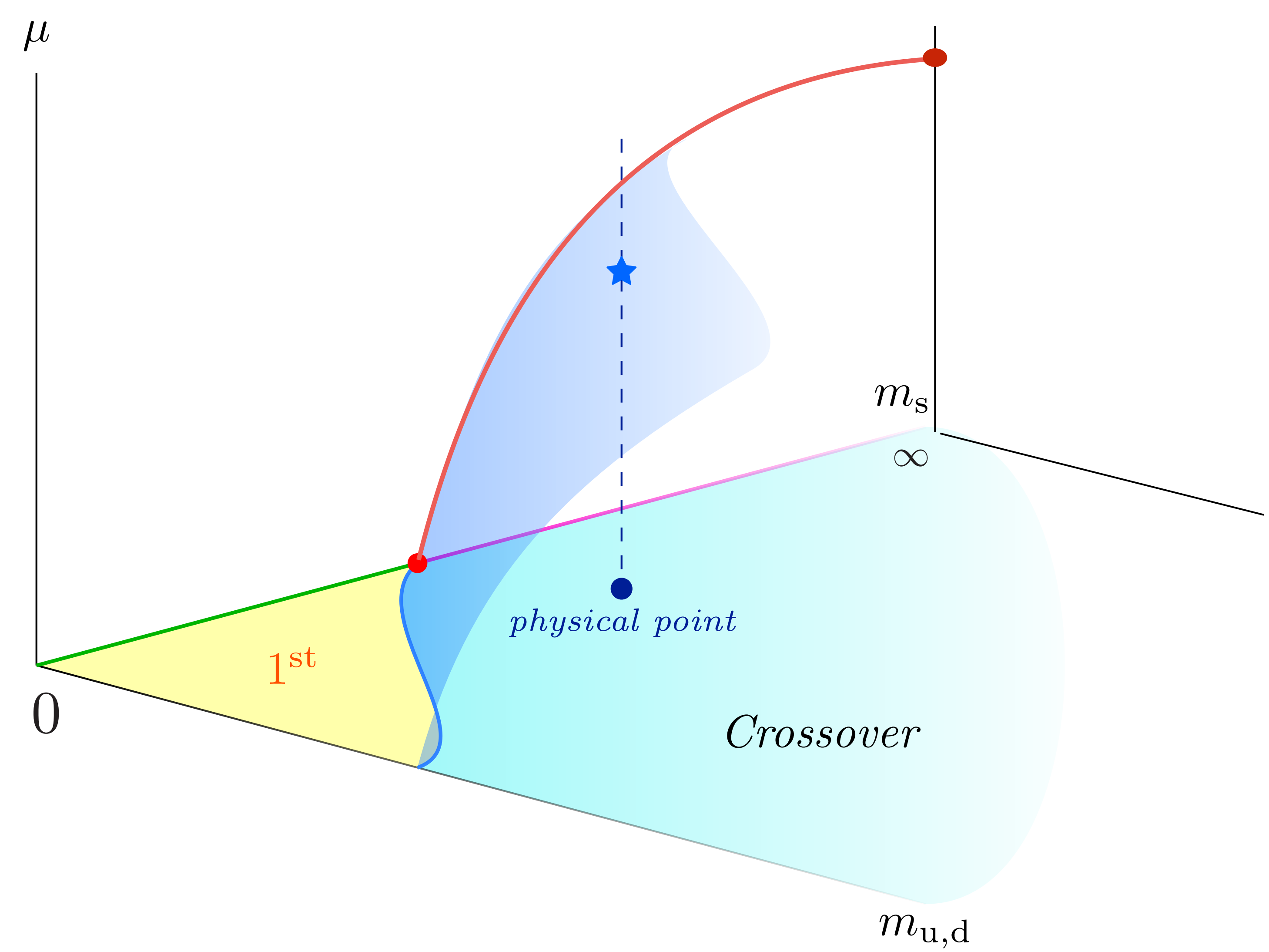
$N_f = 2$



Model predictions, no full QCD information

$N_f = 2 + 1$

[edited from Sciarra, PhD thesis 2016]

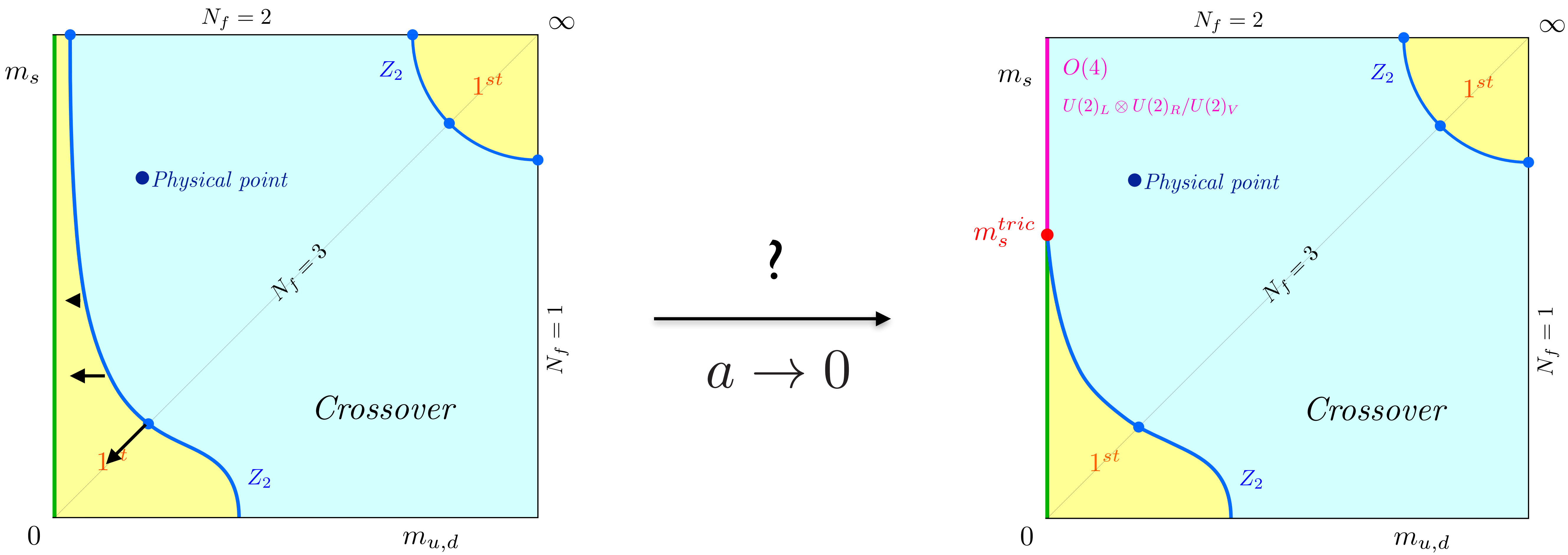


[Stephanov, Rajagopal, Shuryak PRL 98]: (models + early lattice results)

"As m_s is reduced from infinity, the tricritical point ... moves to lower μ until it reaches the T-axis and can be identified with the tricritical point in the (T, m_s)-plane"

The nature of the QCD chiral transition at zero density

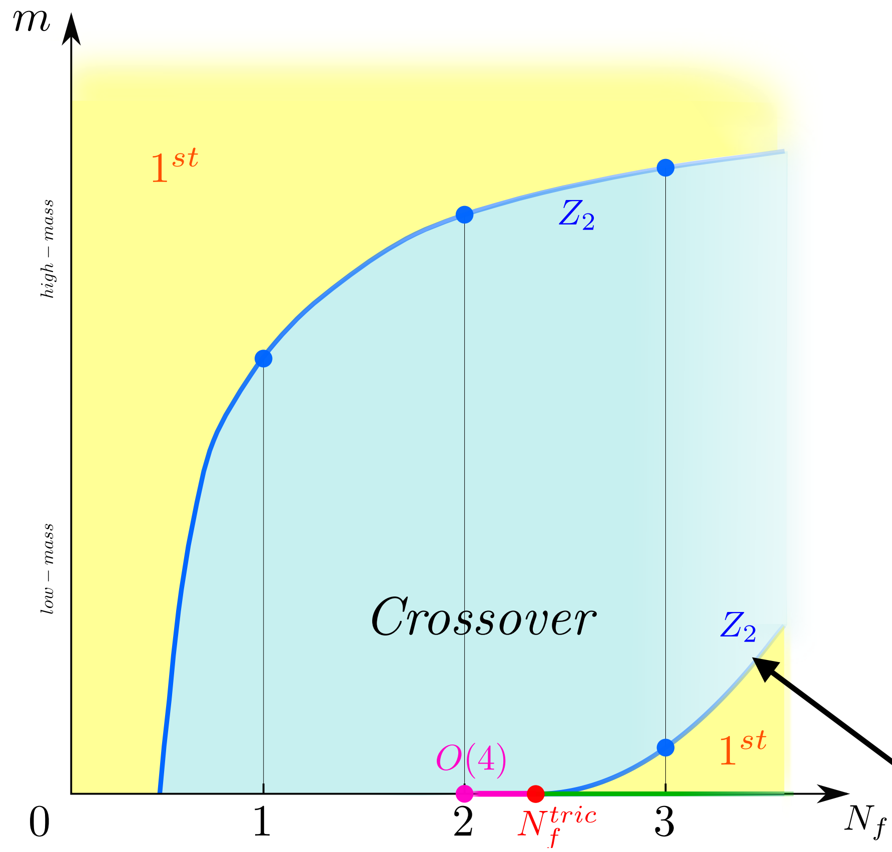
…is elusive, massless limit not simulable!

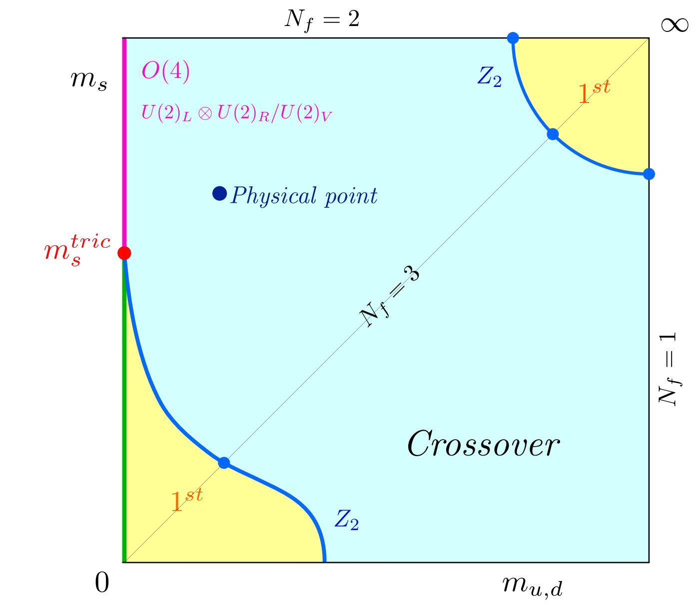


- Coarse lattices with unimproved actions: 1st order for $N_f = 2, 3$
- 1st order region shrinks rapidly as $a \to 0$, no 1st order for improved staggered actions
- Apparent contradictions between different lattice actions?

Details and reference list: [O.P., Symmetry 13, 2021]

Different view point: mass degenerate quarks





$$Z(N_\mathrm{f}, g, m) = \int \mathcal{D}A_\mu \; (\det M[A_\mu, m])^{N_\mathrm{f}} \; e^{-\mathcal{S}_\mathrm{YM}[A_\mu]}$$

$$N_\mathrm{f}^c(am) = N_\mathrm{f}^\mathrm{tric} + \mathcal{B}_1 \cdot (am)^{2/5} + \mathcal{O}\big((am)^{4/5}\big)$$

- Consider analytic continuation to continuous N_f
- Tricritical point guaranteed to exist if there is 1st order at any N_f
- Known exponents for critical line entering tric. point!
- Continuation to $a \neq 0$: Z(2) surface ends in tricritical line

[Cuteri, O.P., Sciarra PRD 18]

Different view point: mass degenerate quarks

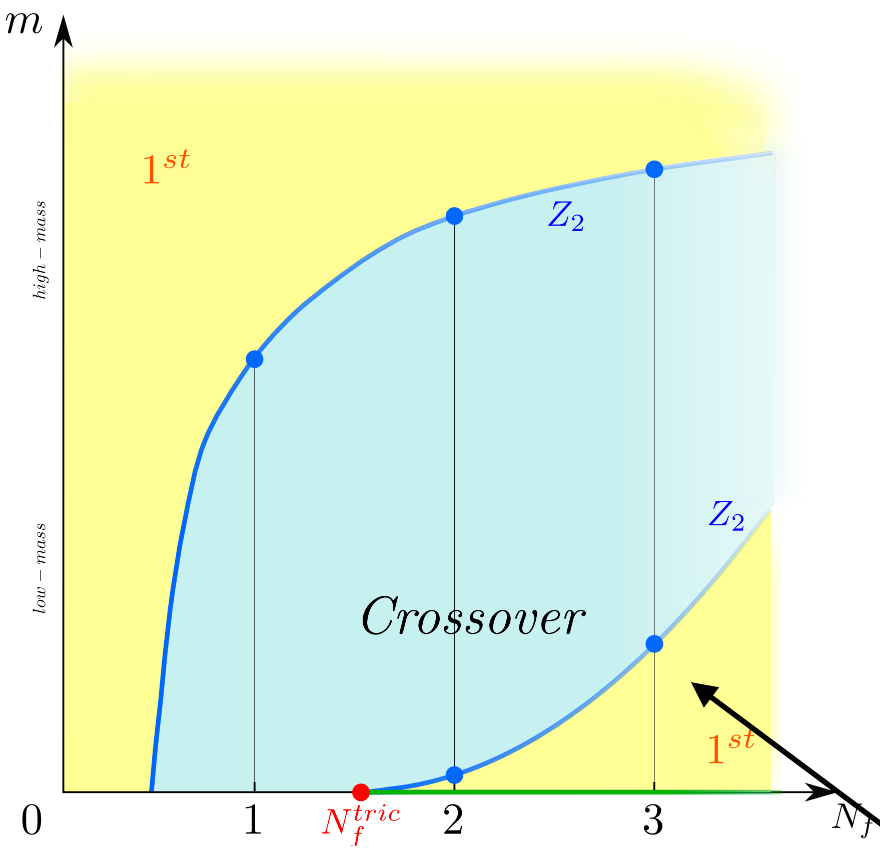
$$Z(N_{\rm f}, g, m) = \int \mathcal{D}A_\mu \; (\det M[A_\mu, m])^{N_{\rm f}} \; e^{-\mathcal{S}_{\rm YM}[A_\mu]}$$

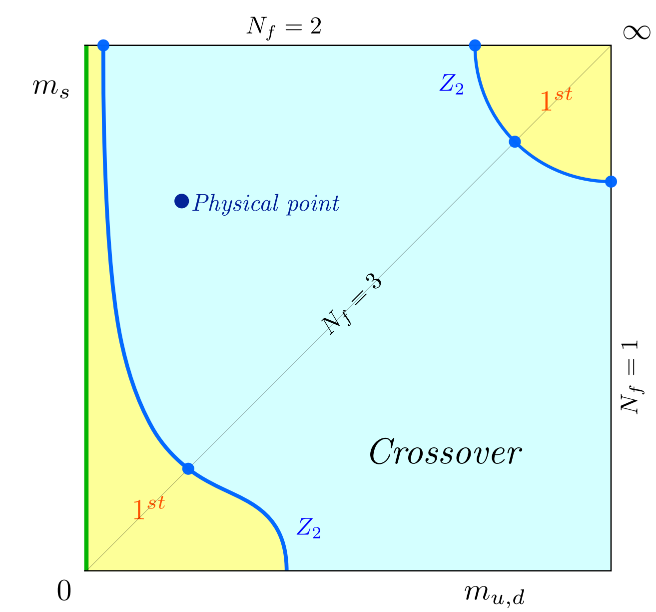
$$N_{\rm f}^c(am) = N_{\rm f}^{\rm tric} + \mathcal{B}_1 \cdot (am)^{2/5} + \mathcal{O}\big((am)^{4/5}\big)$$

- Consider analytic continuation to continuous N_f
- Tricritical point guaranteed to exist if there is 1st order at any N_f
- Known exponents for critical line entering tric. point!
- Continuation to $a \neq 0$: Z(2) surface ends in tricritical line

[Cuteri, O.P., Sciarra PRD 18]

Different view point: mass degenerate quarks





$$Z(N_\mathrm{f}, g, m) = \int \mathcal{D}A_\mu \, (\det M[A_\mu, m])^{N_\mathrm{f}} \, e^{-\mathcal{S}_{\mathrm{YM}}[A_\mu]}$$

$$N_\mathrm{f}^c(am) = N_\mathrm{f}^{\mathrm{tric}} + \mathcal{B}_1 \cdot (am)^{2/5} + \mathcal{O}\big((am)^{4/5}\big)$$

- Consider analytic continuation to continuous N_f
- Tricritical point guaranteed to exist if there is 1st order at any N_f
- Known exponents for critical line entering tric. point!
- Continuation to $a \neq 0$: Z(2) surface ends in tricritical line

[Cuteri, O.P., Sciarra PRD 18]

Methodology to determine order of transition

Finite size scaling of generalised cumulants

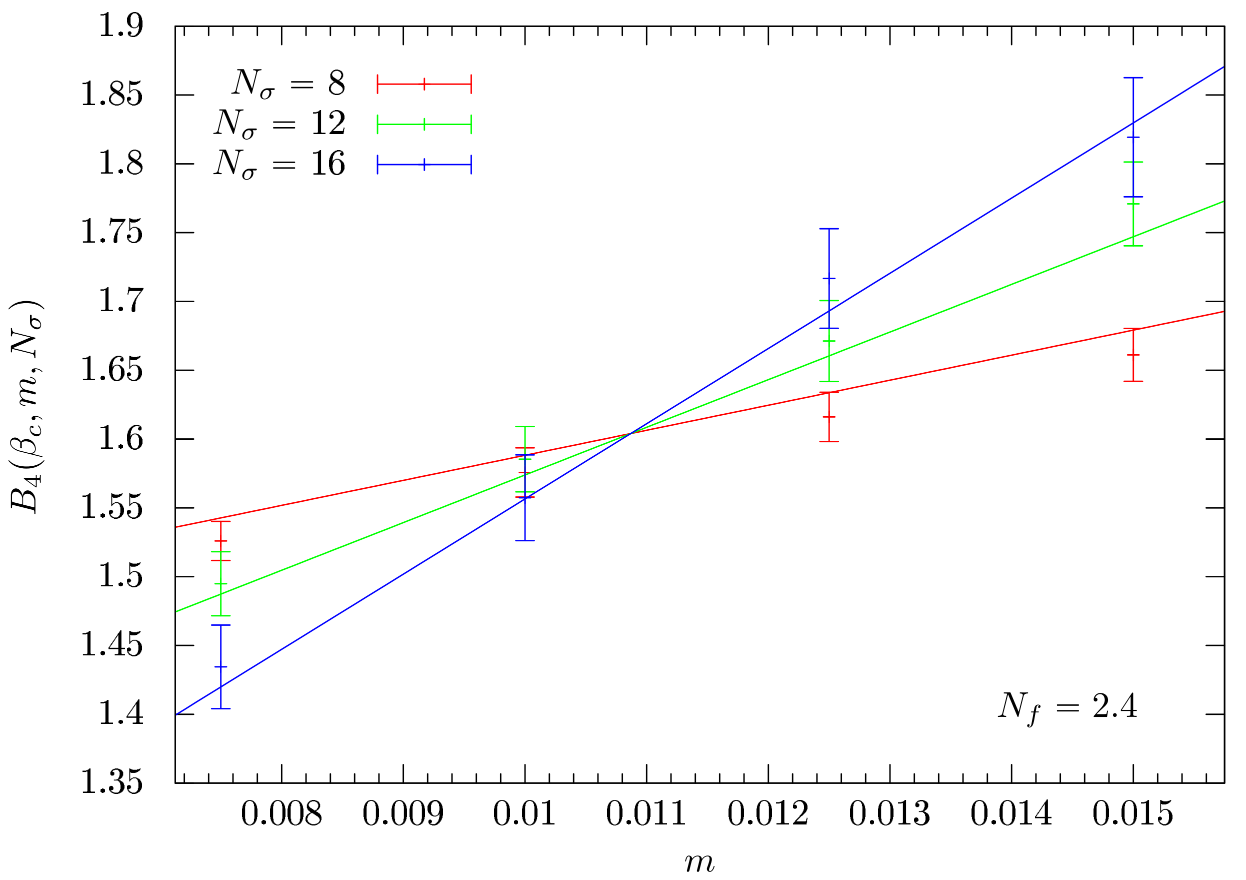
$$B_n = \frac{\langle (\bar{\psi}\psi - \langle \bar{\psi}\psi \rangle)^n \rangle}{\langle (\bar{\psi}\psi - \langle \bar{\psi}\psi \rangle)^2 \rangle^{n/2}}$$

Standard staggered fermions, bare parameters: β, am, N_f, N_τ

(Pseudo-critical) phase boundary: $B_3 = 0$ 3d manifold

Second-order 3d Ising:

2d chiral critical surface separates 1st order from crossover



$$B_4(\beta_c, am, N_\sigma) \approx 1.604 + c\,(am - am_c)\, N_\sigma^{1/0.6301}$$

Methodology to determine order of transition

Finite size scaling of generalised cumulants

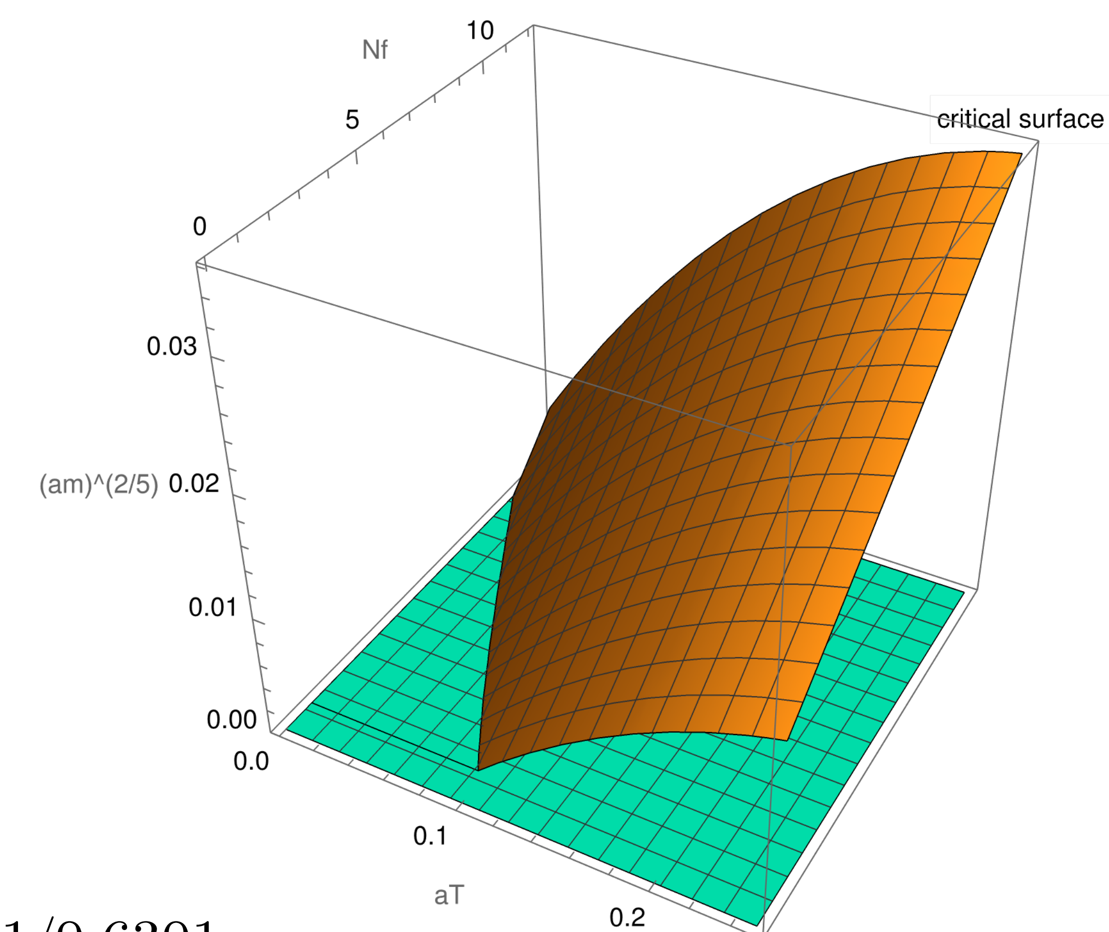
$$B_n = \frac{\langle (\bar\psi\psi - \langle\bar\psi\psi\rangle)^n\rangle}{\langle (\bar\psi\psi - \langle\bar\psi\psi\rangle)^2\rangle^{n/2}}$$

Standard staggered fermions, bare parameters: β, am, N_f, N_τ

(Pseudo-critical) phase boundary: $B_3 = 0$ 3d manifold

Second-order 3d Ising:

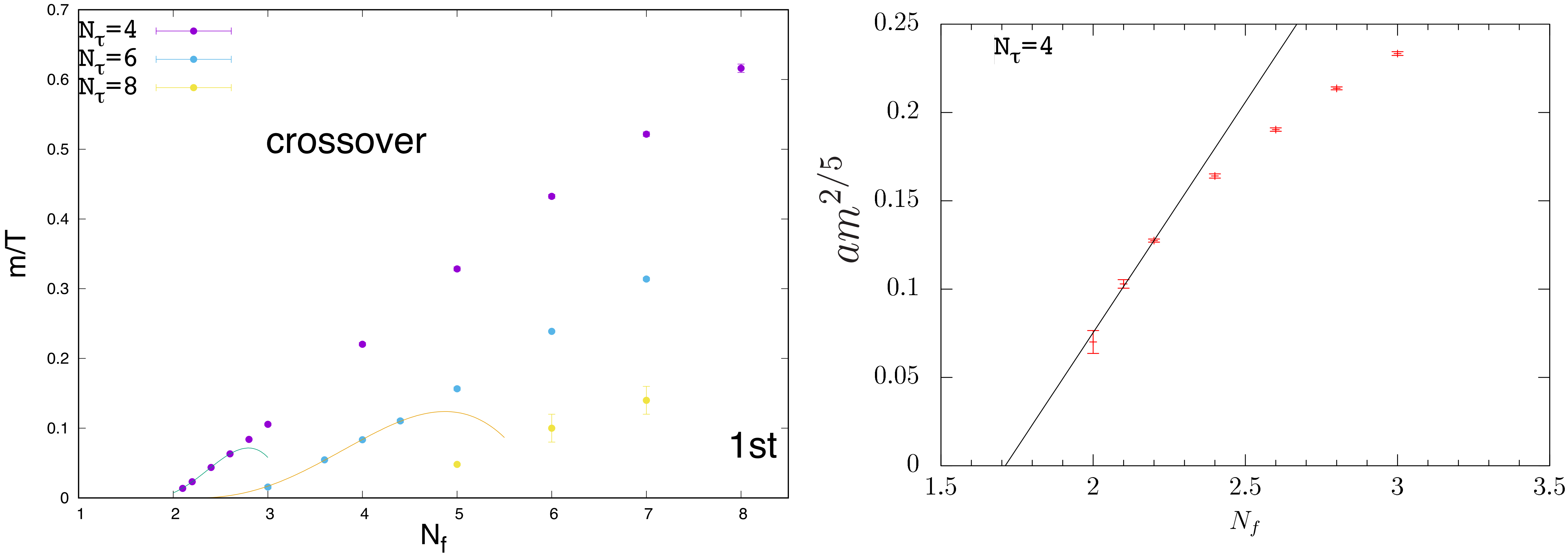
2d chiral critical surface separates 1st order from crossover



$$B_4(\beta_c, am, N_\sigma) \approx 1.604 + c\,(am - am_c)\,N_\sigma^{1/0.6301}$$

Bare parameter space of unimproved staggered LQCD

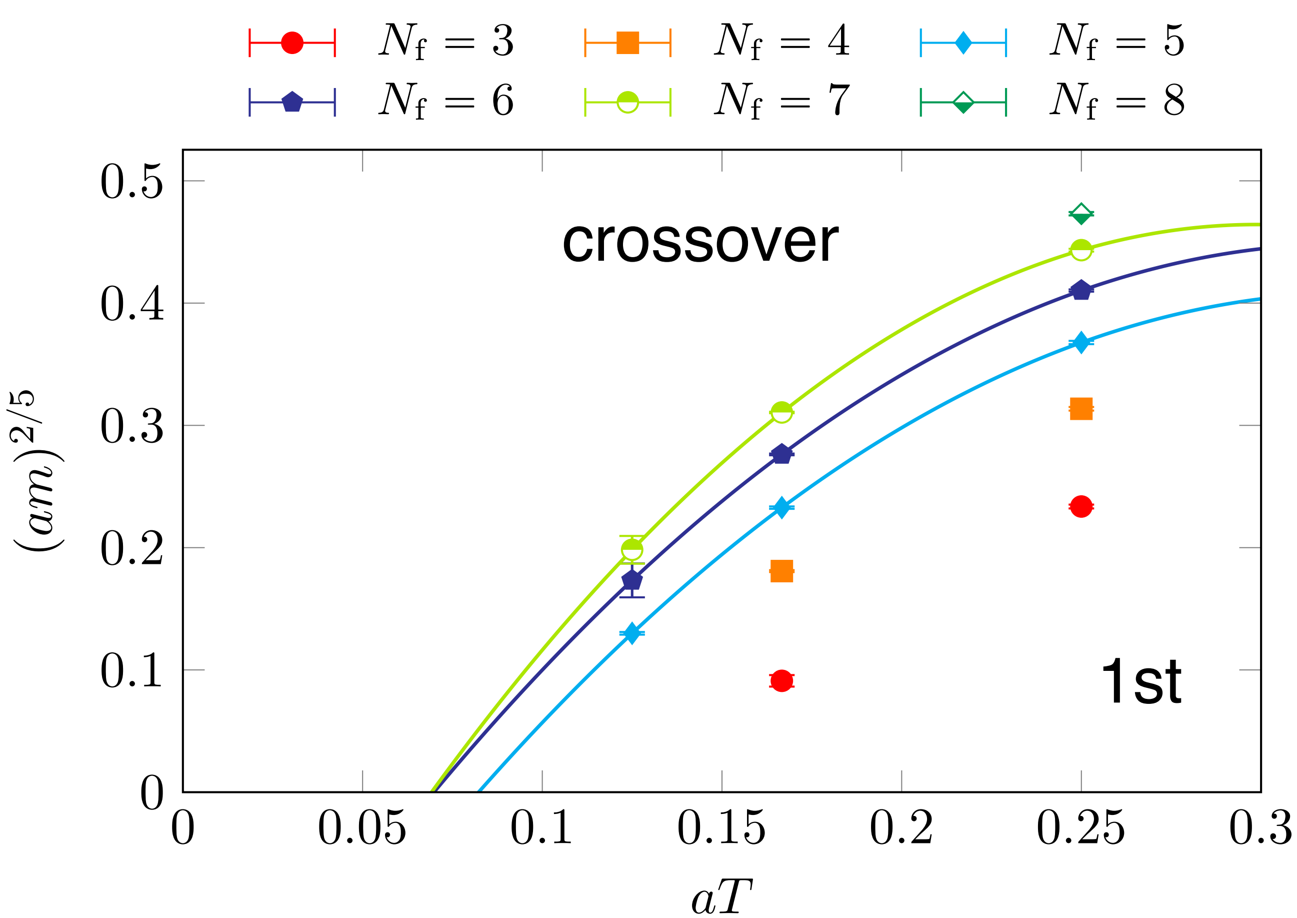
[Cuteri, O.P., Sciarra 21]   ~120 M Monte Carlo trajectories with light fermions, aspect ratios 3,4,5



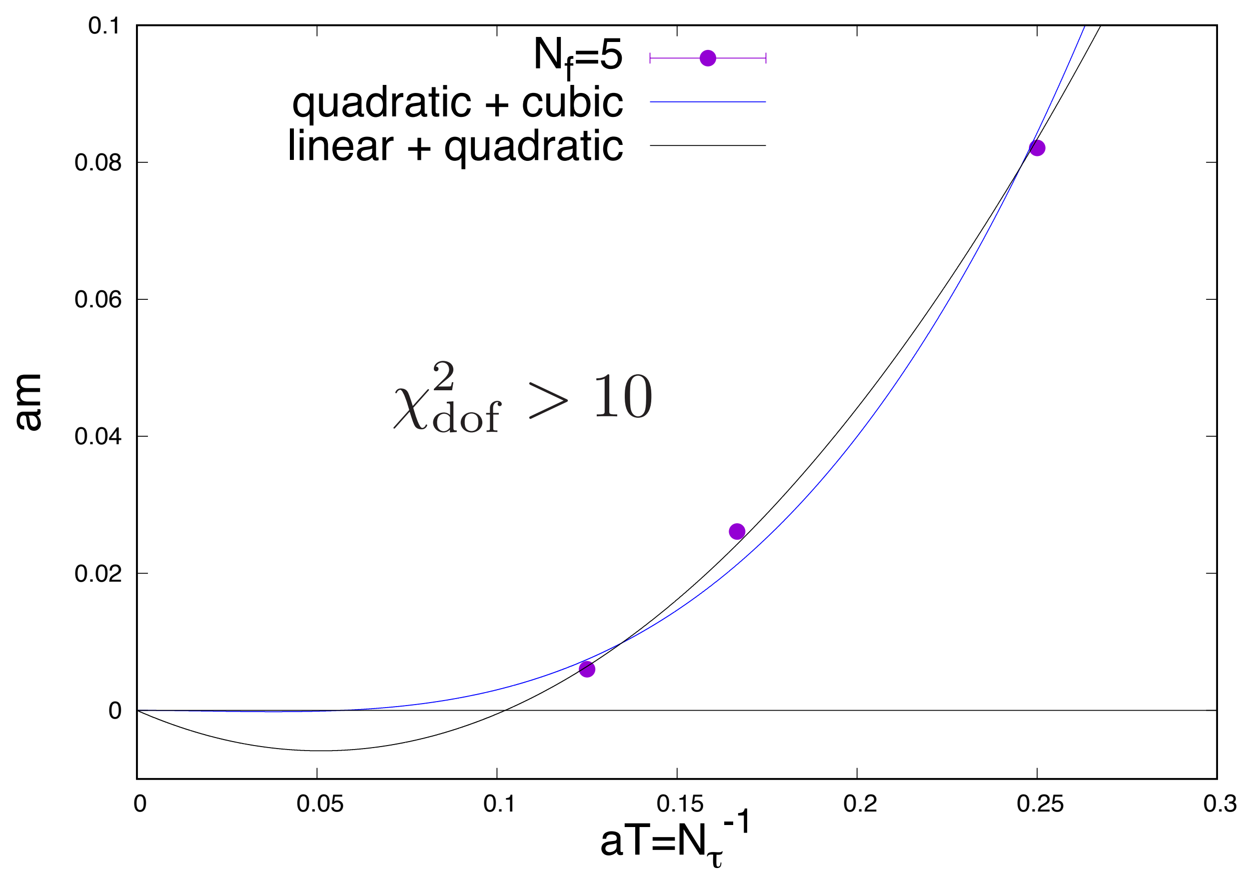
- Tricritical scaling observed in different variable pairings
- Old question: $m_c/T = 0$ or $\neq 0$? Answered for $N_f = 2$
- New question: will $N_f^{\rm tric}$ slide beyond $N_f = 3$?

Bare parameter space of unimproved staggered LQCD

[Cuteri, O.P., Sciarra 21]



1st order scenario does not fit!

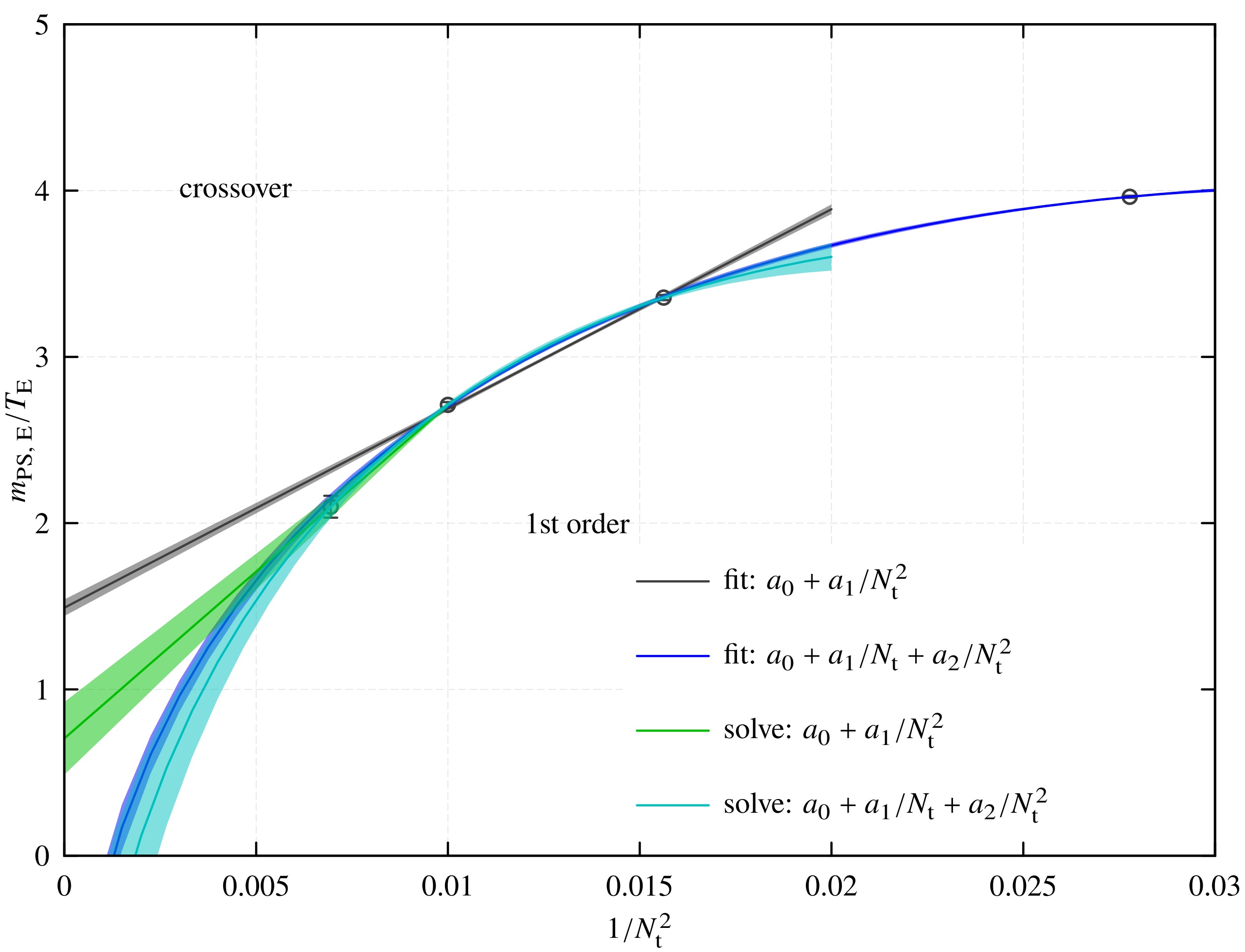


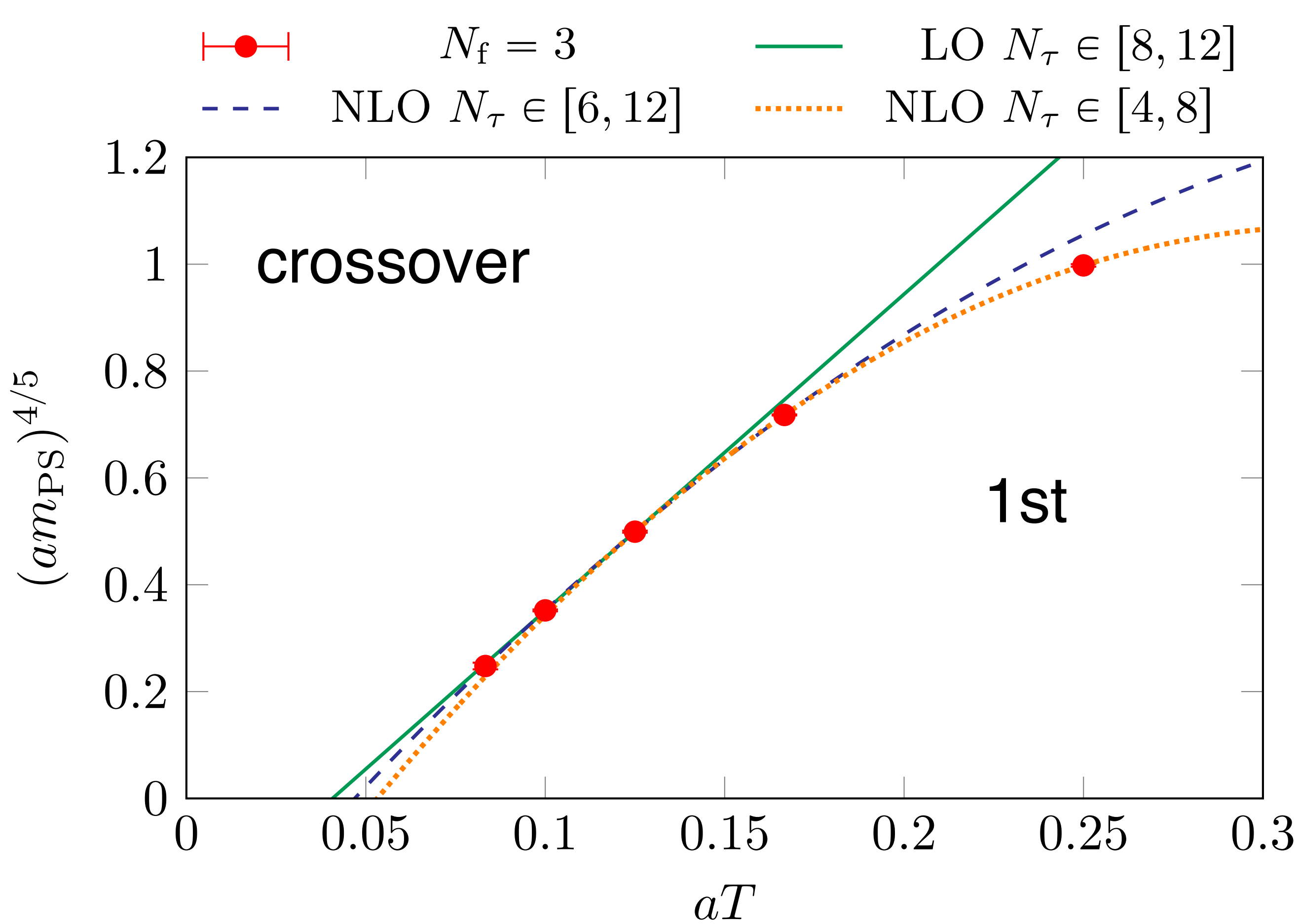
- Tricritical scaling observed also in plane of mass vs. lattice spacing
- Allows extrapolation to lattice chiral limit, tricritical points $N_\tau^{\rm tric}(N_f)$
- 1st order scenario: $m_c(a) = m_c(0) + c_1(aT) + c_2(aT)^2 + \ldots$ incompatible!
- The chiral transition is second order for $N_f = 2 - 6$

Nf=3 O(a)-improved Wilson fermions

[Kuramashi et al. PRD 20] $m_\pi^c \leq 110\ \mathrm{MeV} \quad N_\tau = 4, 6, 8, 10, 12$

Re-analysis using: $am_{PS}^2 \propto am_q$





[Cuteri, O.P., Sciarra, JHEP 21]

Tricritical scaling!

Nf=3 consistent with staggered, 2nd order in chiral continuum limit!

The Columbia plot in the continuum

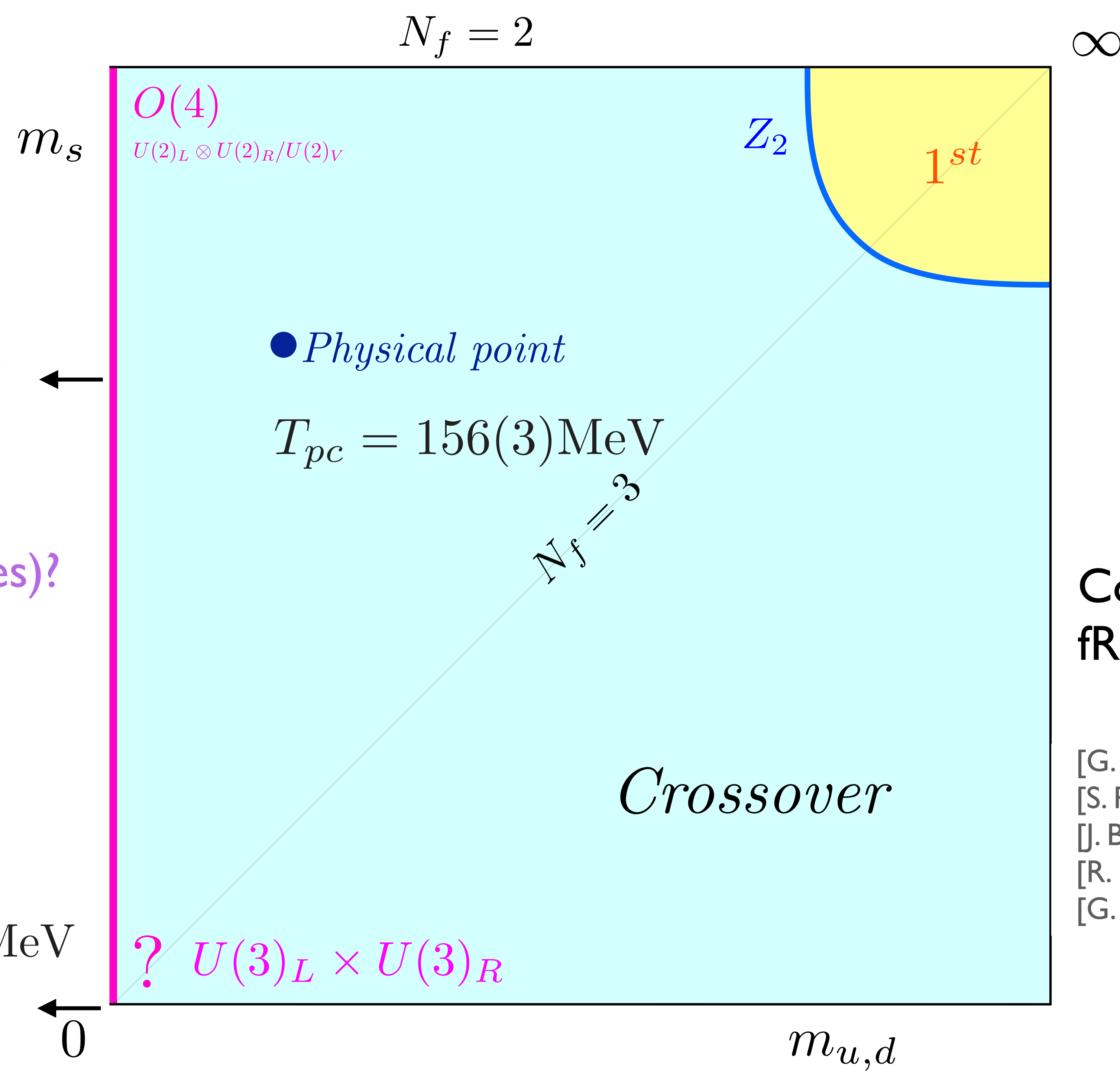
[Cuteri, O.P., Sciarra JHEP 21]

[HotQCD PRL 19]
$T_c^0 = 135(8)\text{MeV}$

Universality class(es)?

[HotQCD PRD 22]

HISQ, crossover down to $m_\pi \approx 80$ MeV
$T_c^0 = 98(6)\text{MeV}$



Consistent with:
fRG, DSE, bootstrap…

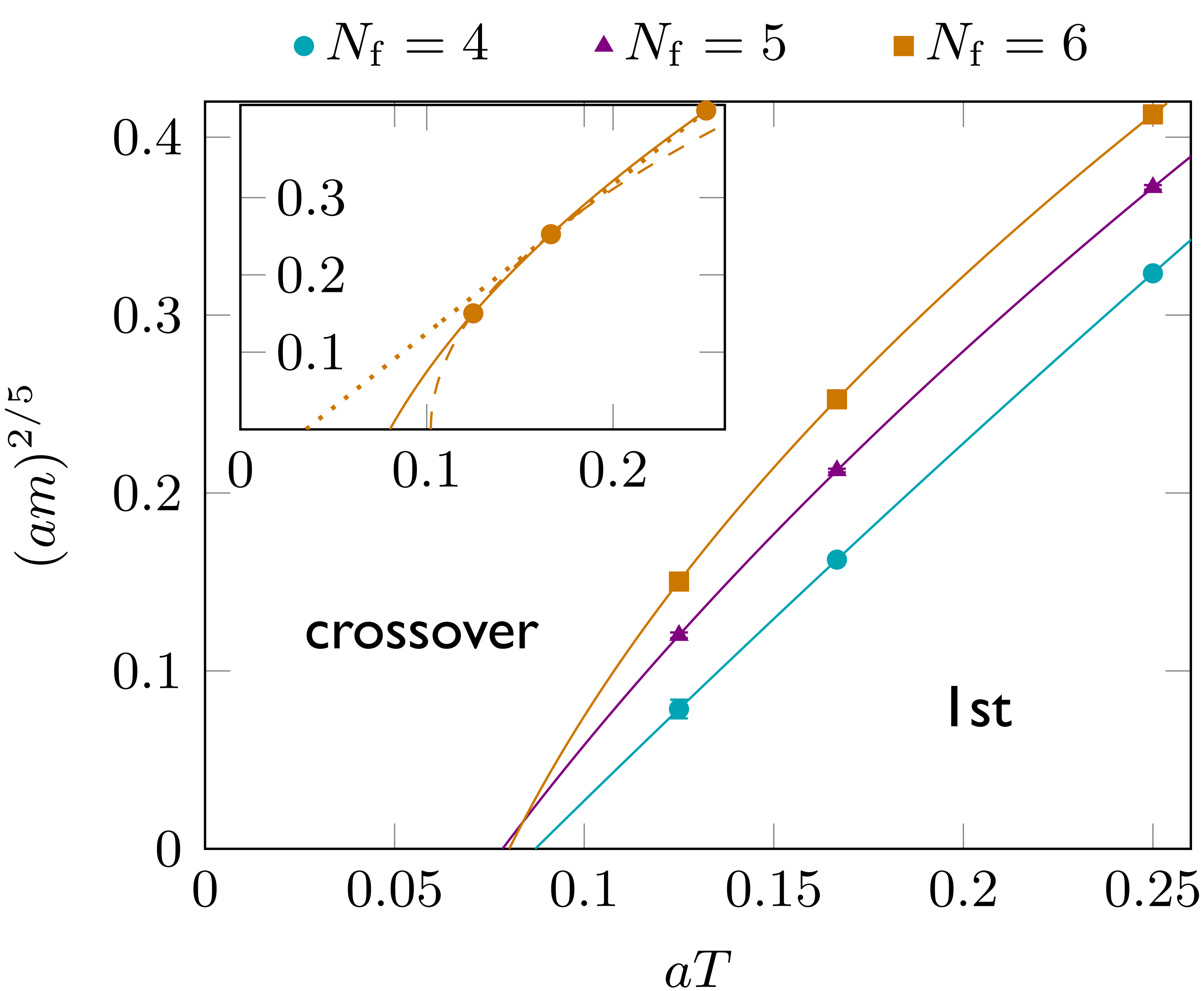
[G. Fejos, PRD 22]
[S. R. Kousvos, A. Stergiou SciPost 23]
[J. Bernhardt, C. S. Fischer PRD 23]
[R. D. Pisarski, F. Rennecke PRL 24]
[G. Fejos, T. Hatsuda PRD 24]

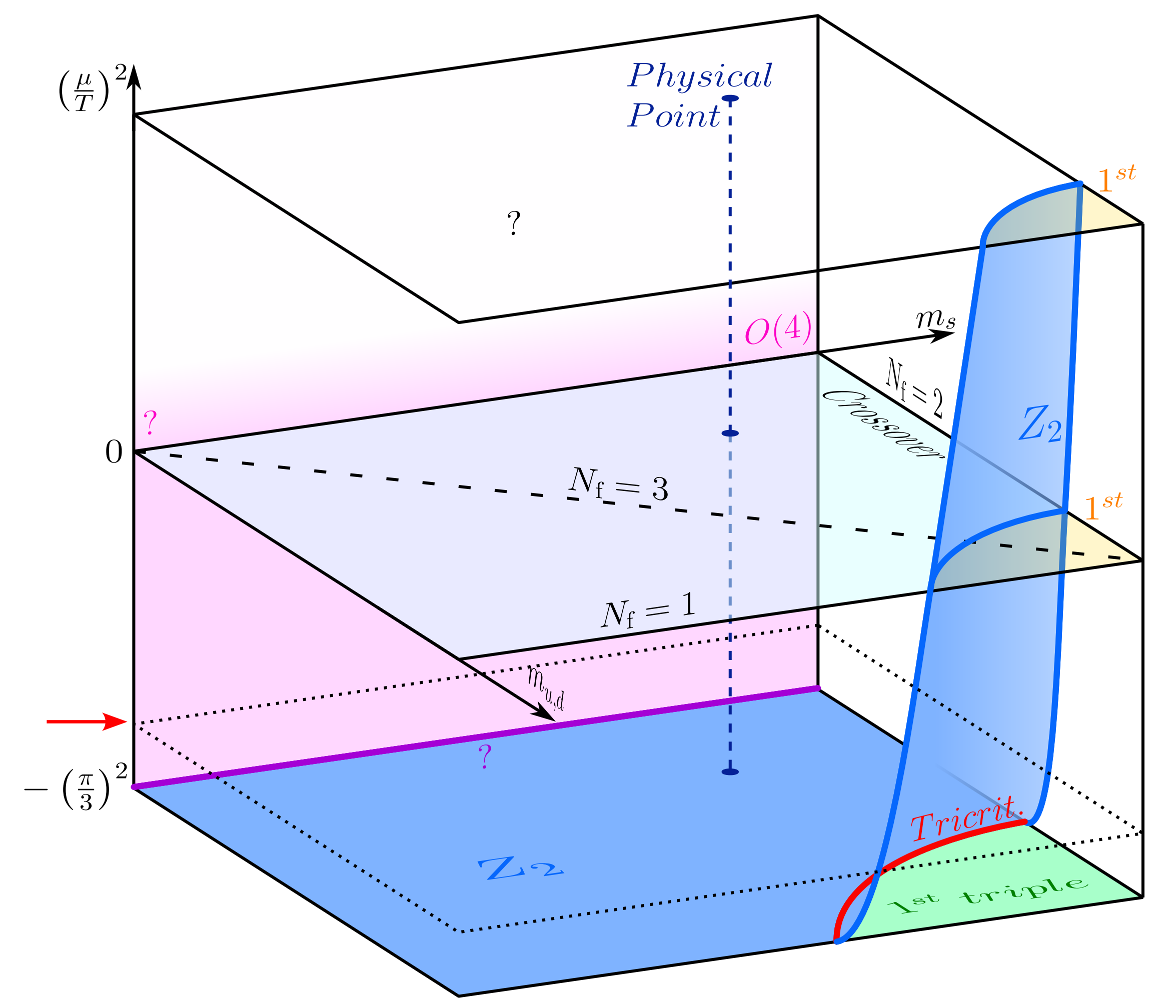
Crossover for DW fermions, Nf=3, $m_q \sim m_{phys}$ [Zhang et al., PoS LAT22, 23]

Imaginary chemical potential

[D'Ambrosio, Kaiser, O.P., PoS LAT 22 + in preparation]

Repeat study of Columbia plot with $\mu = i\,0.81\pi T/3$ Same situation as $\mu = 0$





1st-order region not connected to continuum limit!

Imaginary chemical potential, improved actions

$\mu = i\pi T/3$ Roberge-Weiss boundary

- [Bonati et al., PRD 19] $N_\tau = 4$
  stout-smeared staggered

  quark mass scan down to $m_\pi \approx 50$ MeV
  fixed m_{ud}/m_s

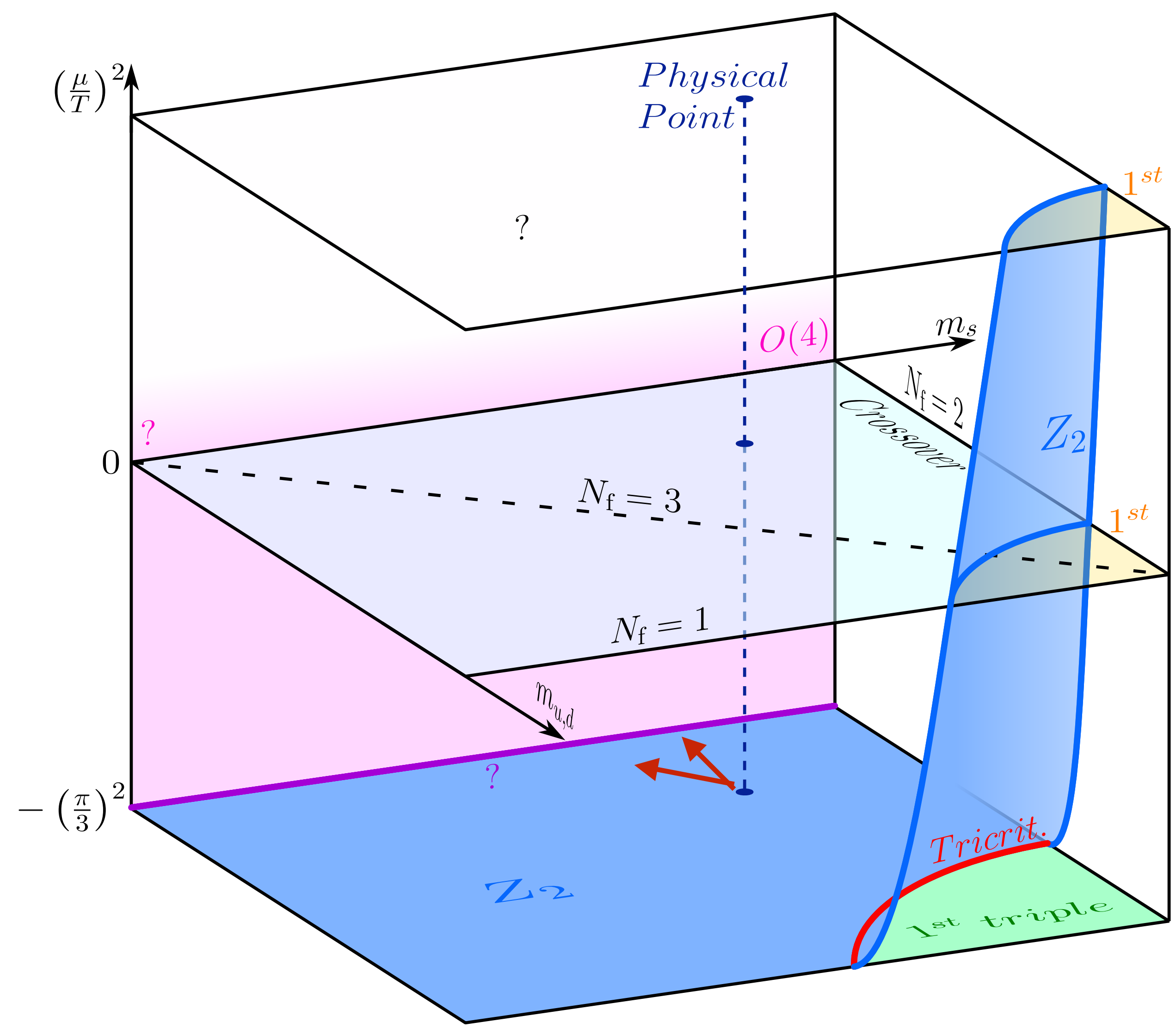
- [Bielefeld+Frankfurt, PRD 22]
  HISQ $N_\tau = 4$

  quark mass scan down to $m_\pi \approx 55$ MeV
  fixed m_s

- No sign of 1st-order phase transition!

- Consistent with DSE approach [Bernhardt, Fischer, PRD 23]

- Entire chiral critical surface moves to massless limit

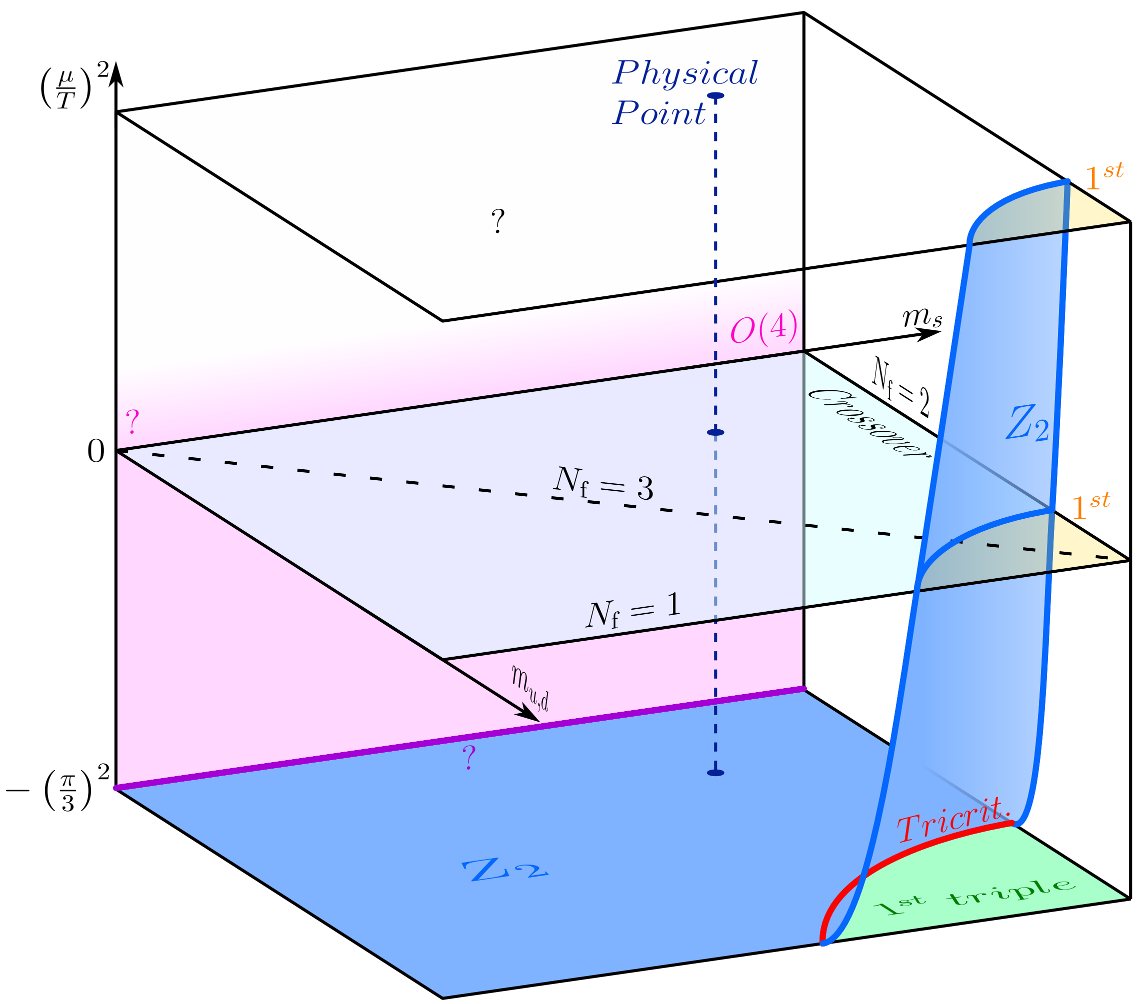


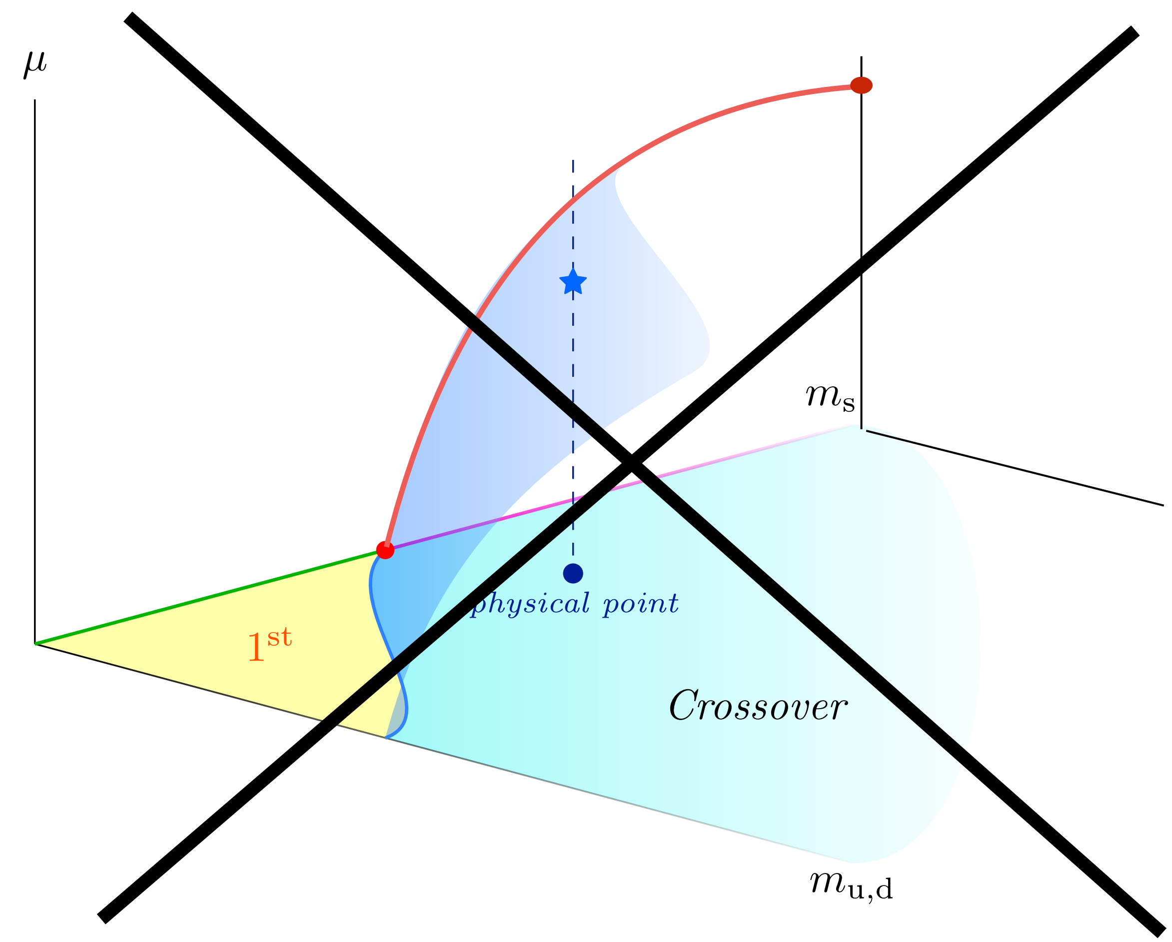
Columbia plot with chemical potential

If we take these results seriously:

Critical point not ruled out,
requires additional critical surface

Class of low energy models now ruled out!

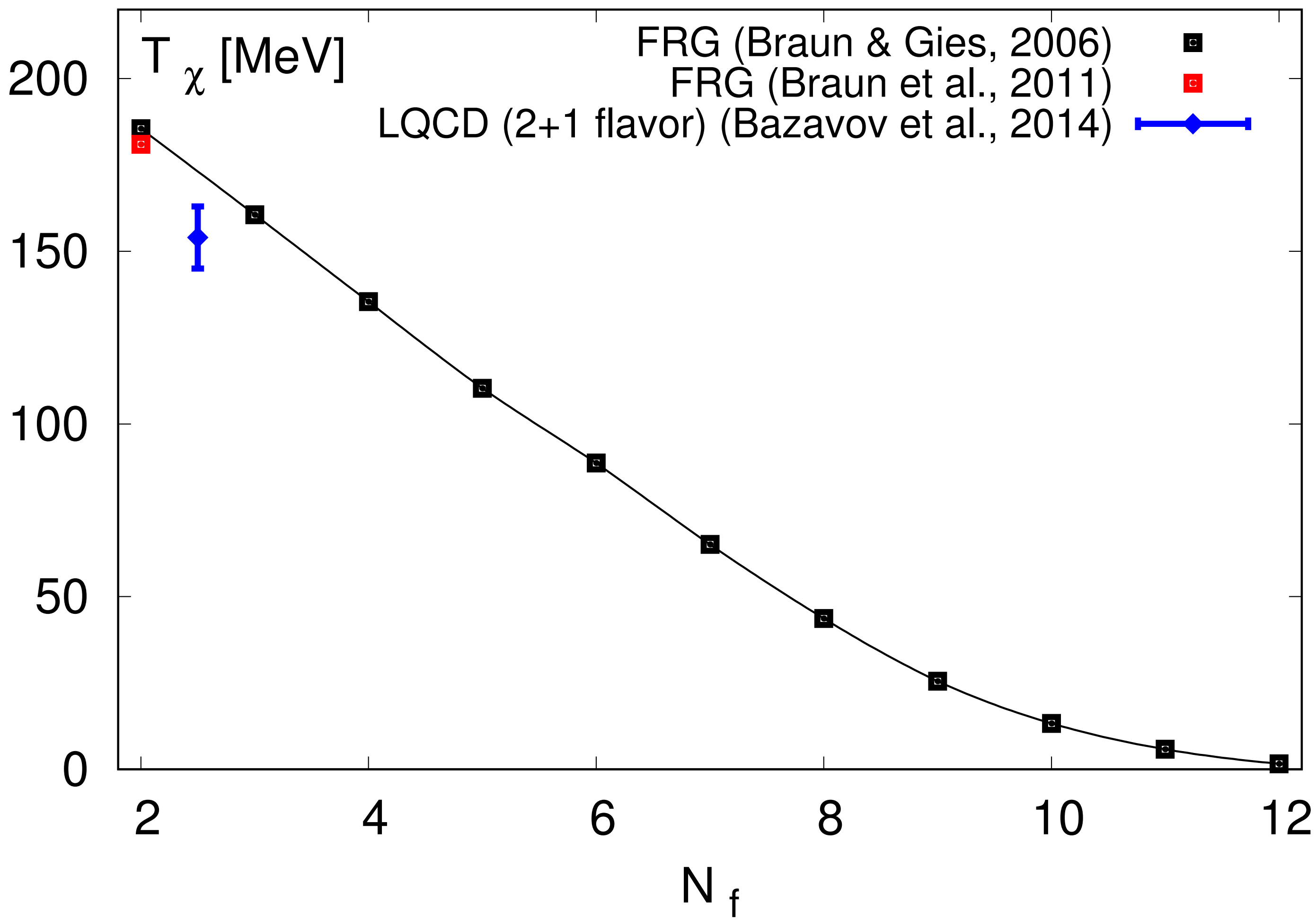




$\nexists$ Tuning of parameters for $N_f = 2 + 1$ theory with critical point at $\mu = 0$!

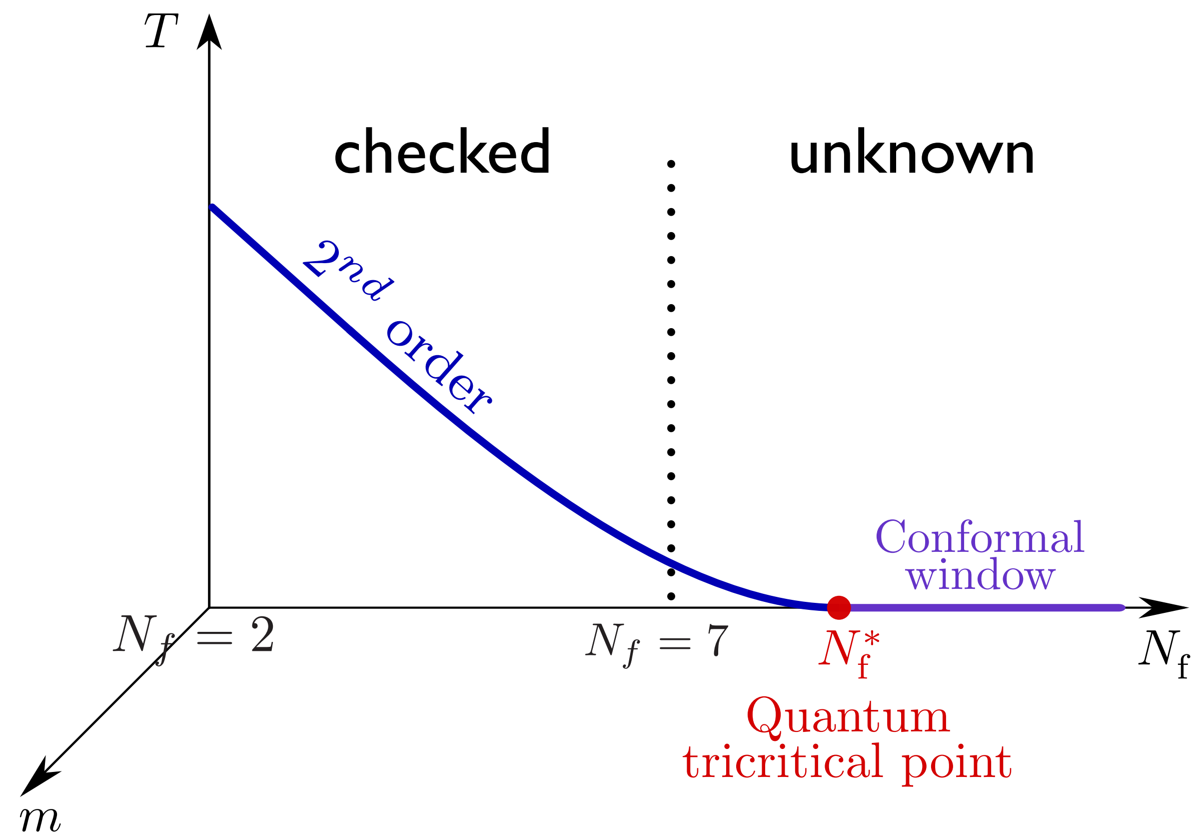
The chiral phase transition for different N_f

Temperature dependence:



For lattice, see [Miura, Lombardo, NPB 13]

Order of the transition:



[Cuteri, O.P., Sciarra, JHEP 21]

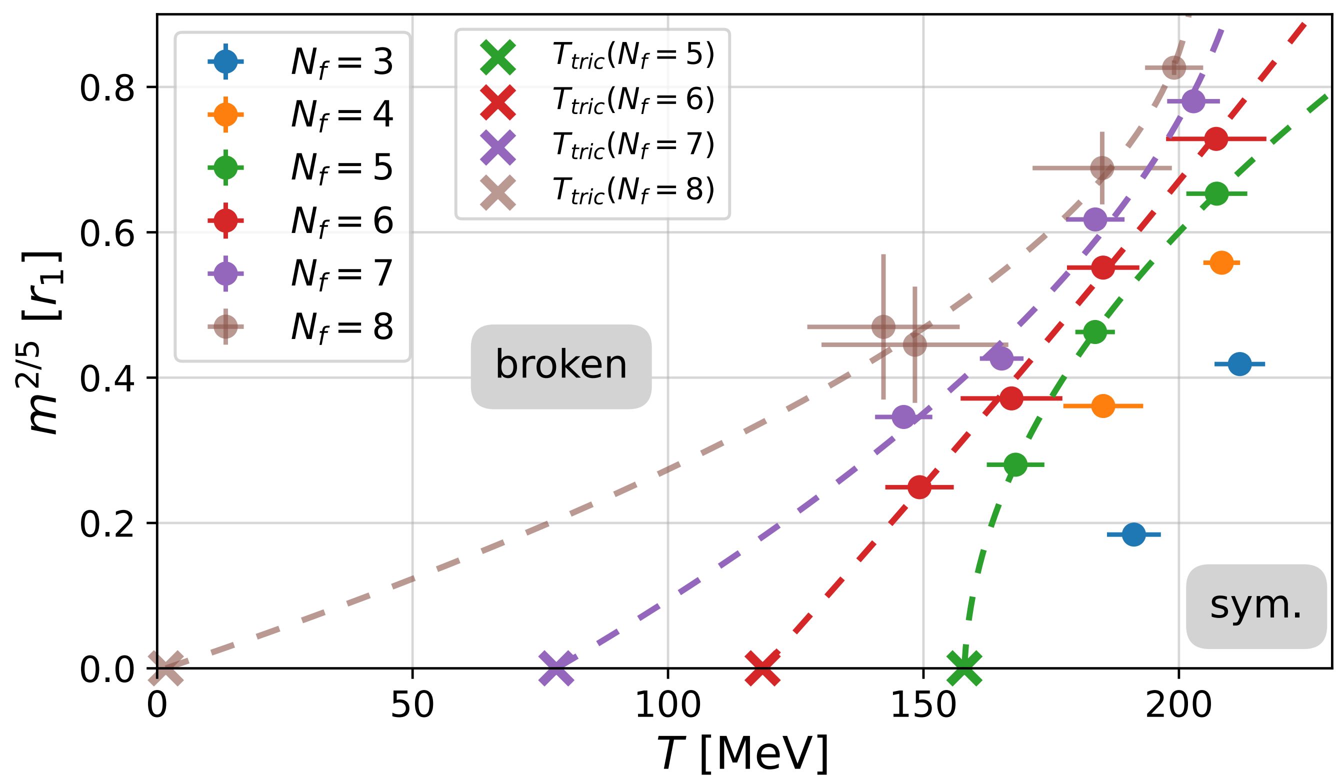
The chiral phase transition in the massless limit is likely second-order for all N_f

Consistent with [Fejos, Hatsuda PRD 24, Pisarski, Renneke PRD 24] with conditions on anomaly

$N_f > 6$, preliminary

- Additional lattice spacing, $N_\tau = 10$
- Scale setting for temperature: Sommer scales r_0, r_1
- Quantitative values of T not important, but when is T=0?
- N_f^* is boundary for tricritical scaling (conformal scaling beyond!)

$$T_c(m) = T_{tric} + A \cdot m^{2/5} + B \cdot m^{4/5}$$



Preliminary result: $7 < N_f^* < 9$

Emergent chiral spin symmetry of QCD

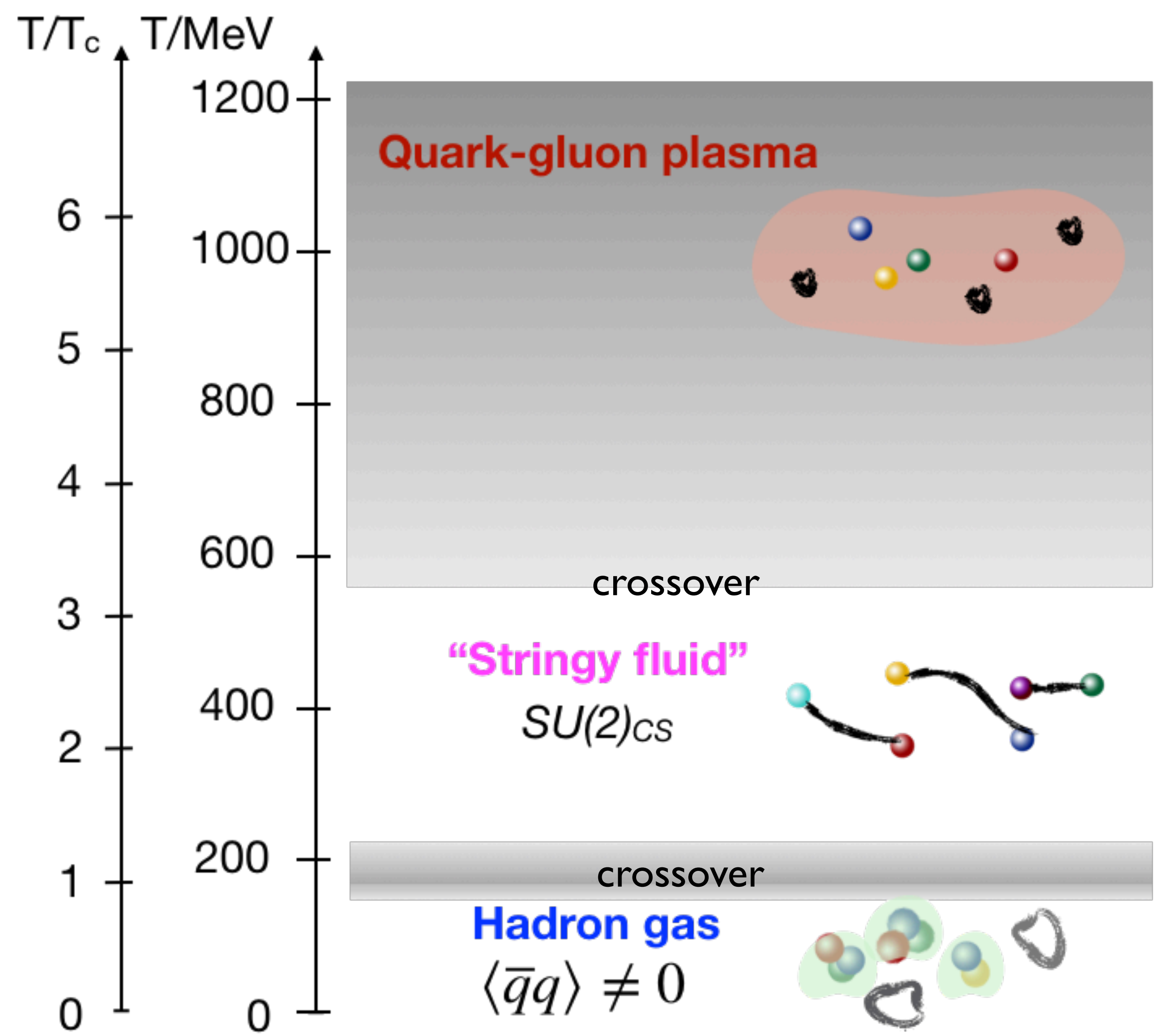
Symmetries (verified):

Degrees of freedom (to be verified):

Chiral symmetry (approximate)

Chiral spin symmetry (approximate)

Chiral symmetry broken



Rohrhofer et al., Phys. Rev. D100 (2019)

Check well-studied observables: screening masses

$$C^s_\Gamma(z) = \sum_{x,y,\tau} C_\Gamma(\tau, \boldsymbol{x}) \quad \stackrel{z\to\infty}{\longrightarrow} \text{const. } e^{-m_{scr} z}$$

Directly related to the partition function and equation of state

by transfer matrices: $T = e^{-aH}, T_z = e^{-aH_z}$

$$e^{pV/T} = Z = \mathrm{Tr}(e^{-aHN_\tau})$$
$$= \mathrm{Tr}(e^{-aH_z N_z}) = \sum_{n_z} e^{-E_{n_z} N_z}$$

Screening masses: eigenvalues of H_z

For T=0 equivalent to eigenvalues of H, for $T \neq 0$ temperature effects

Meson screening masses at high temperatures

[Dalla Brida et al., JHEP 22]

Nf=3, T=1 GeV -160 GeV

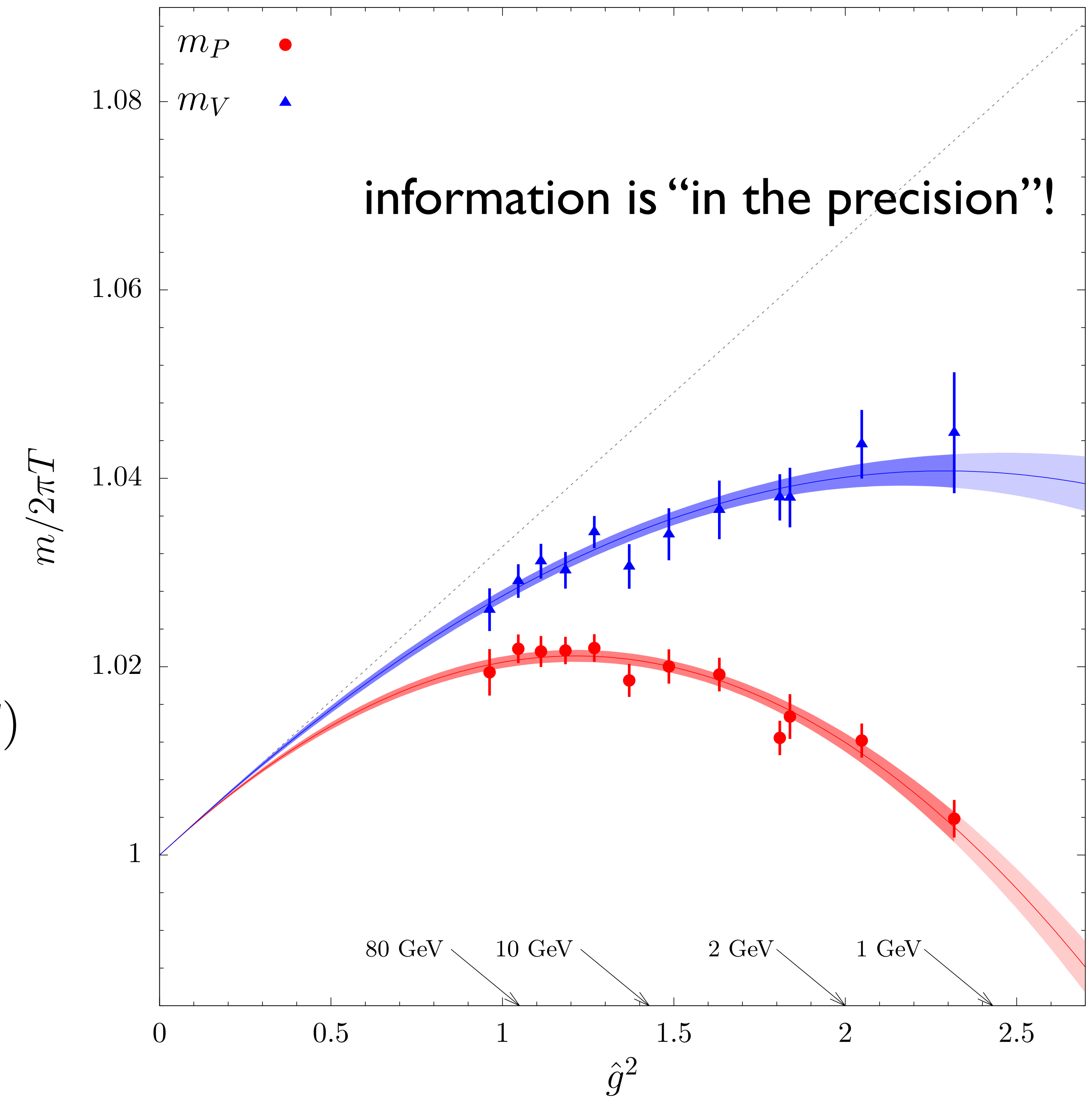
Highly non-trivial technically:
shifted b.c. + step-scaling techniques

$$\frac{m_{PS}}{2\pi T} = 1 + p_2\,\hat{g}^2(T) + p_3\,\hat{g}^3(T) + p_4\,\hat{g}^4(T)$$
$$\frac{m_V}{2\pi T} = \frac{m_{PS}}{2\pi T} + s_4\,\hat{g}^4(T)$$

$p_2 = 0.032739961$
[Laine, Vepsäläinen., JHEP 04]

p_3, p_4, s_4 fitted, excellent $\chi^2_{\rm dof}$

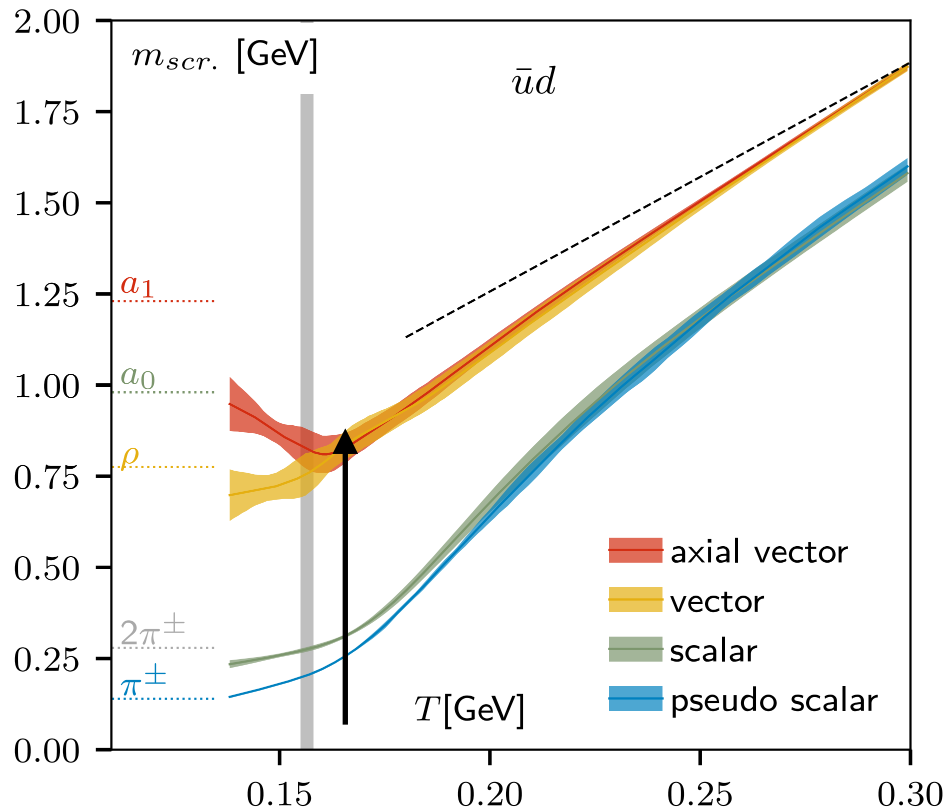
Quark hadron duality holds

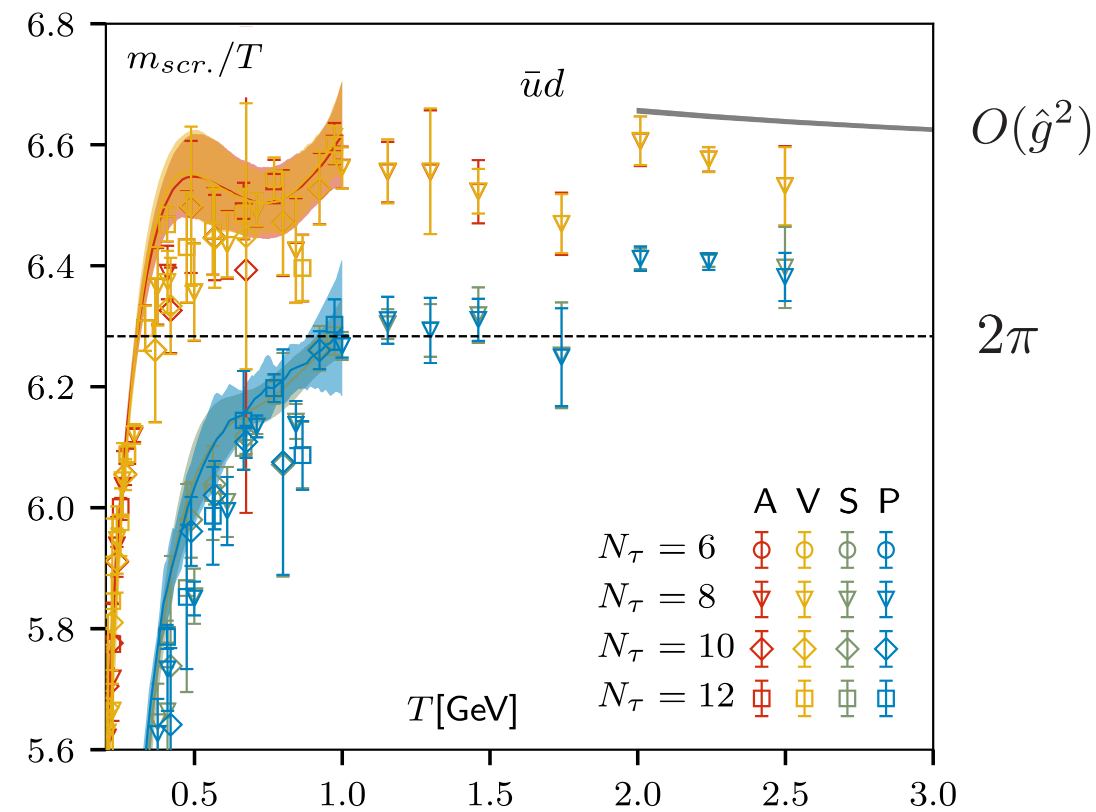


$$\frac{1}{\hat{g}^2(T)} \equiv \frac{9}{8\pi^2}\ln\frac{2\pi T}{\Lambda_{\overline{\rm MS}}} + \frac{4}{9\pi^2}\ln\left(2\ln\frac{2\pi T}{\Lambda_{\overline{\rm MS}}}\right)$$

Meson screening masses at intermediate temperatures

[HotQCD, PRD 19]





Chiral symmetry restoration

Heavy chiral partners "come down" in all flavour combinations

→ pressure increases

Resummed pert. theory:

$$\frac{m_{PS}}{2\pi T} = 1 + p_2\,\hat{g}^2(T) + p_3\,\hat{g}^3(T) + p_4\,\hat{g}^4(T)\ ,$$

$$\frac{m_V}{2\pi T} = \frac{m_{PS}}{2\pi T} + s_4\,\hat{g}^4(T)\ ,$$

Cannot describe the "bend"

No quark hadron duality for T<0.5 GeV in 12 lightest meson channels! CS symmetry!

[Glozman, O.P., Pisarski, EPJA 22]

Spectral functions at finite T

General euclidean correlator:

$$C(\tau, \mathbf{p}) = \int_0^\infty \frac{d\omega}{2\pi} \frac{\cosh(\omega(|\tau| - \beta/2))}{\sinh(\beta\omega/2)} \rho(\omega, \mathbf{p})$$

- Inversion problem ill-defined on a discrete lattice
- Statistical approaches to find "most likely" spectral function
- Alternative: microcausality + KMS [Bros, Buchholz, Ann. Inst. Poincare Phys, Theo 96]

$$\rho(\omega, \mathbf{p}) = \int_0^\infty ds \int \frac{d^3\mathbf{u}}{(2\pi)^2}\, \epsilon(\omega)\, \delta\big(\omega^2 - (\mathbf{p} - \mathbf{u})^2 - s\big)\, \widetilde{D}_\beta(\mathbf{u}, s)$$

← Thermal spectral density

$$\widetilde{D}_\beta(\vec{u}, s) = \widetilde{D}_{m,\beta}(\vec{u})\, \delta(s - m^2) + \widetilde{D}_{c,\beta}(\vec{u}, s)$$

← negligible at low T

"thermoparticle" continuous, scattering, Landau damping etc.

$$C(x_3) = \frac{1}{2} \int_0^\infty ds \int_{|x_3|}^\infty dR\, e^{-R\sqrt{s}} D_\beta(R, s)$$

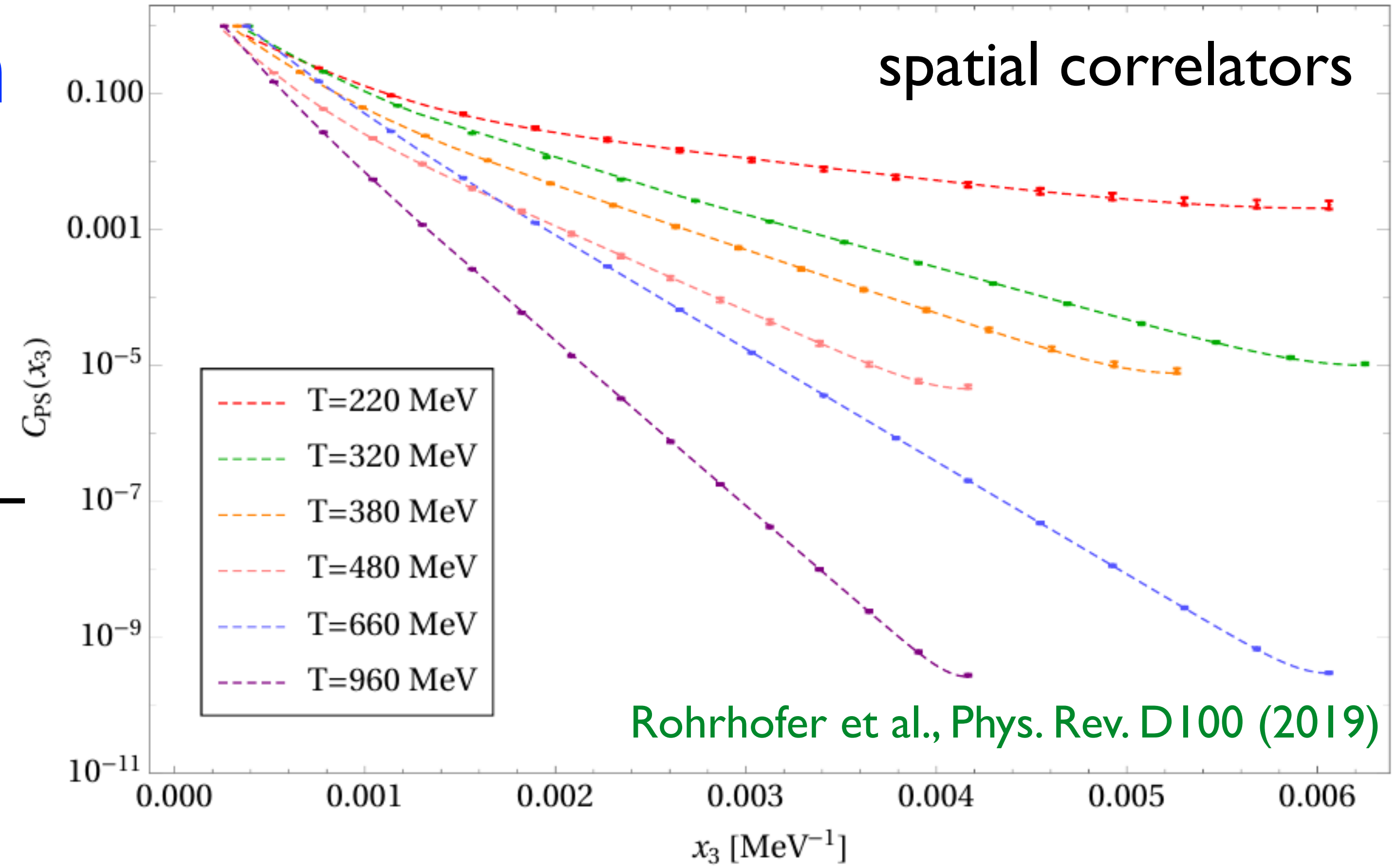
[Lowdon, O.P. JHEP 22]

The pion spectral function

[Lowdon, O.P., JHEP 22]

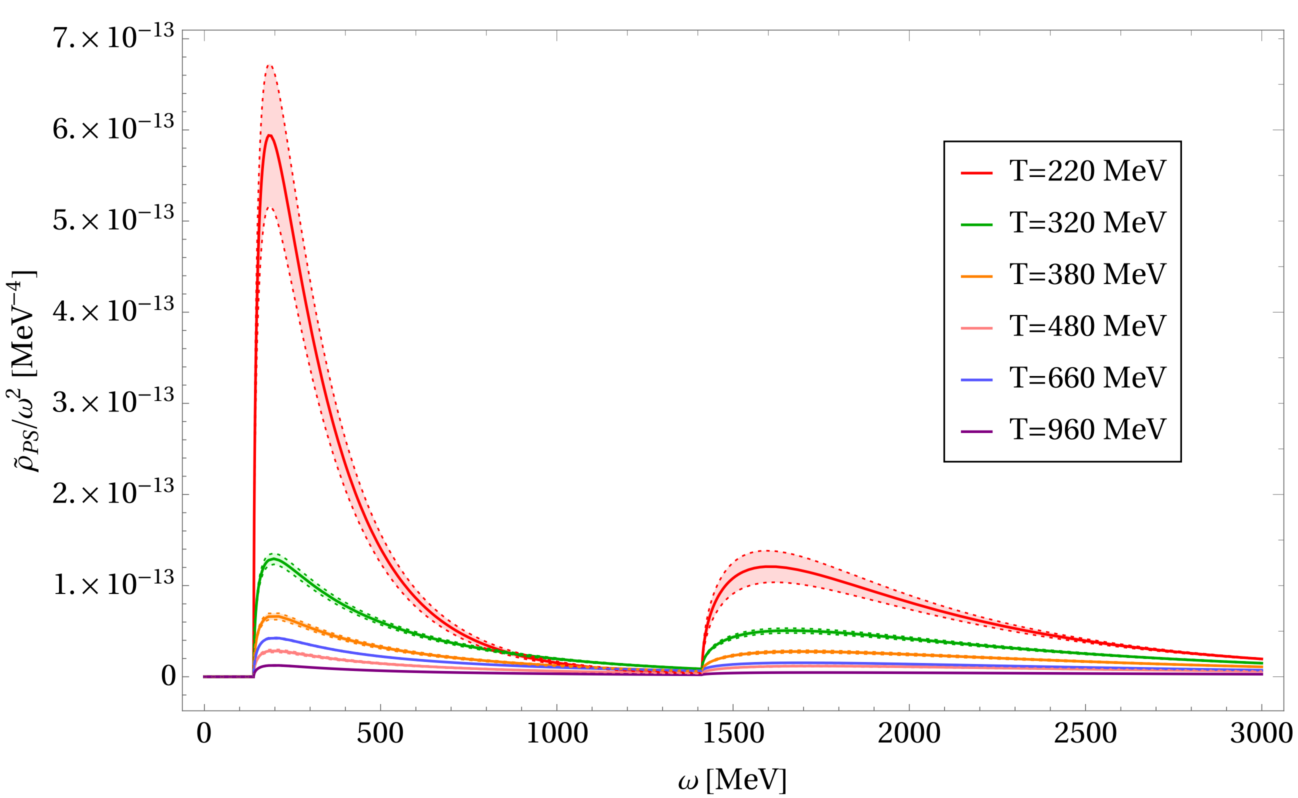
2-state fits π, π^*

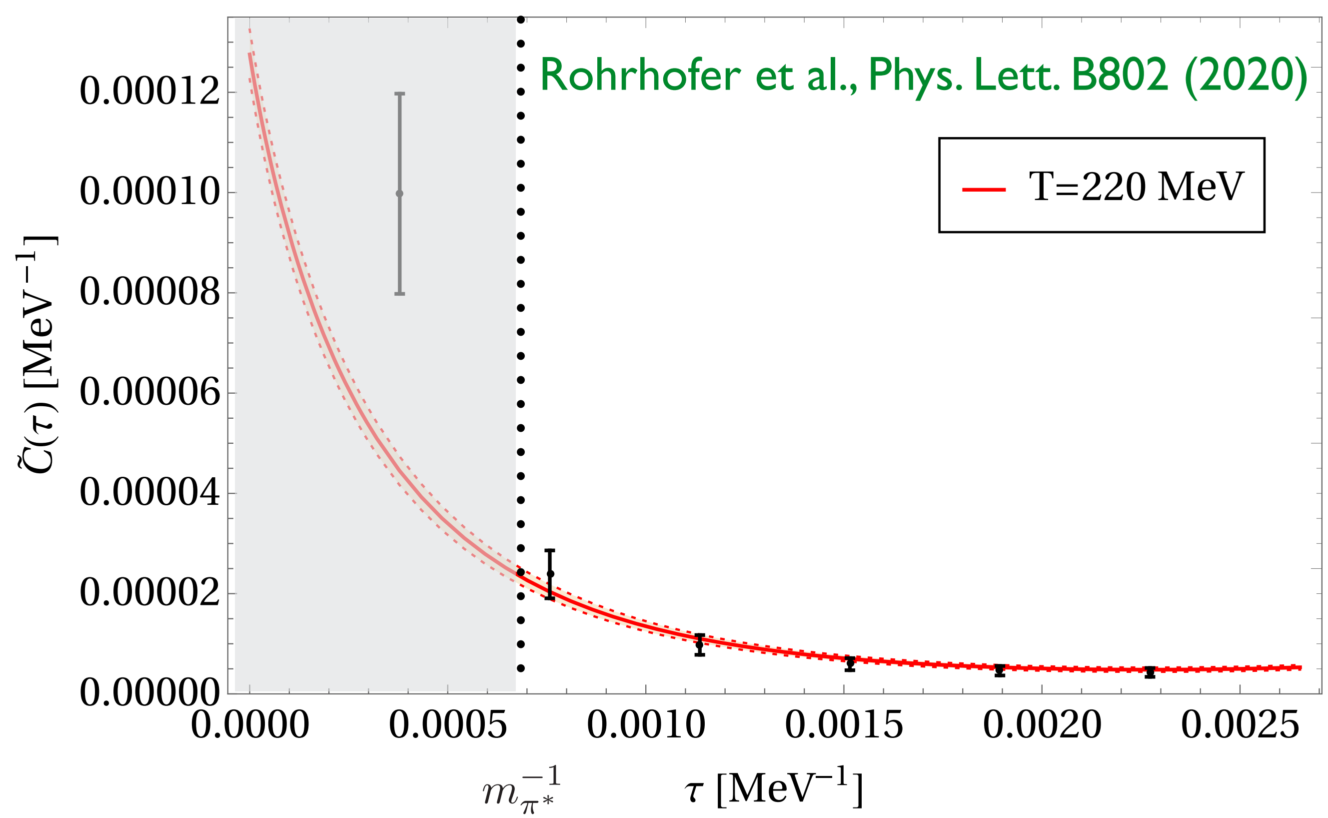
$$D_{m_{\pi^{(*)}},\beta} = \alpha_{\pi^{(*)}} e^{-\gamma_{\pi^{(*)}} x_3}$$



spectral functions

predict temporal correlators, compare with data



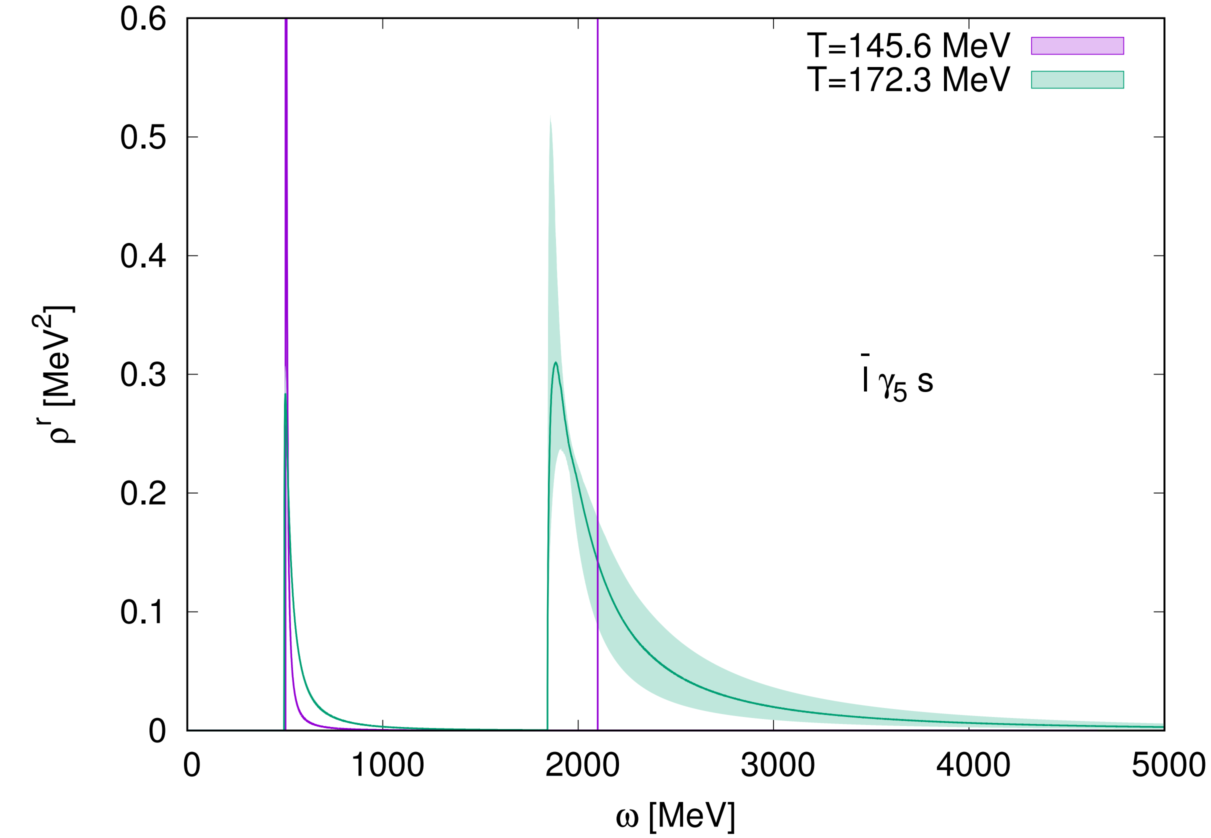


Does QCD deconfine across the chiral crossover ?

[Bala et al., JHEP 24]

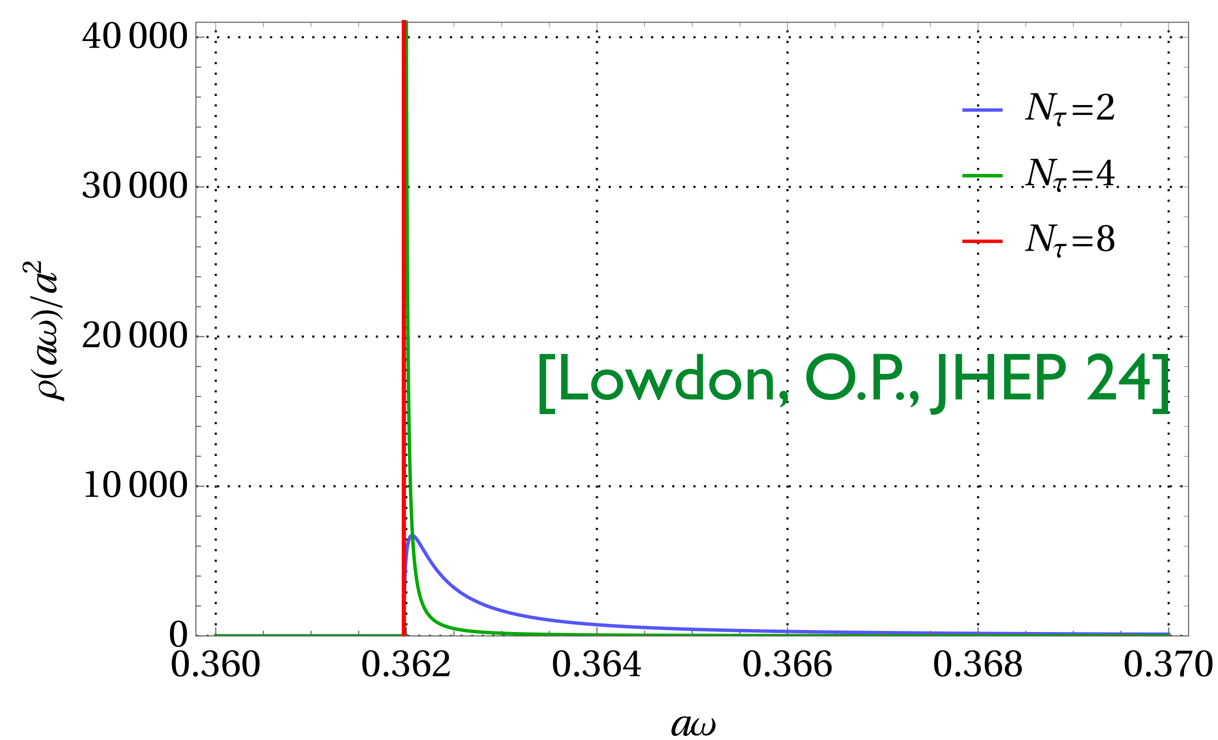
Kaon + Kaon* in full QCD

slightly below and above chiral crossover



Scalar point particle in ϕ^4

no phase transition, no "melting",
only "collisional broadening"



[Lowdon, O.P., JHEP 24]

III. Effective heavy (dense) lattice theory from Wilson action

Pure gauge part: character expansion

$$u(\beta) = \frac{\beta}{18} + \frac{\beta^2}{216} + \ldots < 1$$

$$\beta = \frac{2N_c}{g^2} \qquad T = \frac{1}{aN_\tau}$$

Fermion determinant:

expansion in spatial hops,
temporal hops fully included

$$\kappa = \frac{1}{2am+8}$$

Integrate analytically over spatial links: $W_{\mathbf{x}} = \prod_{\tau=1}^{N_\tau} U_0(\tau, \mathbf{x}), \quad L(\mathbf{x}) = \mathrm{Tr} W(\mathbf{x}), \quad DW = \prod_{\mathbf{x}} dW(\mathbf{x})$

$$Z = \int DW \prod_{<\mathbf{x},\mathbf{y}>} \left[1 + \lambda_1 (L_{\mathbf{x}} L_{\mathbf{y}}^* + L_{\mathbf{x}}^* L_{\mathbf{y}})\right]$$ pure gauge

$$\times \prod_{\mathbf{x}} [1 + h_1 L_{\mathbf{x}} + h_1^2 L_{\mathbf{x}}^* + h_1^3]^{2N_f} [1 + \bar{h}_1 L_{\mathbf{x}}^* + \bar{h}_1^2 L_{\mathbf{x}} + \bar{h}_1^3]^{2N_f}$$ stat. det. $\sim \kappa_s^0$

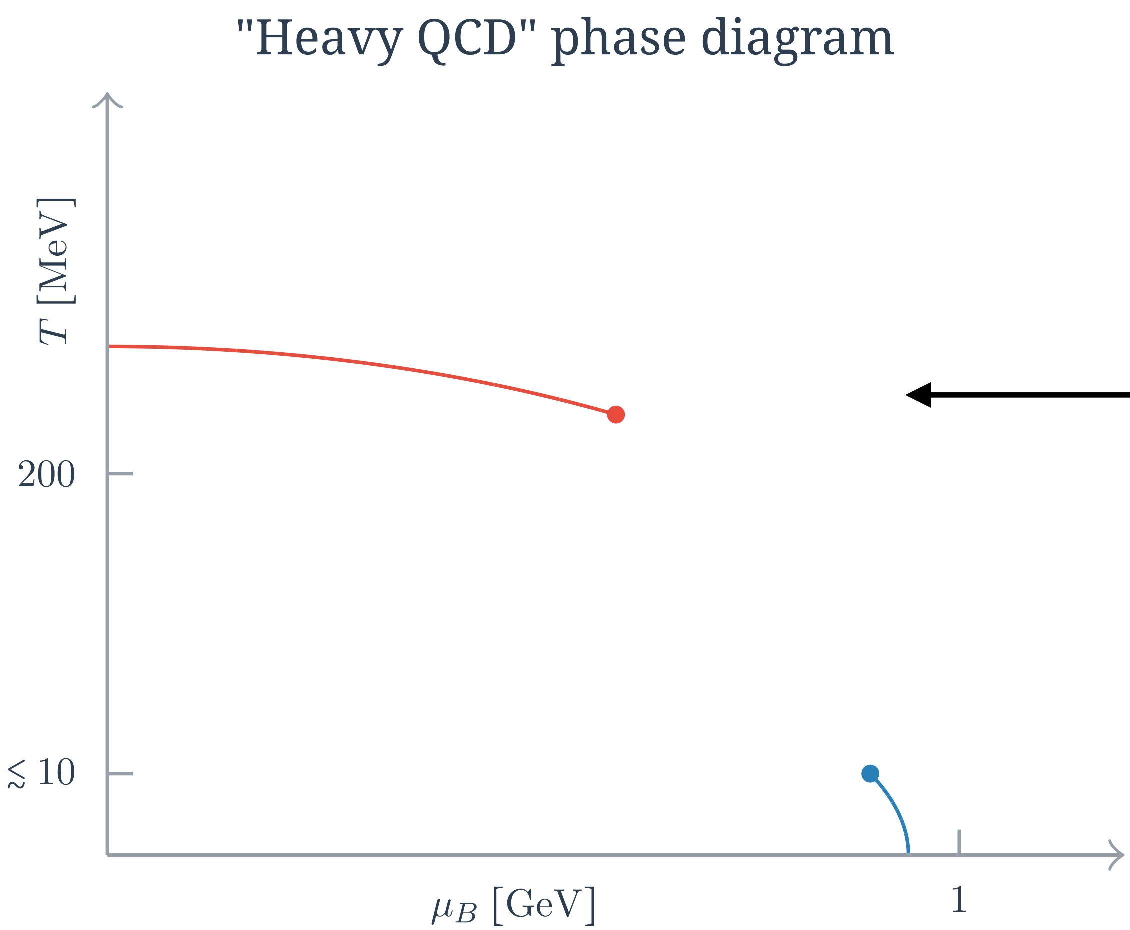
$$\times \prod_{<\mathbf{x},\mathbf{y}>} \left[1 - 2N_f h_2 \left(\mathrm{Tr}\frac{h_1 W_{\mathbf{x}}}{1 + h_1 W_{\mathbf{x}}} - \mathrm{Tr}\frac{\bar{h}_1 W_{\mathbf{x}}^\dagger}{1 + \bar{h}_1 W_{\mathbf{x}}^\dagger}\right) \left(\mathrm{Tr}\frac{h_1 W_{\mathbf{y}}}{1 + h_1 W_{\mathbf{y}}} - \mathrm{Tr}\frac{\bar{h}_1 W_{\mathbf{y}}^\dagger}{1 + \bar{h}_1 W_{\mathbf{y}}^\dagger}\right)\right]$$ kinetic det. $\sim \kappa_s^2$

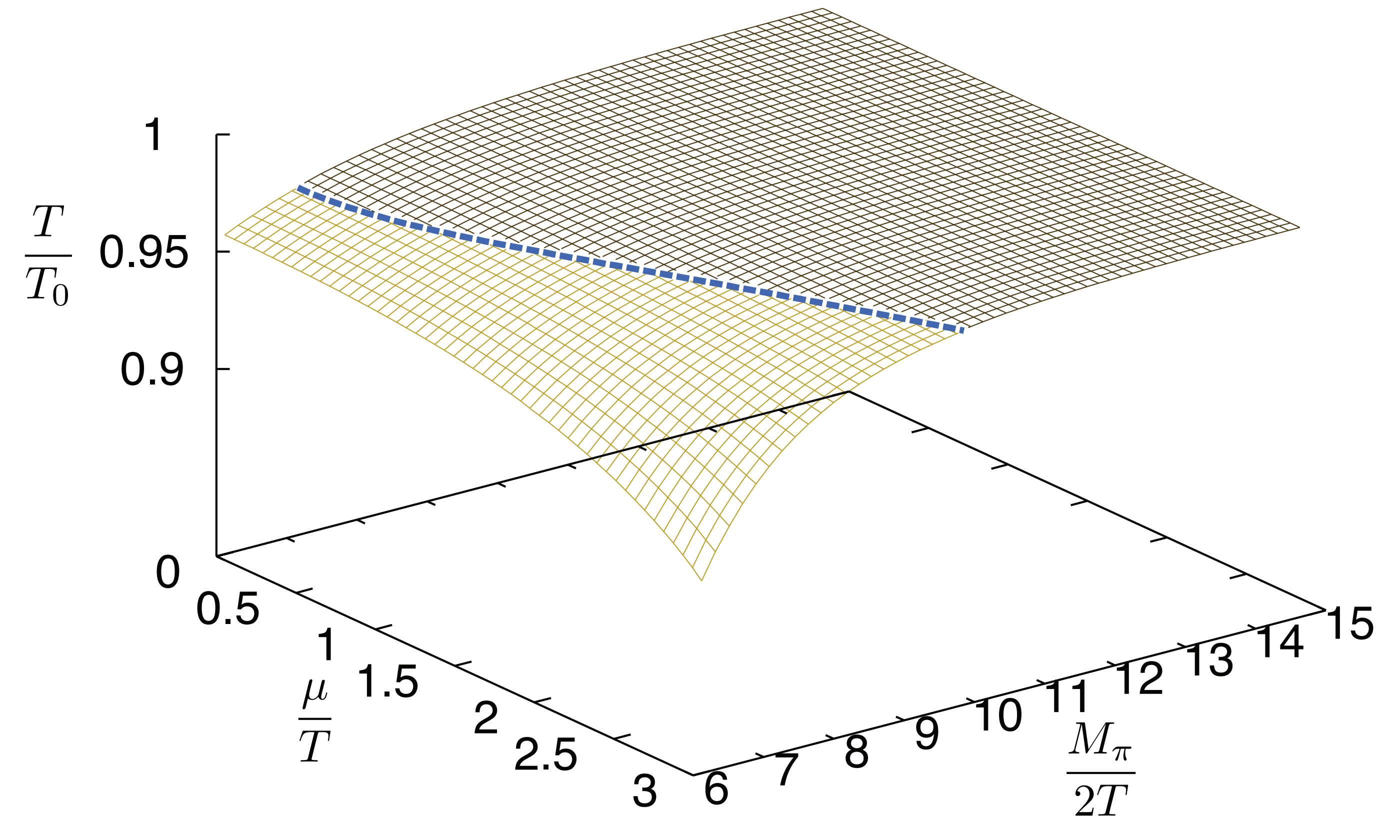
$$\times \ldots$$ $O(\kappa_s^4)$

Example LL* : $\lambda(u, N_\tau \geq 5) = u^{N_\tau} \exp\left[N_\tau \left(4u^4 + 12u^5 - 14u^6 - 36u^7 + \frac{295}{2}u^8 + \frac{1851}{10}u^9 + \frac{1055797}{5120}u^{10}\right)\right]$

The deconfinement transition at finite density

(no continuum limit yet)



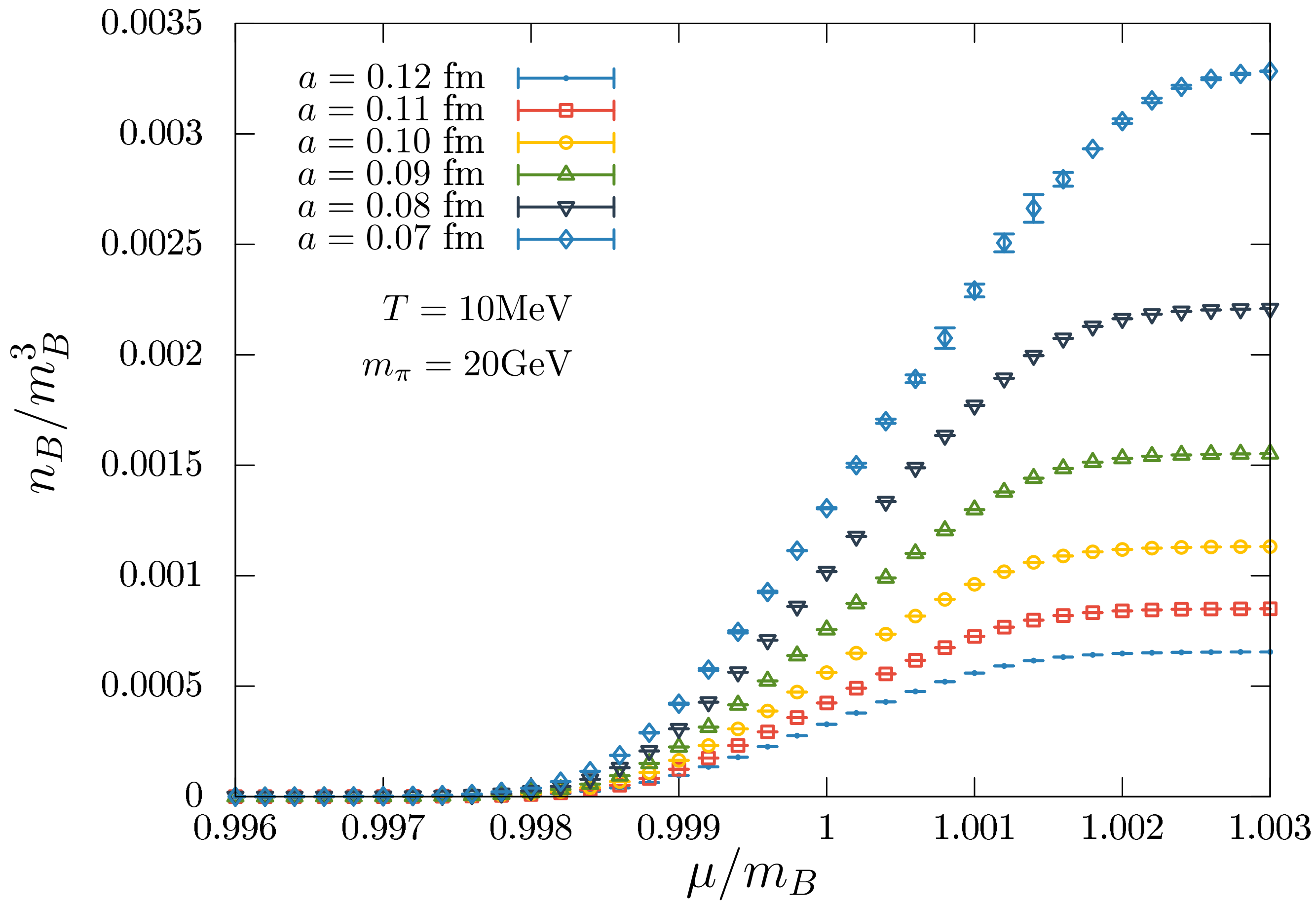


Zero density agrees within 10% with full lattice simulations on Nt=6!

[Fromm, Langelage, Lottini, O.P. JHEP (2012)]

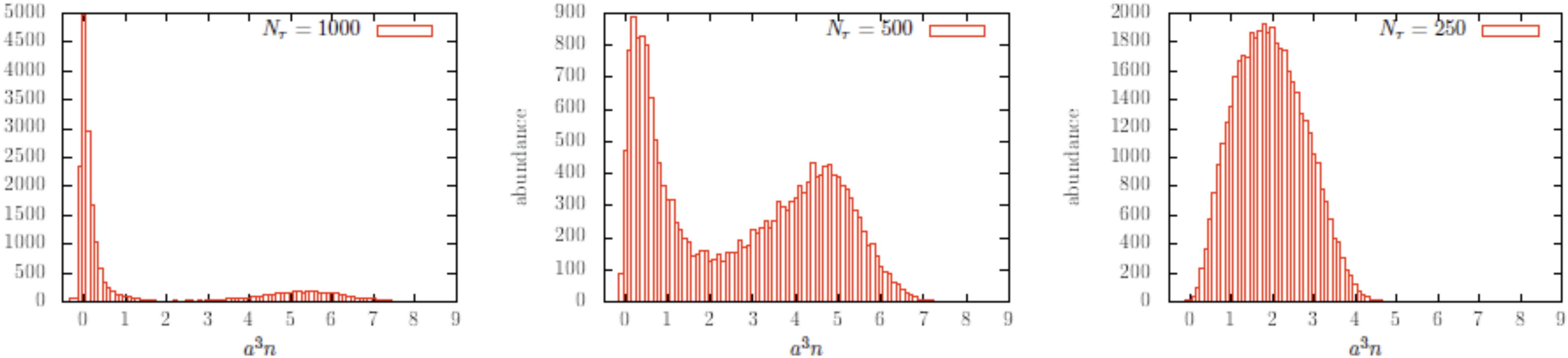
Cold and dense regime

[Fromm, Langelage, Lottini, Neuman, O.P. , PRL 13, Glesaaen, Neuman, O.P. , JHEP 15]

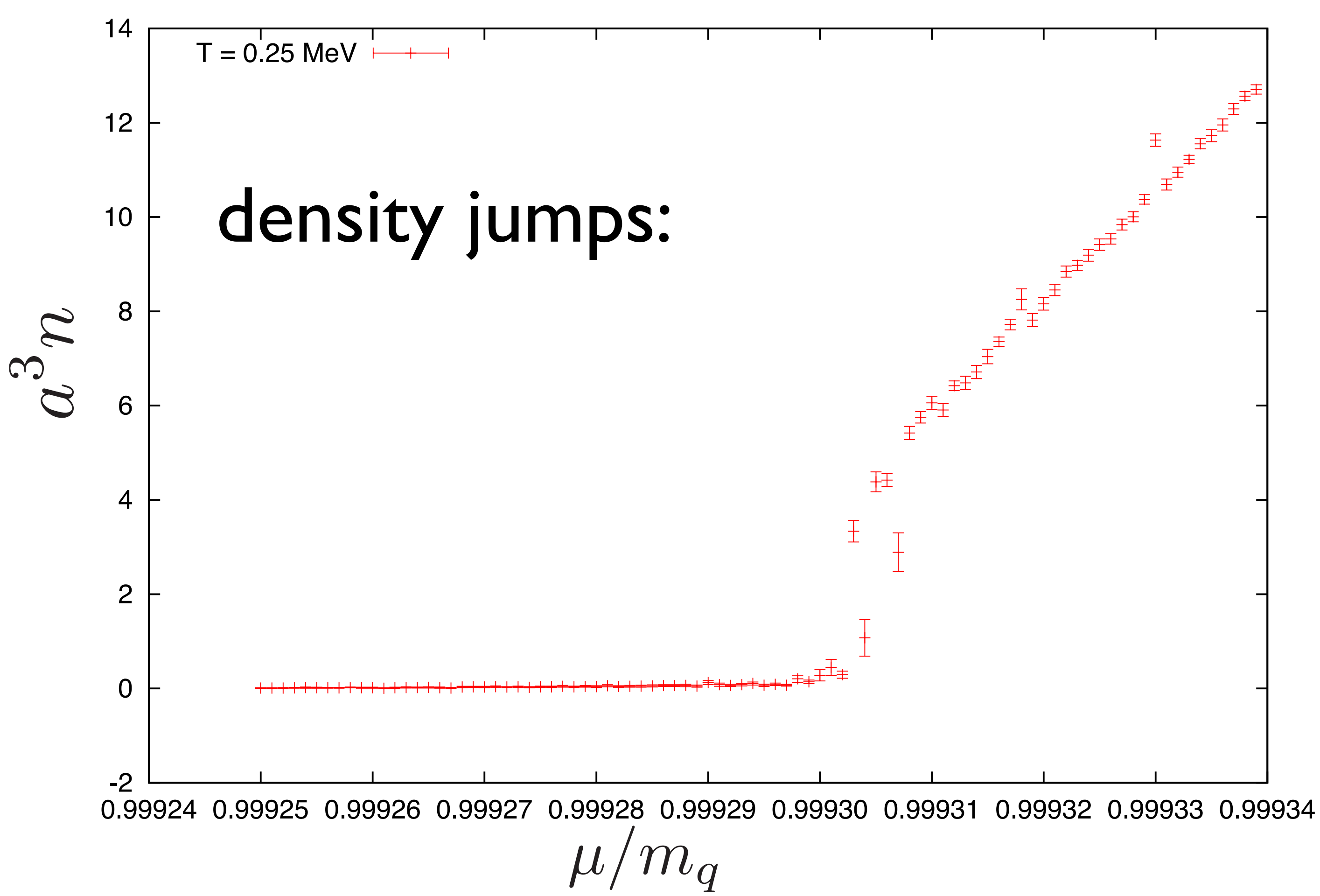


- Cut-off effects grow rapidly beyond onset transition: lattice saturation!
- Finer lattice necessary for larger densities!

Light quarks: first order transition + endpoint

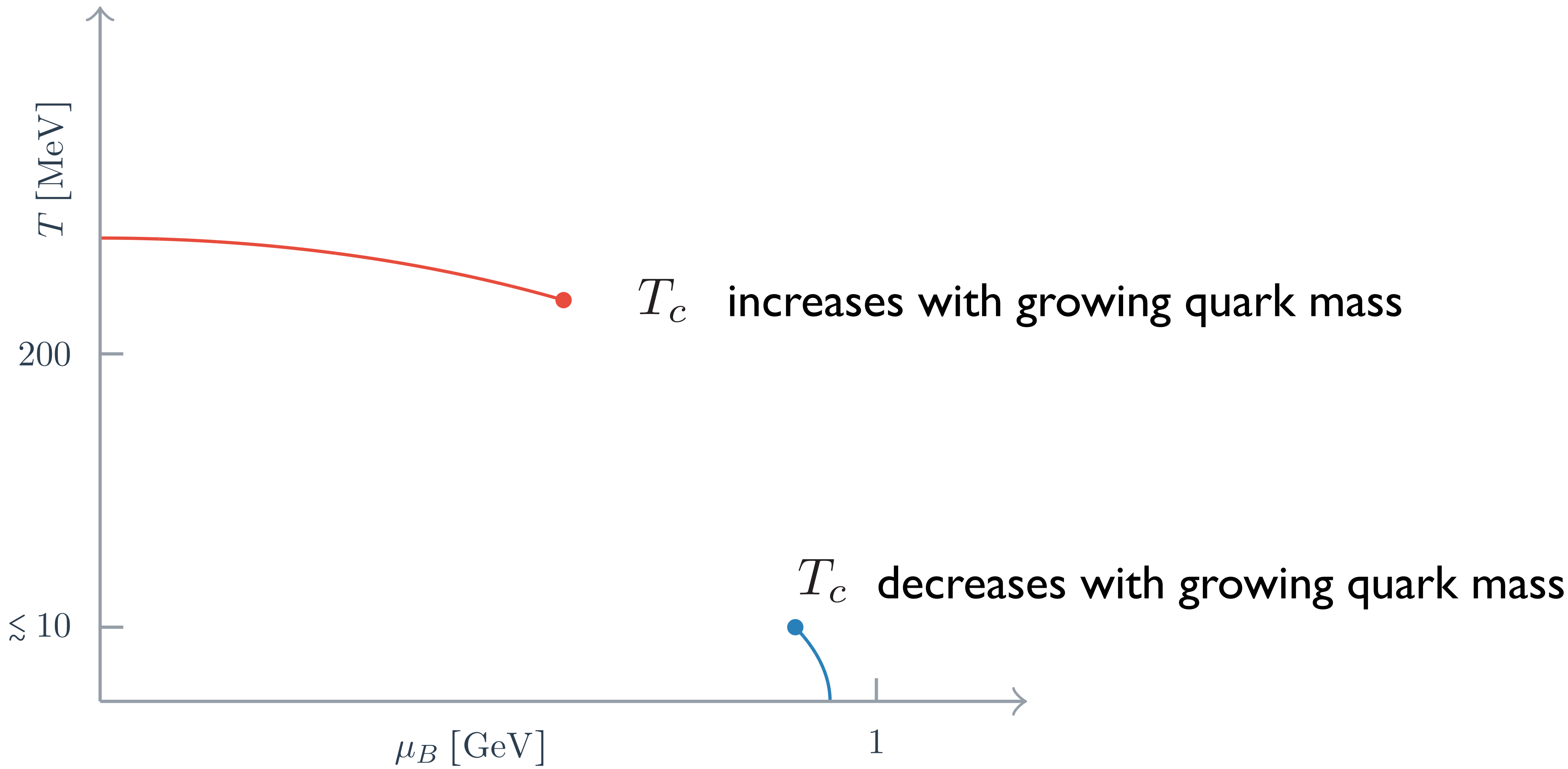


density jumps:



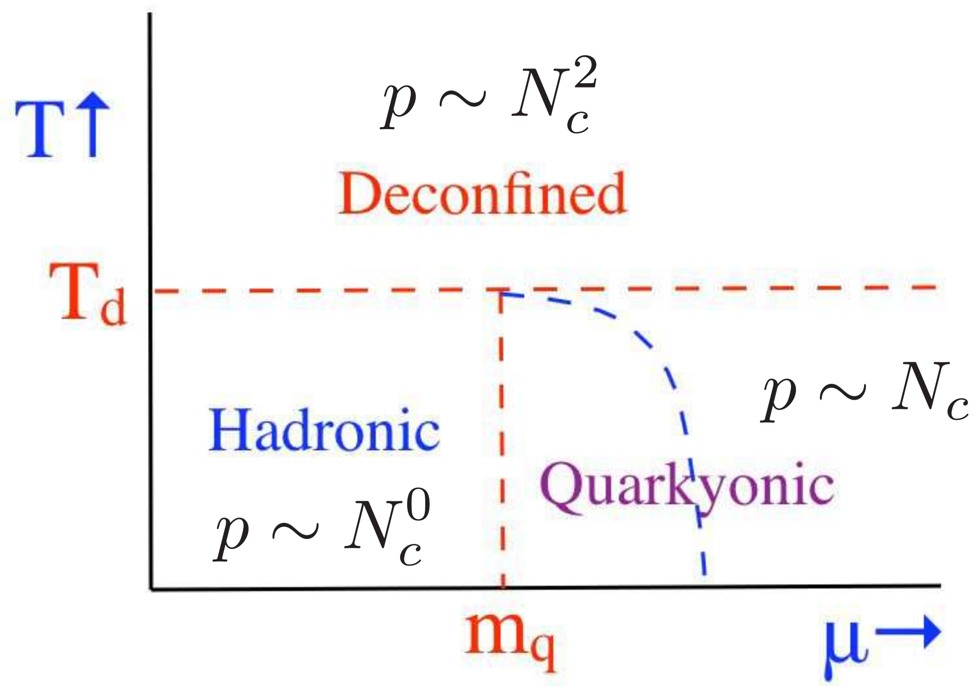
- phase coexistence: first order
- for higher $T = \frac{1}{aN_\tau}$ crossover
- nuclear liquid gas transition!

Phase diagram of heavy quark QCD



The large N_c QCD phase diagram

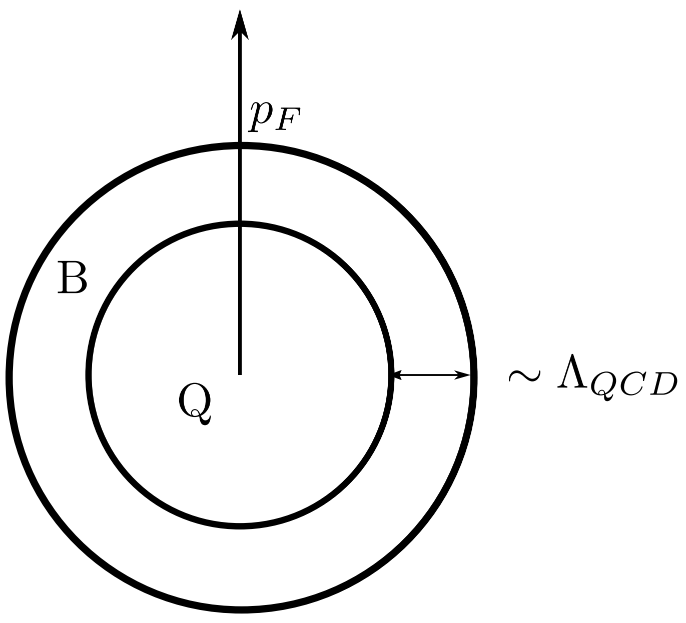
[McLerran, Pisarski NPA (2007), …]



Quarkyonic matter in momentum space:

Fermi sea of quarks, surrounded by Fermi shell of baryons;

$p_F \sim \mu$ can interpolate from purely baryonic to quark matter



From conjecture to calculation: eff. theory for general N_c

Strong coupling limit

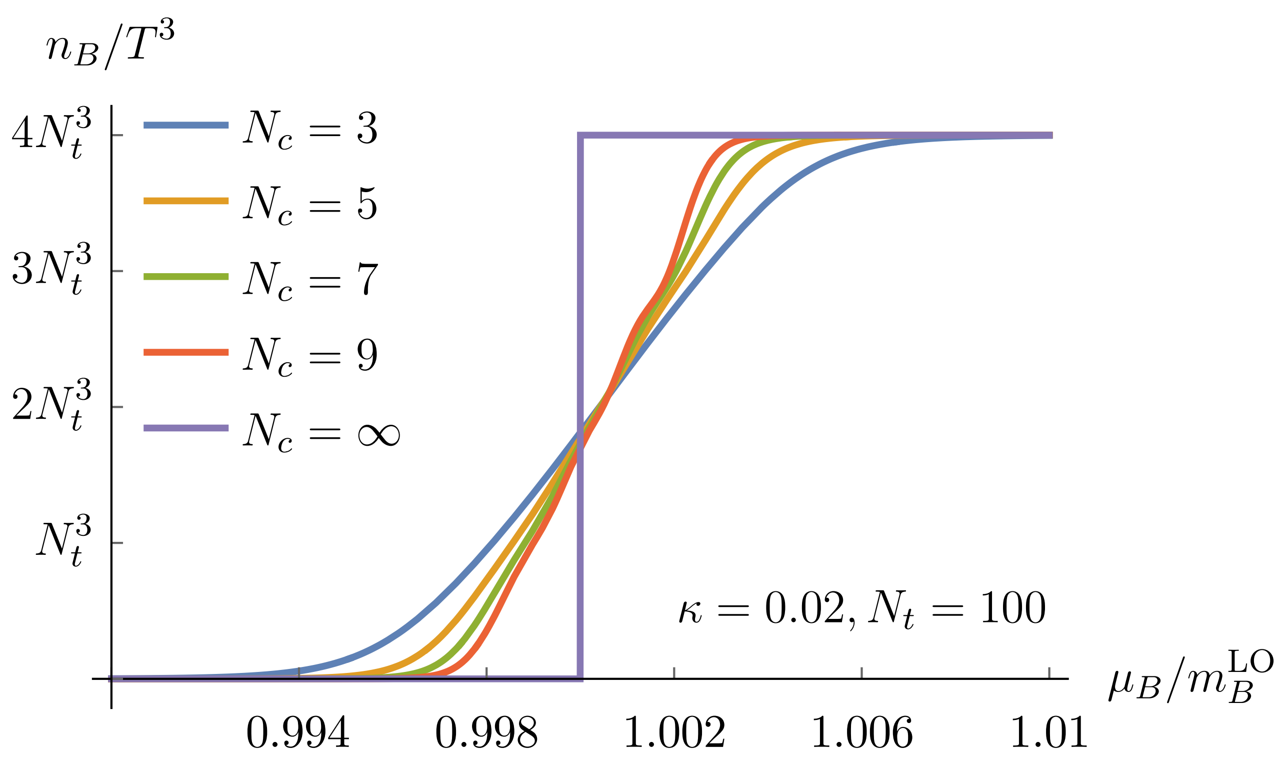
[O.P., Scheunert JHEP (2019)]

Order hopping expansion		κ^0	κ^2	κ^4
$h_1 < 1$ $(\mu_B < m_B)$	$a^4 p$	$\sim \frac{1}{6N_\tau} N_c^3 h_1^{N_c}$	$\sim -\frac{1}{48} N_c^7 h_1^{2N_c}$	$\sim \frac{3N_\tau \kappa^4}{800} N_c^8 h_1^{2N_c}$
	$a^3 n_B$	$\sim \frac{1}{6} N_c^3 h_1^{N_c}$	$\sim -\frac{N_\tau}{24} N_c^7 h_1^{2N_c}$	$\sim \frac{(9N_\tau+1)N_\tau}{1200} N_c^8 h_1^{2N_c}$
	$a^4 e$	$\sim -\frac{\ln(2\kappa)}{6} N_c^4 h_1^{N_c}$	$\sim \frac{N_\tau \ln(2\kappa)}{48} N_c^8 h_1^{2N_c}$	
	ϵ	0	$\sim -\frac{1}{4} N_c^3 h_1^{N_c}$	
$h_1 > 1$ $(\mu_B > m_B)$	$a^4 p$	$\sim \frac{4\ln(h_1)}{N_\tau} N_c$	$\sim -12 N_c$	$\sim 198 N_c$
	$a^3 n_B$	~ 4	$\sim -N_\tau \frac{N_c^4}{h_1^{N_c}}$	$\sim -\frac{(59N_\tau-19)N_\tau}{20} \frac{N_c^5}{h_1^{N_c}}$
	$a^4 e$	$\sim -4\ln(2\kappa) N_c$	$\sim 24 \ln(2\kappa) N_c$	
	ϵ	0	~ -6	

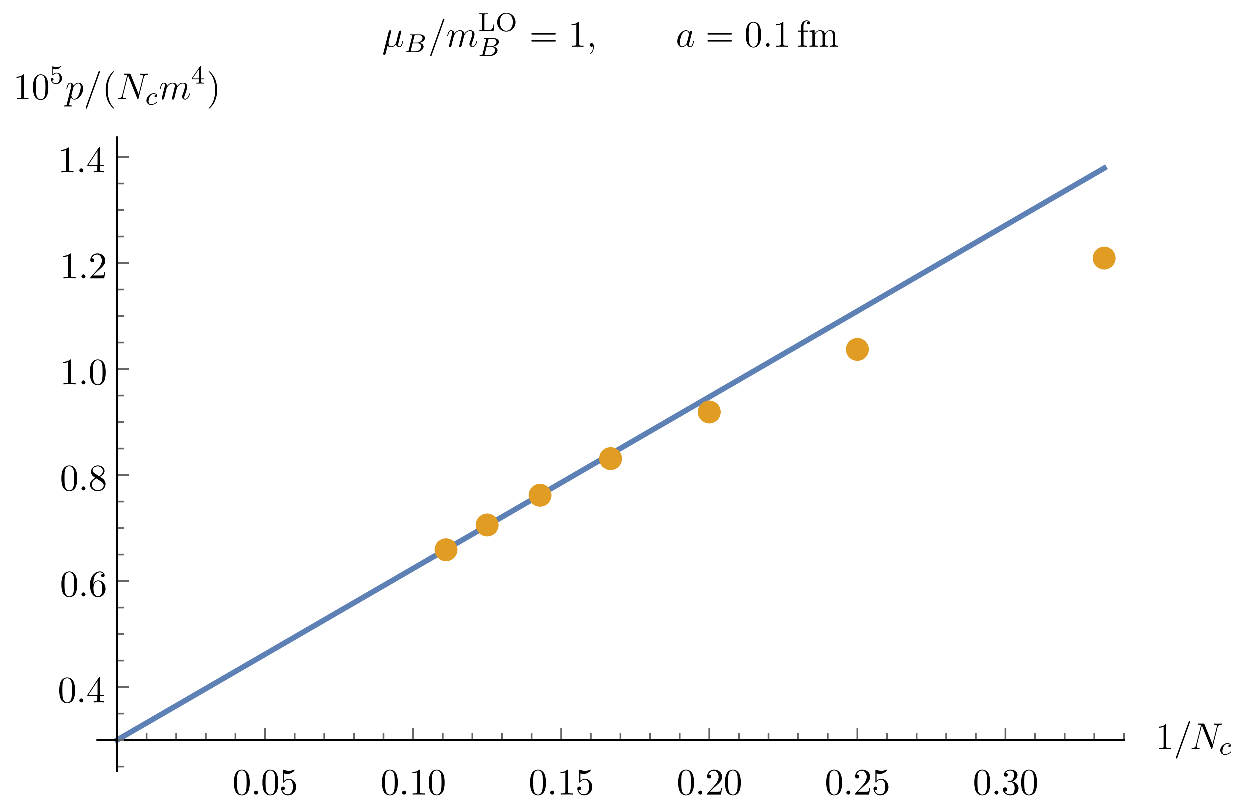
Beyond the onset transition: $p \sim N_c$ definition of quarkyonic matter!

The baryon onset transition for growing N_c

Transition becomes more strongly 1st-order for every T!

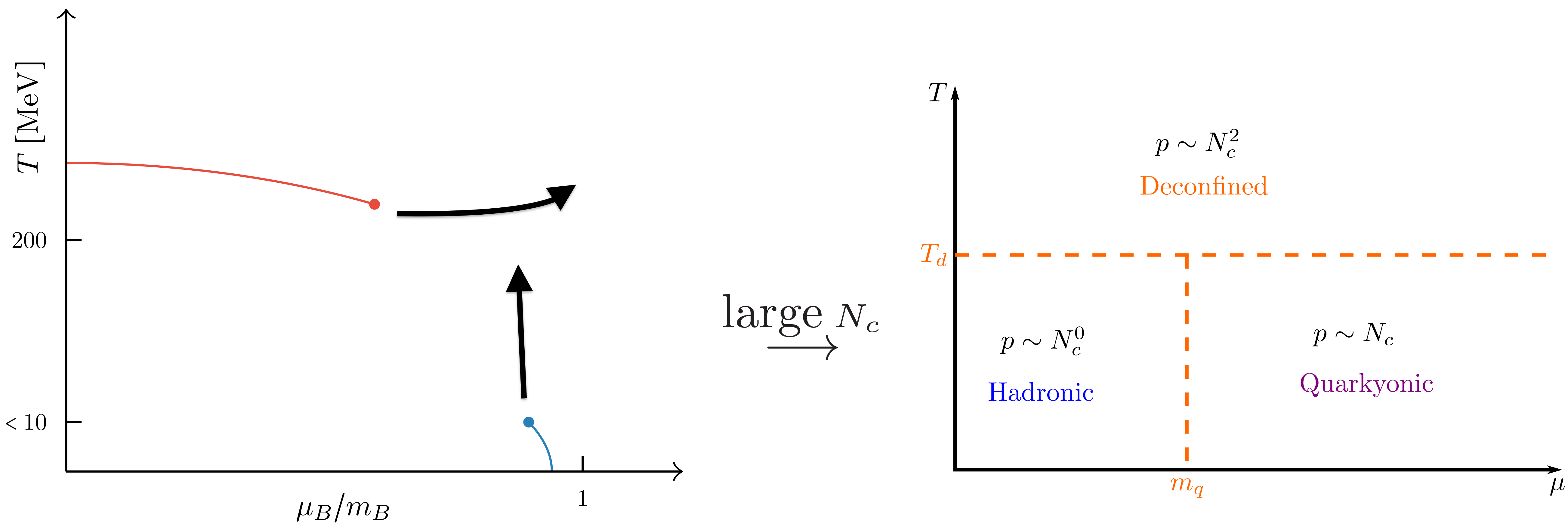


Pressure scaling right after onset



$$p \sim N_c(1 + \mathrm{const.} N_c^{-1})$$

Altogether:



- Large N_c phase diagram emerges smoothly
- Varying N_c: dense QCD is consistent with quarkyonic matter
- Should also hold for light quarks, N_c-scaling is property of expansion coefficients!

Baryon matter is special case of quarkyonic matter

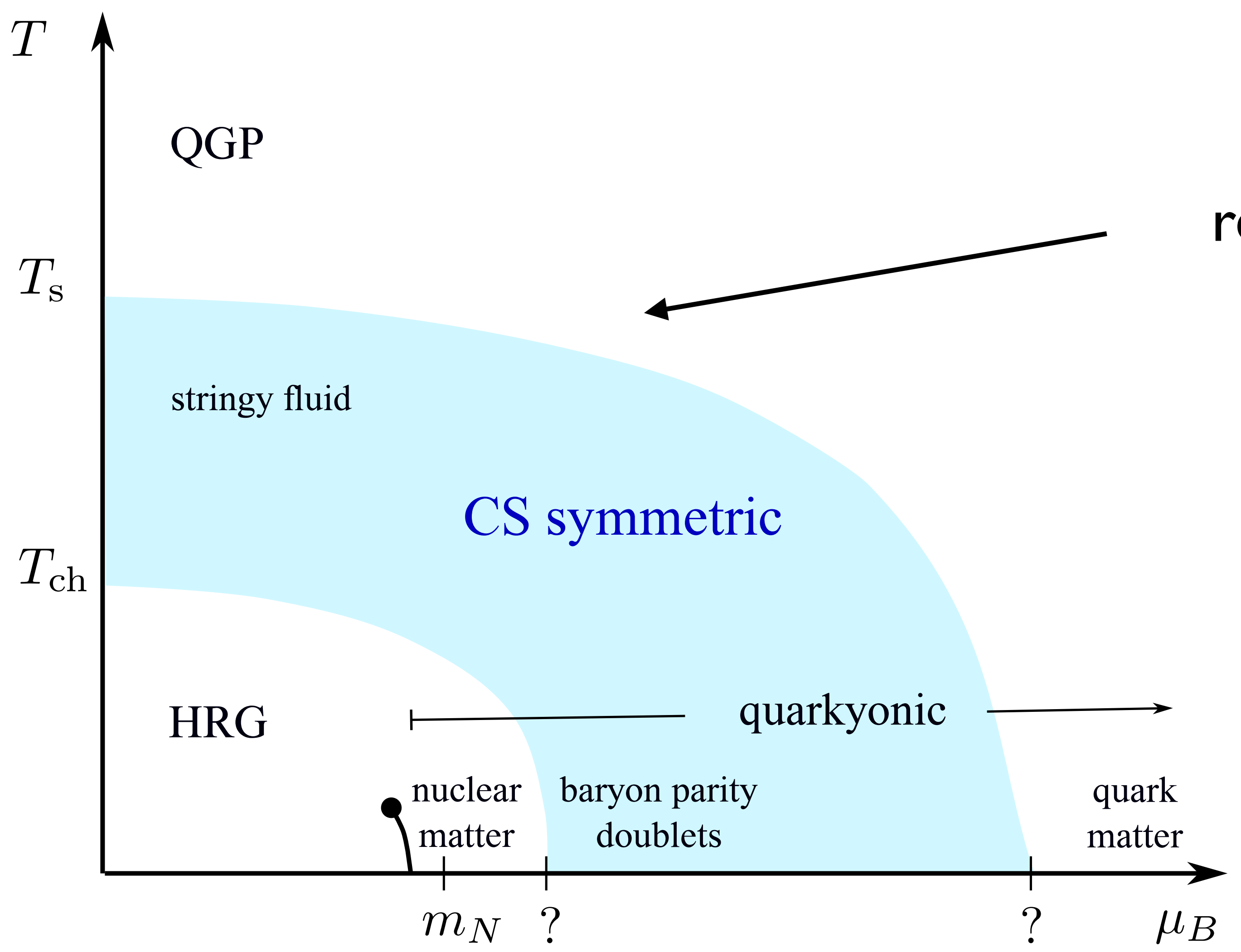
Phenomenological evidence: [Koch, McLerran, Miller, Vovchenko, PRC 24]

Implications for physical QCD?

One viable scenario (more possibilities): [L. Glozman, R. Pisarski, O.P., EJPA 22]

$\mu\bar{\psi}\gamma_0\psi$ is CS symmetric

resonance-like states



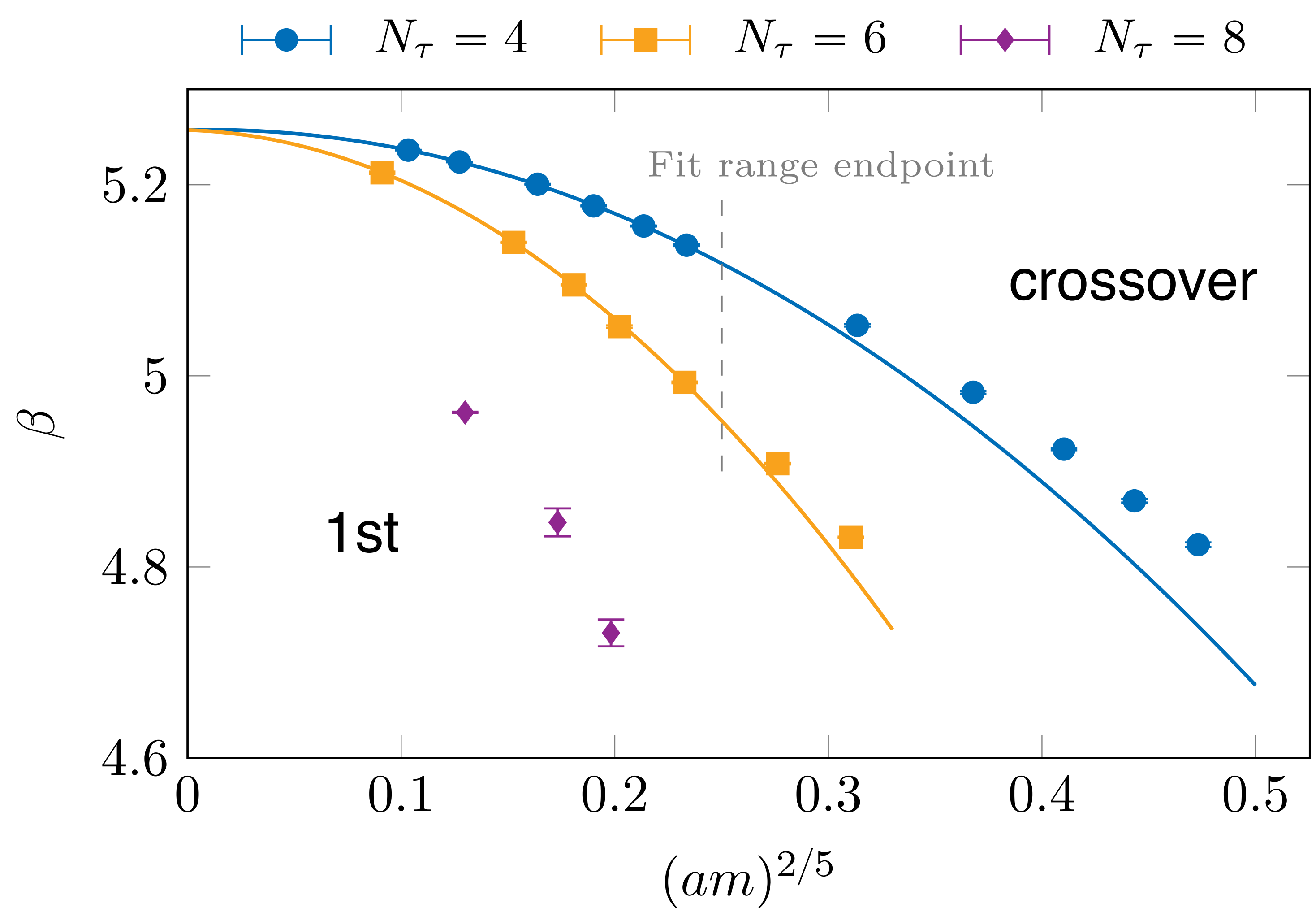
Conclusions

- Chiral transition at zero density is likely 2nd order for Nf=2-7 massless quark flavours
- Imaginary chemical potential has no effect on the order of the chiral transition
- Evidence for chiral-spin symmetry:
  screening masses, spectral functions resonance-like
- Heavy mass LQCD consistent with quarkyonic matter and quark-hadron continuity

Backup slides

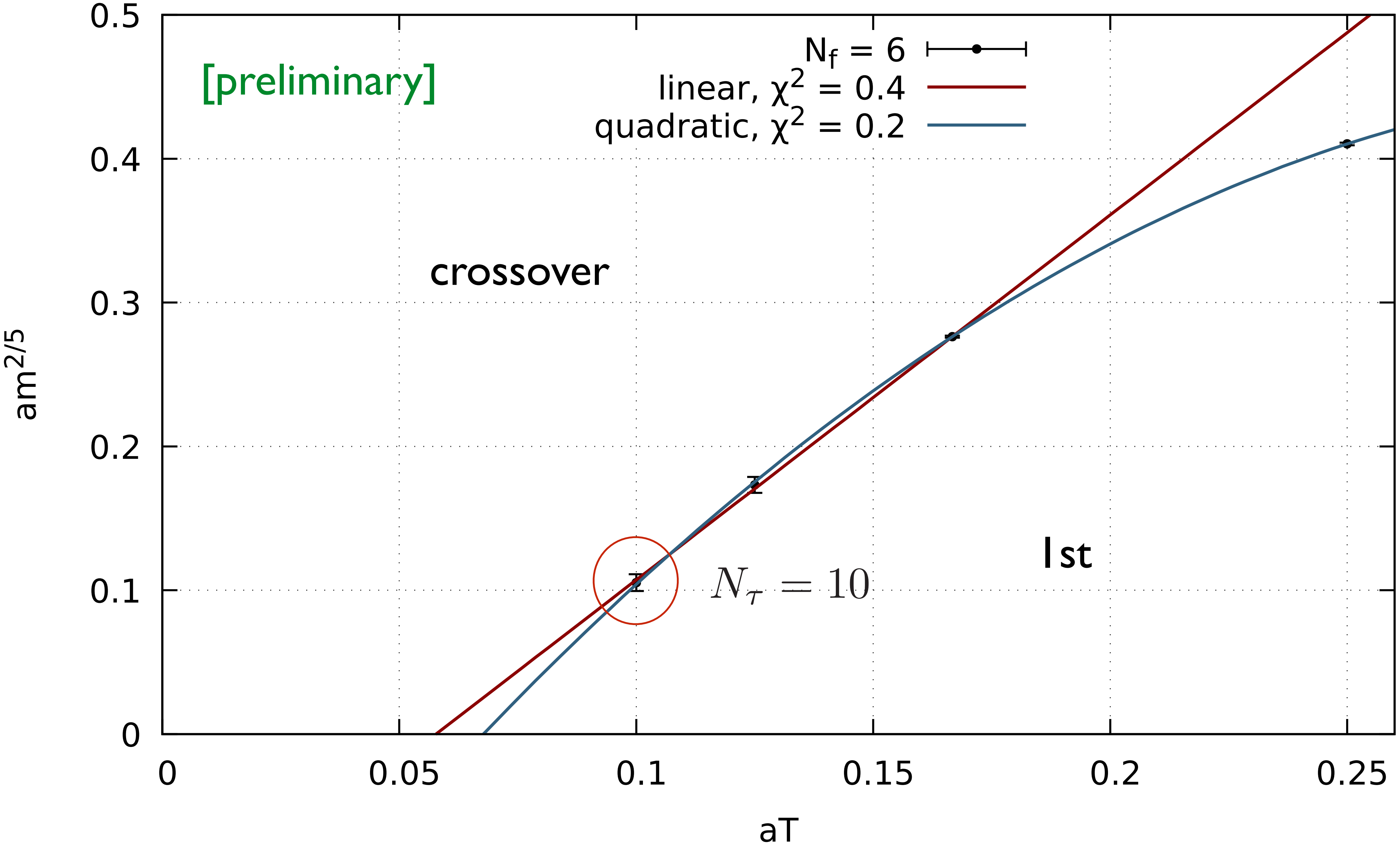
Bare parameter space of unimproved staggered LQCD

[Cuteri, O.P., Sciarra JHEP 21] ~120 M Monte Carlo trajectories with light fermions, aspect ratios 3,4,5



- Data points implicitly labeled by Nf
- Tricritical scaling observed in lattice bare parameter space
- Tricritical extrapolation always possible!

progressing to finer lattices



New $N_\tau = 10$ result on predicted scaling curve!

Thermal spectral density + thermoparticles

- The thermal spectral density $\widetilde{D}_{\beta}(\boldsymbol{u},s)$ holds the key to understanding in-medium phenomena, but what structure does it have?
- A natural decomposition [Bros, Buchholz, *NPB* 627 (2002)] is:

$$\widetilde{D}_{\beta}(\vec{u},s) = \widetilde{D}_{m,\beta}(\vec{u})\,\delta(s-m^2) + \widetilde{D}_{c,\beta}(\vec{u},s)$$

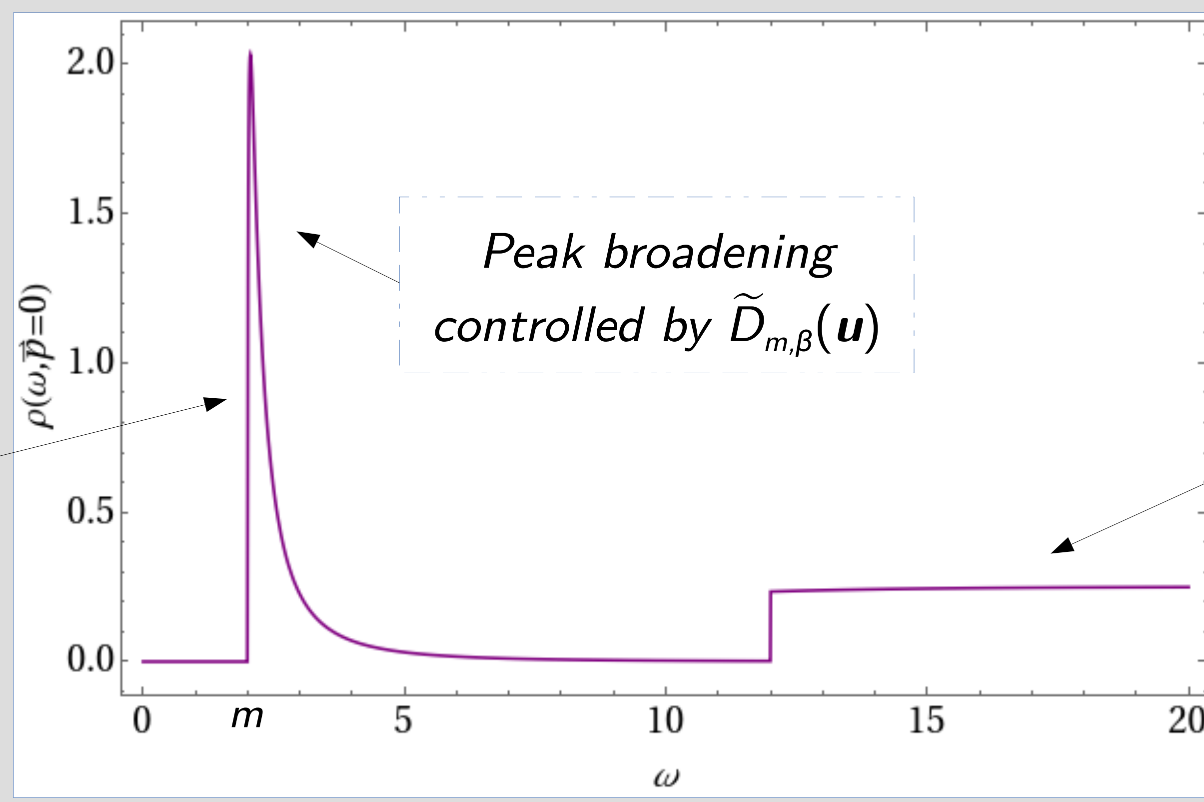
(Negligible at low T)

Thermoparticle component

"Damping factor"

Continuous component

Causes $T=0$ mass pole m to be ***screened*** *by thermal effects*



Fixes T-dependence of continuous spectral contributions

[slide by P. Lowdon]

Comparison with plasmon ansatz

Bros+Buchholz Ansatz

Perturbative plasmon: Breit-Wigner shape

Both fit spatial correlator

$$\rho_{\rm PS}(\omega,\mathbf{p}=0)=\epsilon(\omega)\left[\theta(\omega^2-m_\pi^2)\,\frac{4\,\alpha_\pi\,\gamma_\pi\sqrt{\omega^2-m_\pi^2}}{(\omega^2-m_\pi^2+\gamma_\pi^2)^2}+\theta(\omega^2-m_{\pi^*}^2)\,\frac{4\,\alpha_{\pi^*}\,\gamma_{\pi^*}\sqrt{\omega^2-m_{\pi^*}^2}}{(\omega^2-m_{\pi^*}^2+\gamma_{\pi^*}^2)^2}\right]$$

$$\rho_{PS}^{BW}(\omega,\mathbf{p}=0)=\frac{4\alpha_\pi\omega\Gamma_\pi}{(\omega^2-m_\pi^2-\Gamma_\pi^2)^2+4\omega^2\Gamma_\pi^2}+\frac{4\alpha_\pi^*\omega\Gamma_{\pi^*}}{(\omega^2-m_{\pi^*}^2-\Gamma_{\pi^*}^2)^2+4\omega^2\Gamma_{\pi^*}^2}$$

Predicted temporal correlators: