

## The helium-4 abundance BSM: without solutions

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### I. THE EXERCISE: CALCULATING THE HELIUM-4 ABUNDANCE AFTER THE BIG BANG

The neutron abundance in the early Universe is governed by the following processes:

$$n \leftrightarrow p + e^- + \bar{\nu}_e \quad (1a)$$

$$n + \nu_e \leftrightarrow p + e^- \quad (1b)$$

$$n + e^+ \leftrightarrow p + \bar{\nu}_e \quad (1c)$$

Based on dimensional analysis, the rate of these interactions roughly scale as  $\Gamma \simeq G_F^2 T^5$  and hence they will freeze-out at  $T \sim 1$  MeV. In the Standard Model, effectively all of the neutrons in the plasma end up forming  ${}^4\text{He}$  because it is the most tightly bounded light nuclei. This means that if we are able to track the neutron abundance of the Universe we will know how much  ${}^4\text{He}$  should have been generated after the Big Bang. By contrasting this prediction with observations of the helium mass fraction in the Universe,  $Y_p \equiv \rho_{{}^4\text{He}}/\rho_B = 0.245 \pm 0.003$  [PDG-24], we can set bounds on an array of BSM models, including to the parameter  $N_{\text{eff}}$ . This exercise will guide you through the basics of the calculation.

1. Assuming that there are only neutrons and protons in the plasma, write down the Boltzmann equation for the neutron and proton number densities given processes that interconvert them. Write these equations as a function of the rate of neutron conversion  $\lambda_{np}$  and proton conversion  $\lambda_{pn}$ .
2. Given these expressions write the equation for the fraction of neutrons in the plasma, i.e.  $X_n \equiv n_n/(n_n + n_p)$ .
3. Assuming that the rates of proton-to-neutron and neutron-to-proton conversion are very efficient, this is, we are in thermal and chemical equilibrium, obtain the relationship between the  $\lambda_{pn}$  and  $\lambda_{np}$  reaction rates. *Tip:* For this you want to use chemical equilibrium of neutron-proton interactions (i.e.  $\mu_n = \mu_p$ ) and the formula for Maxwell-Boltzmann number densities.
4. The neutron-to-proton interaction rates require a phase space integration over the distribution functions of  $e$  and  $\nu_e$ . A simplified (but still rather accurate) formula for it is given by:

$$\lambda_{n+\nu_e \rightarrow p+e^-} = \lambda_{n+e^+ \rightarrow p+\bar{\nu}_e} \simeq 15 \frac{1}{\tau_n} \frac{T^4 \Delta m}{m_e^5}, \quad (2)$$

where  $\tau_n = 878.4 \pm 0.5$  s is the neutron lifetime,  $\simeq 15$  mins. Note that at  $T \sim$  MeV these rates are  $\sim 10^3$  times faster than neutron decay  $\Gamma_n = 1/\tau_n$ .

Remember that we can write the Hubble parameter as

$$H = 1.66 \sqrt{g_\star} \frac{T^2}{M_{\text{Pl}}}. \quad (3)$$

With this, pose the differential equation for  $X_n$  as a function of  $T$  as the time variable. One can use the simplification that  $dT/dt = -HT$ .

5. Numerically solve this system using  $g_\star = 10.75 + 2\frac{7}{8}\Delta N_{\text{eff}}$  and find the resulting solution for  $X_n$  without taking into account for the moment  $n$  decay. What do you find for  $\Delta N_{\text{eff}} = -1, 0, +1$ ? What is the fractional change to this neutron fraction?
6. Finally, we need to relate the neutron abundance in the Universe to the abundance of Helium-4. In order to form Helium-4, a substantial amount of D, T and  ${}^3\text{He}$  needs to be synthesized before hand. Critically, any ‘‘heavy’’ element formation requires the presence of substantial amounts of deuterium and this means that BBN will happen quickly once the abundance of deuterium is  $\mathcal{O}(1)$ . We can then approximate the abundance of helium-4

as given by the abundance of all the neutrons in the plasma once there is a large number of deuterium in the plasma and BBN starts. This is:

$$n_{4\text{He}} = n_n/2, \quad (4)$$

since what is typically quoted is the mass fraction of helium then one finds:

$$Y_P = 4 \frac{n_{4\text{He}}}{n_B} = 2X_n|_{T=T_{\text{BBN}}}. \quad (5)$$

The goal now is to estimate this temperature,  $T_{\text{BBN}}$ . In the plasma there are processes  $n + p \leftrightarrow D + \gamma$  occurring efficiently and this tells us that  $\mu_n + \mu_p = \mu_D$ . Again, by using the Maxwell-Boltzmann number density formula we can then relate the number densities of  $n$ ,  $p$ ,  $\gamma$  and  $D$ . Find this relationship in terms of the deuterium fraction, i.e.  $X_D = n_D/n_B$  where  $n_B$  is the number of baryons. Note that  $g_D = 3$  (because Deuterium is a vector state). Once you have that formula solve for when  $X_D \simeq 10^{-2}$  to find the temperature at which BBN starts. You should find  $T_{\text{BBN}} \simeq 0.072 \text{ MeV}$ .

7. Given that we know the temperature at which BBN starts we can calculate the final helium abundance using Eq. (15). The only thing that changes between neutron-proton freeze-out is that the neutrons are decaying. This can be easily incorporated by taking:

$$Y_P = 2X_n(T \rightarrow 0) \times e^{-t_{\text{BBN}}/\tau_n} \quad (6)$$

The only piece we have left is simply to get the time at which this corresponds. Remember that in a radiation dominated Universe we have  $t \simeq 1/(2H)$ . Taking into account entropy conservation find that at  $T_{\text{BBN}} \simeq 0.072 \text{ MeV}$  almost all  $e^+e^-$  have annihilated and therefore  $g_\star \simeq 3.36 + 0.46\Delta N_{\text{eff}}$ . With this, calculate  $Y_P$  for  $\Delta N_{\text{eff}} = -1, 0, +1$ . What is the fractional change to this neutron fraction? Given the observed value of  $Y_p = 0.245 \pm 0.003$  what bound can be derived on  $\Delta N_{\text{eff}}$ ?