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Quantum Simulation of the Quantum Hall Effect

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In collaboration with:

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Cold atomic quantum Hall effect – Why?

Systems with well controlled Hamiltonians

Quantum Hall solver

- Hard problem: Competition between different phases

Novel quantum Hall phases

- Quantum Hall effect of bosons
- *Interacting* integer quantum Hall phases
- Ideal phases (parent Hamiltonians)

Exploring anyonic properties

- First experimental detection of fractional statistics
- Braiding of (non-Abelian) anyons
- Anyon technologies

Outline

1. Quantum Hall Physics – in general:

- Single-particle physics: Landau levels
- Many-body effects: Trial states

2. Fractional quantum Hall physics of spin-1/2 bosons:

- Abelian vs. Non-Abelian phases
- Numerical results: ambiguous
- Prospects of a quantum simulation

3. *Integer* quantum Hall physics of spin-1/2 bosons:

- Interactions are crucial
- Edge spectrum as a fingerprint

4. Anyon braiding in small systems

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Quantum Hall Systems

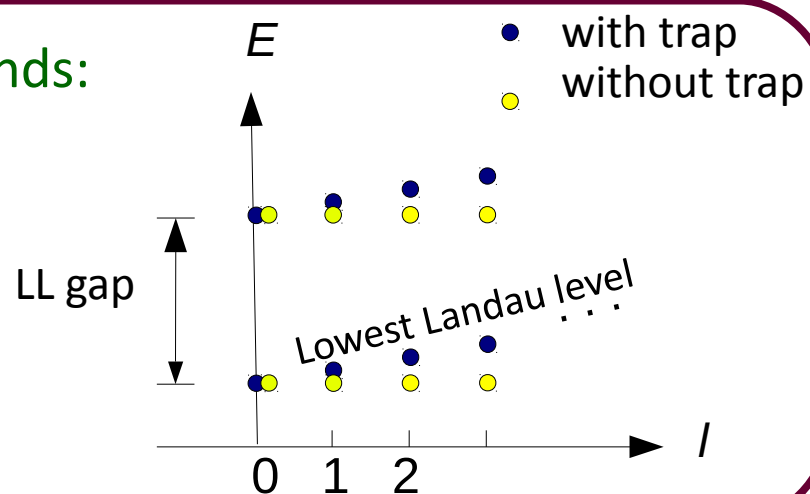
$$H = \frac{(\mathbf{p} + \mathbf{A})^2}{2M} + \frac{M}{2} \omega^2 \mathbf{r}^2$$

Gauge potential:

$$\mathbf{A} = \frac{B}{2}(y, -x)$$

Trapping potential

Flat bands:



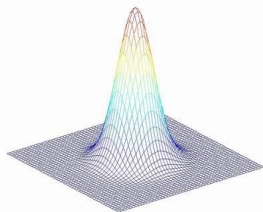
Fermions:

Integer filling:
 → effectively non-interacting
 → Integer Quantum Hall Phases

Fractional filling:
 → Gapped phases due to interactions
 → Fractional Quantum Hall Phases

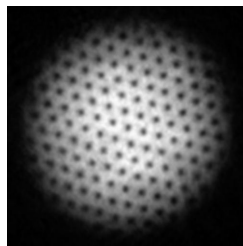
Bosons:

Condensate



→
symmetry breaking

Vortex lattice



→
melting

Interacting Quantum Hall Phases:

→ fractional
 (like the fermionic ones, with or without spin)

→ integer
 (no fermionic counterpart, needs spin)

Trial states: Laughlin and Halperin

Wave functions with “zeros” for all particle pairs:

→ Laughlin wave function (spinless system)

$$\Psi_L^{(q)} = \prod_{i < j} (z_i - z_j)^q \exp[-\sum_i |z_i|^2/2] \quad \text{filling } \nu = 1/q$$

$(z = x + iy)$



→ Halperin wave function (two-component system)

$$\Psi_H^{(lmn)} \sim \prod_{1 \leq i < j \leq N_\uparrow} (z_{i\uparrow} - z_{j\uparrow})^l \prod_{1 \leq i < j \leq N_\downarrow} (z_{i\downarrow} - z_{j\downarrow})^m \prod_{\substack{1 \leq i \leq N_\uparrow \\ 1 \leq j \leq N_\downarrow}} (z_{i\uparrow} - z_{j\downarrow})^n$$

fillings $\nu_\uparrow = \frac{l - n}{lm - n^2}$ and $\nu_\downarrow = \frac{m - n}{lm - n^2}$



Exact zero-energy solutions in contact potential!

Trial states: Pairing states

- 1) Divide system into k clusters.
- 2) Each cluster forms a Laughlin/Halperin state.
- 3) (Anti-)Symmetrize over all possible clusters.

→ Read-Rezayi series (spinless):

$$\Psi_{\text{RR}}^{(k)} \sim \mathcal{S}[\Psi_{\text{L}}^{(2)}(z_{i_1}, \dots, z_{i_M}) \Psi_{\text{L}}^{(2)}(z_{i_{M+1}}, \dots, z_{i_{2M}}) \dots]$$

filling $\nu = k/2$

Moore/Read (1991)
Read/Rezayi (1999)

→ Non-Abelian spin singlet (NASS) series

Ardonne/Schoutens (1999)

$$\Psi_{\text{NASS}}^{(k)} \sim \mathcal{S}[\Psi_{\text{H}}^{(221)}(z_{i_1\uparrow}, \dots, z_{i_M\uparrow}, z_{i_1\downarrow}, \dots, z_{i_M\downarrow}) \Psi_{\text{H}}^{(221)}(z_{i_{M+1}\uparrow}, \dots, z_{i_{2M}\uparrow}, z_{i_{M+1}\downarrow}, \dots, z_{i_{2M}\downarrow}) \dots]$$

filling $\nu = 2k/3$

Exact ground states for $(k+1)$ -body contact interactions!

Trial states: Composite fermion states

Construction Recipe:

1. Composite fermion = particle + m magnetic fluxes

→ Jastrow factor: $J(z) = \prod_{i>j} (z_i - z_j)^m$

2. CFs fill Landau levels at modified magnetic field

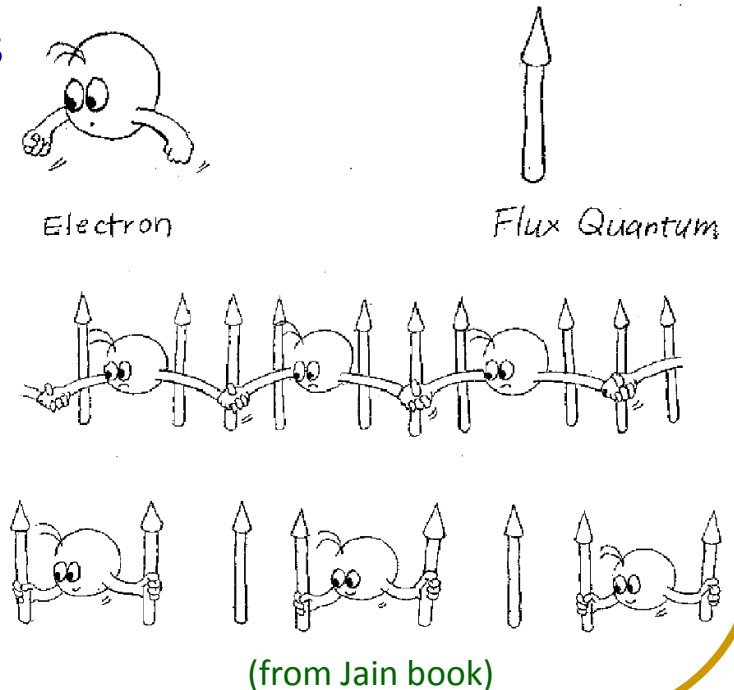
→ Slater determinant ϕ of filled LLs

3. Project back into low-energy space:

Lowest Landau level of the original system

$$\Psi_{\text{CF}} = \mathcal{P}_{\text{LLL}} \Phi(z) J(z)$$

Jain/Kawamura (1995)
Cooper/Wilkin (1999)



Construction works for fermionic and bosonic systems with or without spin, at filling factors $\nu = \frac{n}{mn \pm 1}$ where the number m of attached fluxes per particle must be even for fermions or odd for bosons.

Trial states: Overview

	Spinless fermions	Spinless bosons	Two-component bosons (fully unpolarized)
Abelian Fractional Quantum Hall States	<p>Laughlin</p> $\nu = 1/q, q \text{ odd}$	<p>Laughlin</p> $\nu = 1/q, q \text{ even}$	<p>Halperin</p> $\nu = \frac{2}{m+n}, m \text{ even}$
	<p>CF states</p> $\nu = \frac{n}{mn+1}, m \text{ even}$	<p>CF states</p> $\nu = \frac{n}{mn+1}, m \text{ odd}$	<p>CF states</p> $\nu = \frac{n}{n \pm 1} \notin \mathbb{N}$
Non-Abelian Fractional Quantum Hall States	<p>Read-Rezayi</p> $\nu = \frac{k}{k+2}$	<p>Read-Rezayi</p> $\nu = \frac{k}{2}$	<p>NASS</p> $\nu = \frac{2k}{3}$
Integer Quantum Hall States	<p>trivial</p>	<p>×</p>	<p>CF state</p> $\nu = 2$

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4. Anyon braiding in small systems

The System

We now focus on:

- bosons
- pseudospin-1/2
- in the lowest Landau level
- with contact interactions:

$$H = \sum_{i < j} \left[g_{\uparrow\uparrow} \delta(z_{i\uparrow} - z_{j\uparrow}) + g_{\downarrow\downarrow} \delta(z_{i\downarrow} - z_{j\downarrow}) + g_{\uparrow\downarrow} \delta(z_{i\uparrow} - z_{j\downarrow}) + g_{\uparrow\downarrow} \delta(z_{i\downarrow} - z_{j\uparrow}) \right]$$

- SU(2)-symmetric: $g_{\uparrow\uparrow} = g_{\uparrow\downarrow} = g_{\downarrow\downarrow}$

Numerical studies on different geometries:

Disk	Torus	Sphere
<ul style="list-style-type: none">• most realistic• edge effects	<ul style="list-style-type: none">• Purely bulk physics• Complicated wave functions	<ul style="list-style-type: none">• Purely bulk physics• Relatively simple wave func.• Shifted filling factors

NASS series on the torus?

RAPID COMMUNICATIONS

PHYSICAL REVIEW A **86**, 031604(R) (2012)

Quantum Hall states in rapidly rotating two-component Bose gases

Shunsuke Furukawa and Masahito Ueda

Department of Physics, University of Tokyo, 7-3-1 Hongo, Bunkyo-ku, Tokyo 113-0033, Japan

RAPID COMMUNICATIONS

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Non-Abelian spin-singlet states of two-component Bose gases in artificial gauge fields

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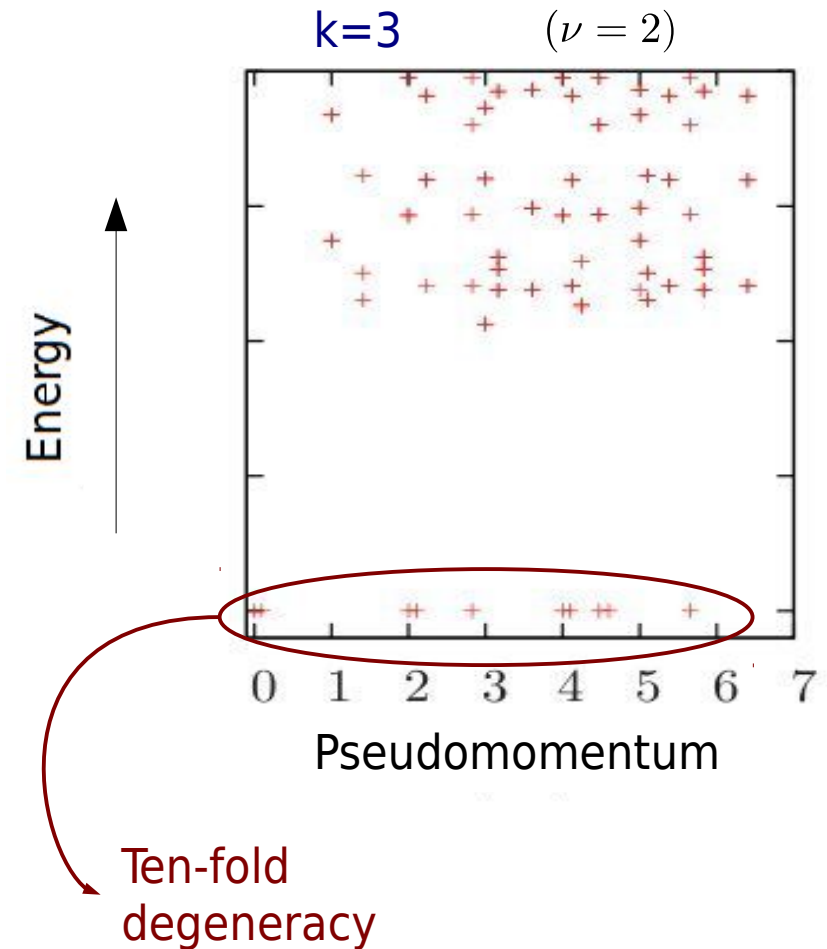
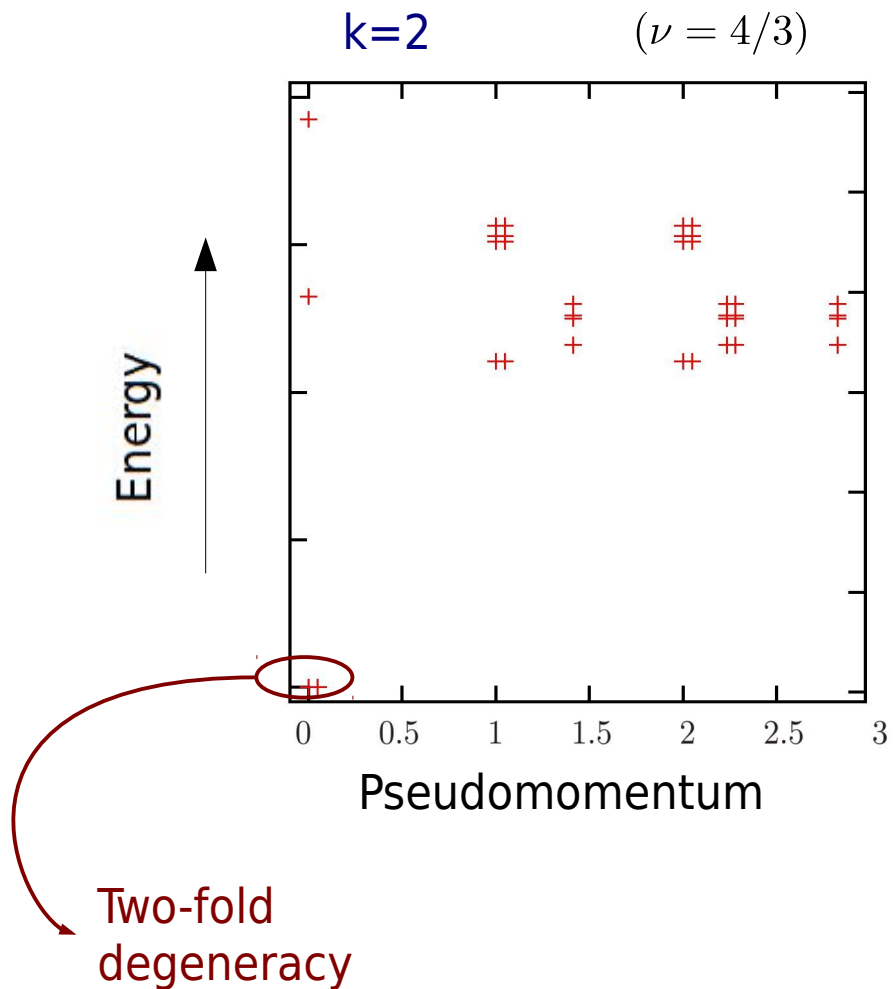
³*ICREA-Institució Catalana de Recerca i Estudis Avançats, 08010 Barcelona, Spain*

Exact diagonalization on the torus:

- Evidence of incompressible (gapped) phases at $\nu = \frac{2k}{3}$ for $k = 1, 2, 3$.
- NASS series?

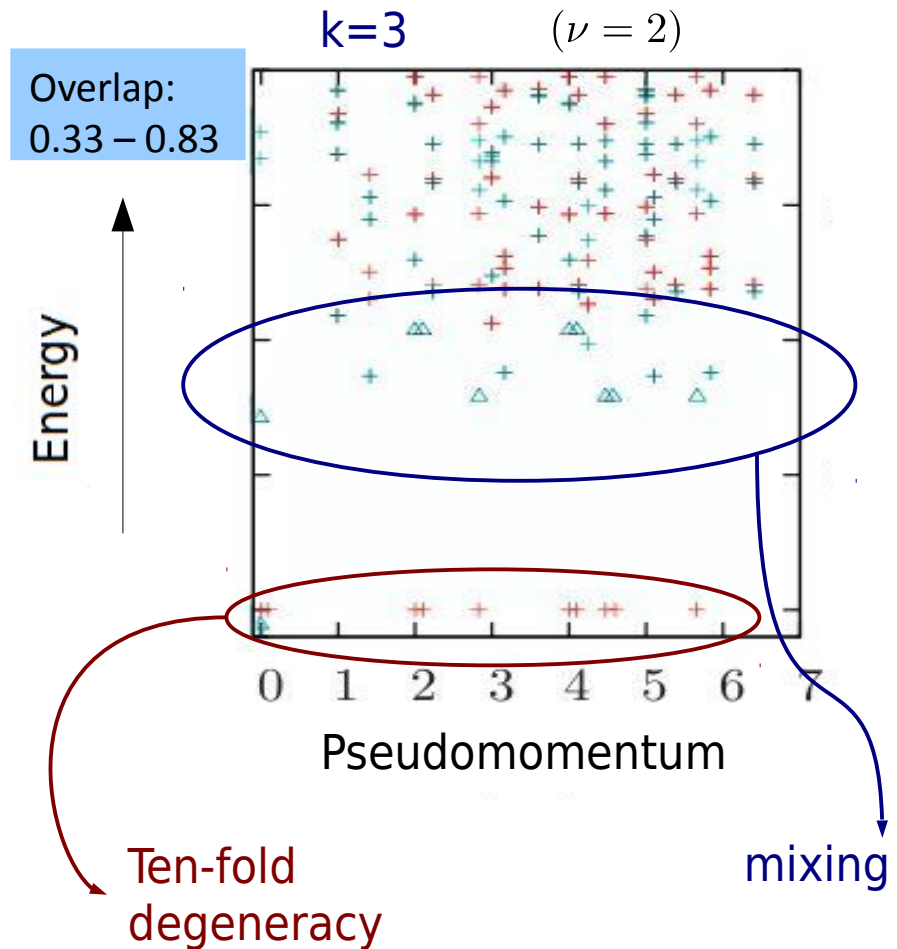
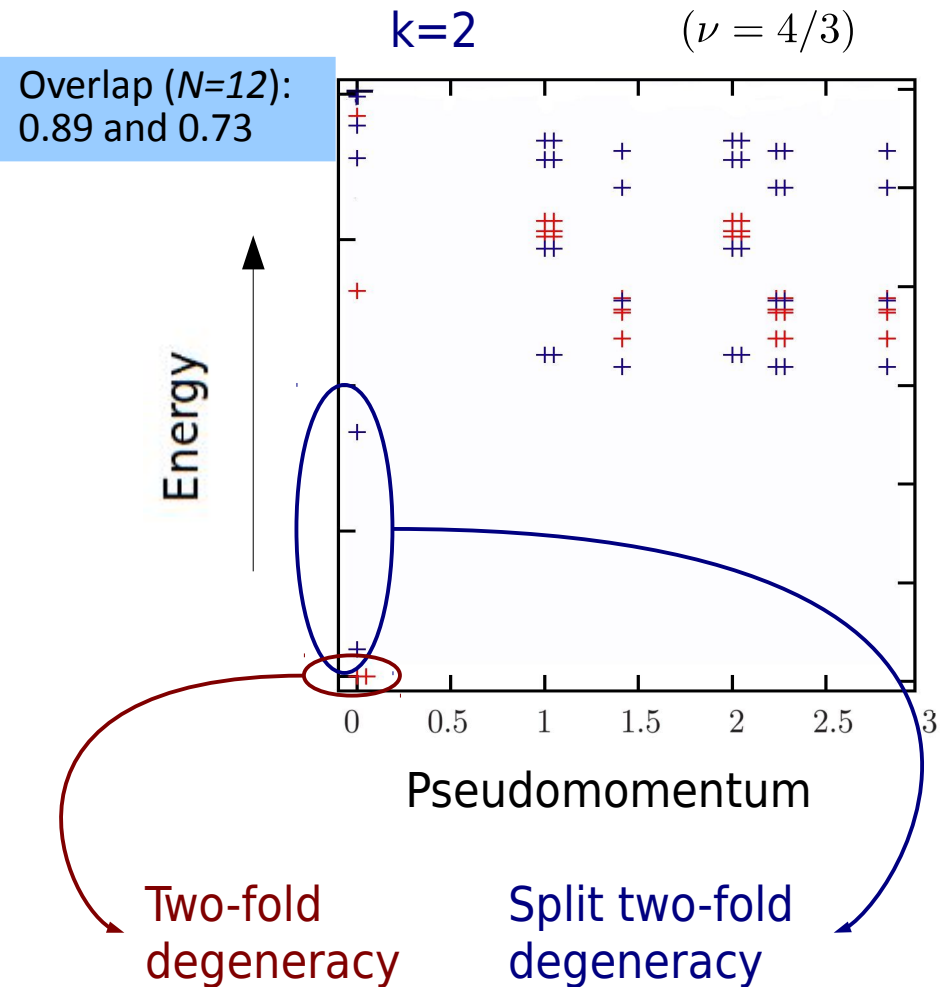
NASS series on the torus?

Spectra of $(k+1)$ -body contact interaction



NASS series on the torus?

Spectra of $(k+1)$ -body contact interaction
versus
Spectra of two-body contact interaction



CF states on the sphere?

ED on torus:

NASS phase
at $\nu=4/3$

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Different picture on the sphere!

PHYSICAL REVIEW B **87**, 245123 (2013)

Quantum Hall effect of two-component bosons at fractional and integral fillings

Ying-Hai Wu and Jainendra K. Jain

Department of Physics, The Pennsylvania State University, University Park, Pennsylvania 16802, USA

CF states on the sphere?

Overlaps on the sphere:

- with NASS state: 0.918
- with CF state: 0.985

for $N=12$ at filling $\nu=4/3$.

BUT: Filling factor is biased on the sphere.

$$\nu = \frac{N}{N_V} + \delta$$

Direct competition between NASS and CF is not possible on the sphere.
(Neither on small disks!)

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To solve the competition, one might study:

- Overlaps on the torus
- System on a disk → Quantum simulation!

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Cold atom quantum simulation

All ingredients of the Hamiltonian are available:

- ✓ Synthetic magnetic fields
- ✓ 2-body contact potential

But how could a quantum simulation distinguish between different states?

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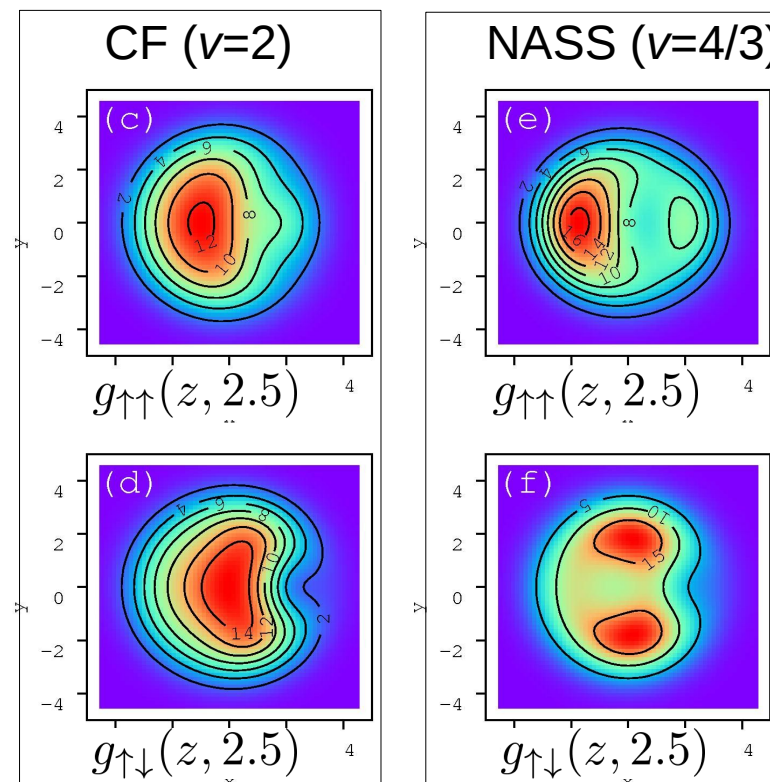
But how could a quantum simulation distinguish between different states?

Example:

$$N_{\uparrow} = 4$$

$$N_{\downarrow} = 4$$

$$L = 16$$



Correlation functions: $C(z_1, z_2) = \langle \Psi | \hat{\psi}^\dagger(z_1) \hat{\psi}^\dagger(z_2) \hat{\psi}(z_1) \hat{\psi}(z_2) | \Psi \rangle$

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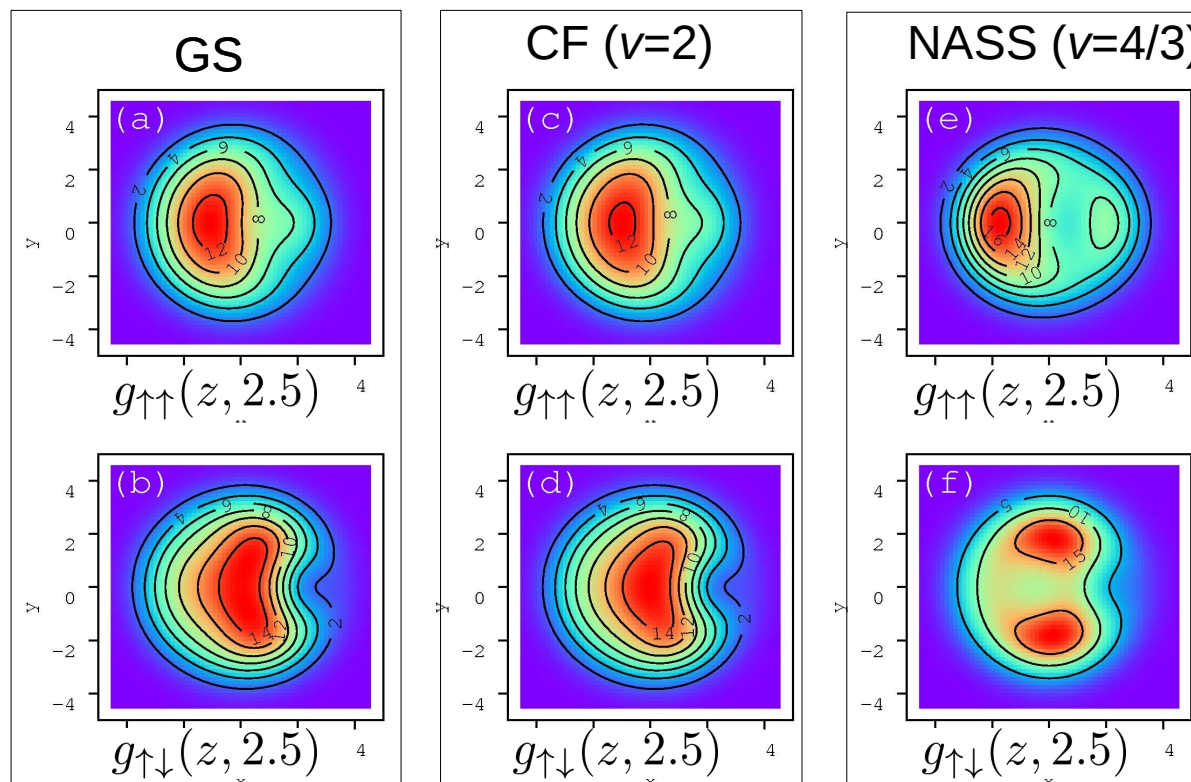
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What happens at $\nu=2$?

PRL 110, 046801 (2013)

PHYSICAL REVIEW LETTERS

week ending
25 JANUARY 2013

Integer Quantum Hall Effect for Bosons

T. Senthil¹ and Michael Levin²

¹*Department of Physics, Massachusetts Institute of Technology, Cambridge, Massachusetts 02139, USA*

²*Department of Physics, Condensed Matter Theory Center, University of Maryland, College Park, Maryland 20742, USA*

Effective field theory:
Possibility of an
interacting integer
quantum Hall effect
for two-component
bosons at $\nu=2$.

PRL 111, 090401 (2013)

PHYSICAL REVIEW LETTERS

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30 AUGUST 2013

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PHYSICAL REVIEW B 87, 245123 (2013)

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RAPID COMMUNICATIONS

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PHYSICAL REVIEW B 88, 161106(R) (2013)

Microscopic model for the boson integer quantum Hall effect

N. Regnault^{1,2} and T. Senthil³

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³*Department of Physics, Massachusetts Institute of Technology, Cambridge, Massachusetts 02139, USA*

PHYSICAL REVIEW B 89, 045114 (2014)

Quantum Hall phases of two-component bosons

T. Graß,¹ D. Raventós,² M. Lewenstein,^{1,3} and B. Juliá-Díaz^{1,2}

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What happens at $\nu=2$?

Torus:

Grass, Julia-Diaz, Barberan, Lewenstein (PRA, 2012)
Regnault & Senthil (PRB, 2013)
Furukawa & Ueda (PRL, 2013)

→ no NASS phase
→ unique, gapped GS

Sphere:

Furukawa & Ueda (PRL, 2013)
Wu & Jain (PRB, 2013)

Entanglement spectra:
→ edge physics of iIQHE

Overlap: with CF state
0.888 ($N=14$)

Disk:

Wu & Jain (PRB, 2013)
Grass, Raventos, Julia-Diaz, Lewenstein (PRB, 2014)

Edge spectrum agrees with IQH theory.

Overlap with CF state:
0.970 ($N=8, L=16$)

Edge spectrum at $\nu=2$

Effective edge Hamiltonian of singlet state

[J.E. Moore, F.D.M. Haldane, PRB **55** 7818 (1997)]

$$H_{\text{edge}} \propto v_s (S_z^2 + \sum_l l b_l^\dagger b_l) + v_c \sum_l l c_l^\dagger c_l$$

TABLE I. Number of modes of H_{edge} with $v_s < 0$ and $v_c > 0$.

ΔL_z	-4	-3	-2	-1	+1	+2	+3	+4
Number of singlets	2	1	1	0	1	2	3	5
Number of triplets	2	2	1	1	0	0	0	0
Number of quintets	1	0	0	0	0	0	0	0

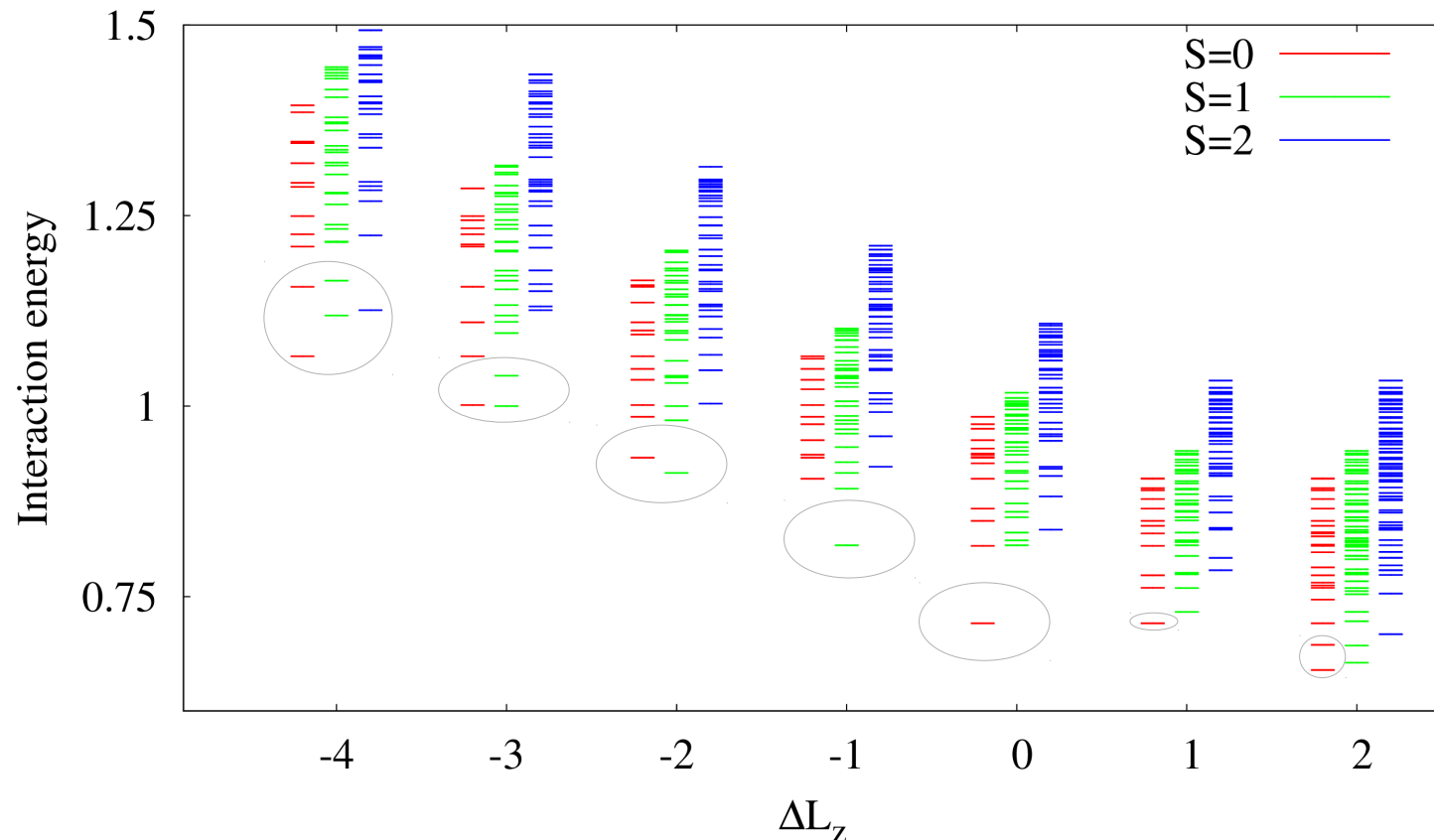
Numerical results on a disk

[T. Grass *et al.*, PRB **89** 045114 (2014)]

$$N_\uparrow = 4$$

$$N_\downarrow = 4$$

$$L_z = 16 + \Delta L_z$$



Edge spectrum at $\nu=2$

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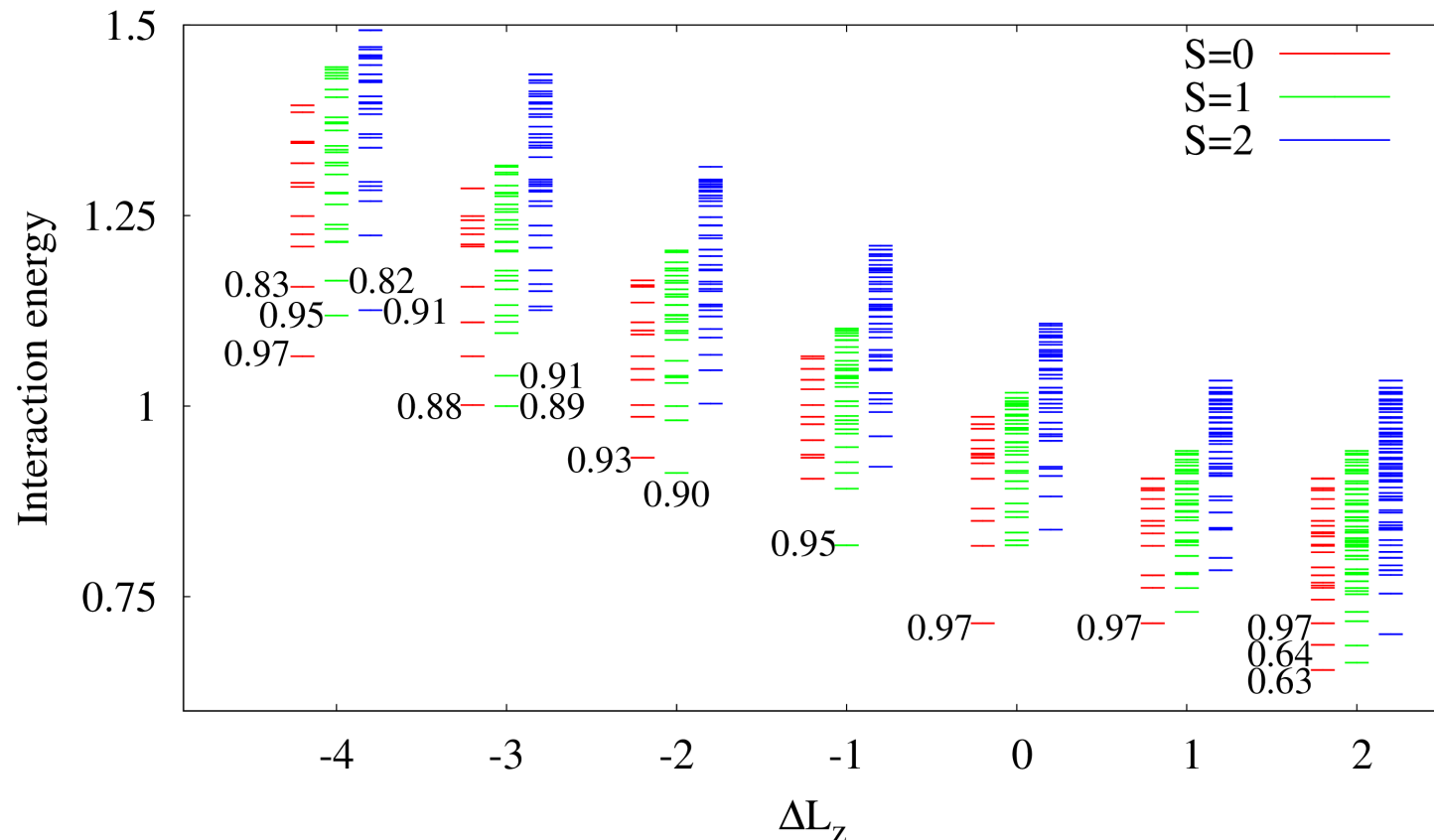
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$$L_z = 16 + \Delta L_z$$

Explicit construction of edge states provides the same counting, and good overlaps:

- **Backward states:**
Excite CFs in flux-reversed LIs
- **Forward states:**
Symmetric polynomial



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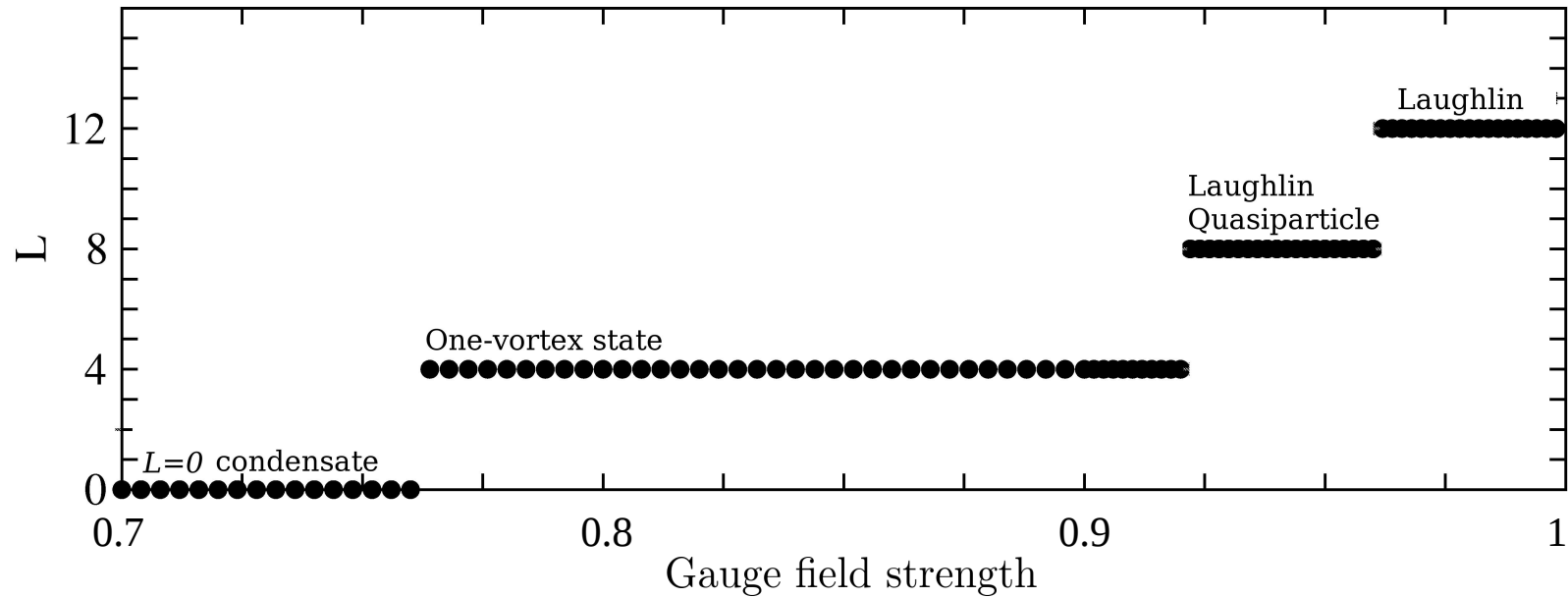
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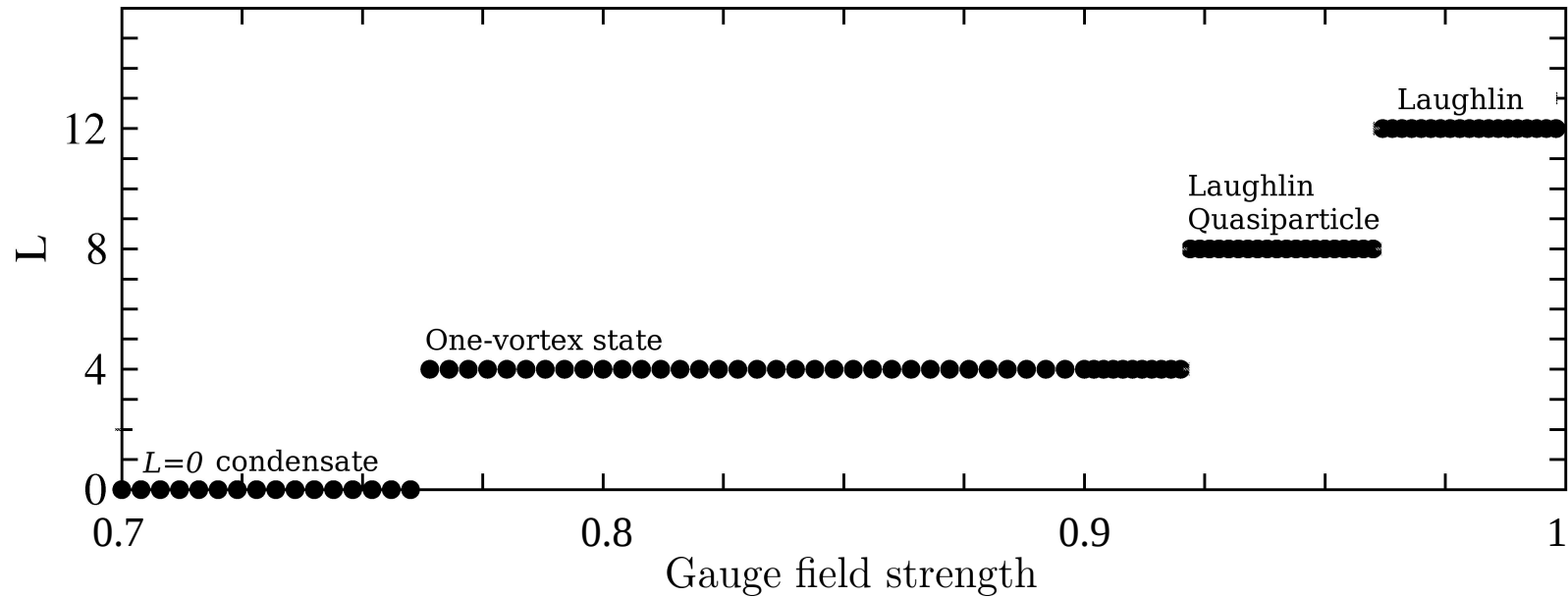
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One-component Bose gas



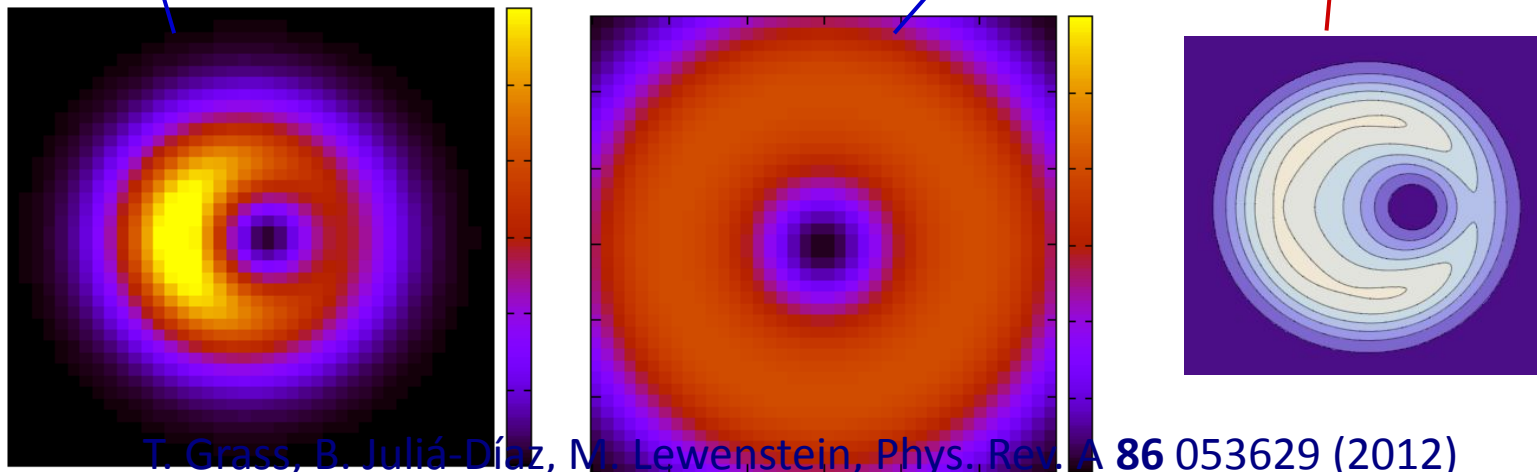
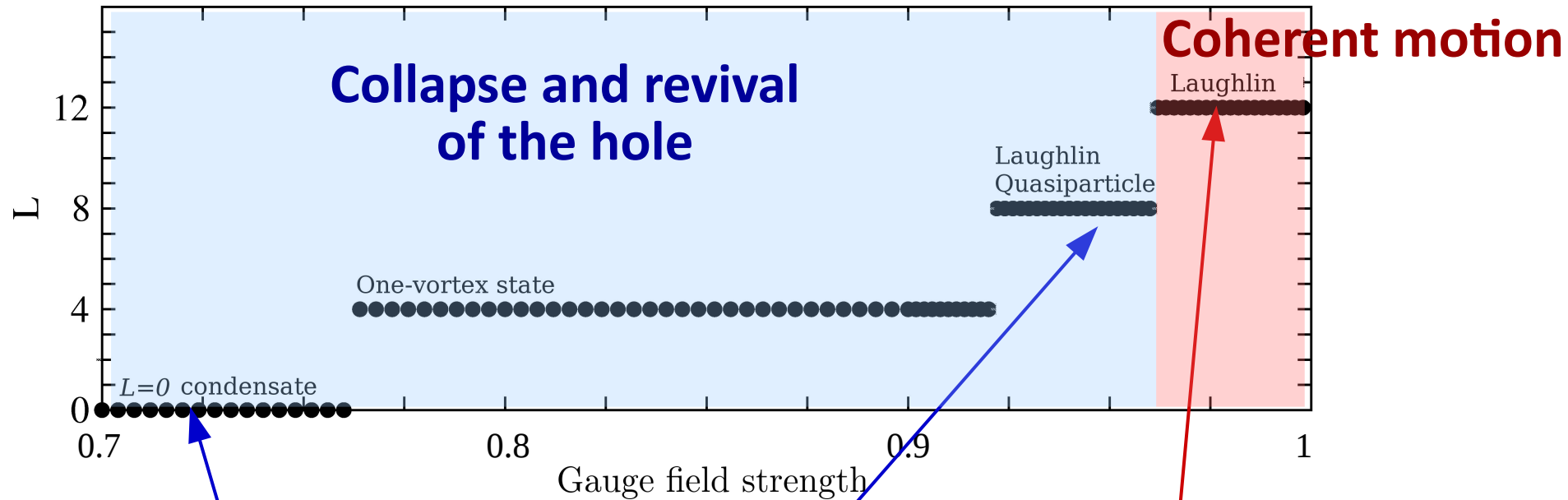
One-component Bose gas

Pierce holes $\Psi \rightarrow \prod_i (z_i - \xi) \Psi$ for laser potential $V_L \propto \delta(\xi)$.



One-component Bose gas

Pierce holes $\Psi \rightarrow \prod_i (z_i - \xi) \Psi$ for laser potential $V_L \propto \delta(\xi)$.

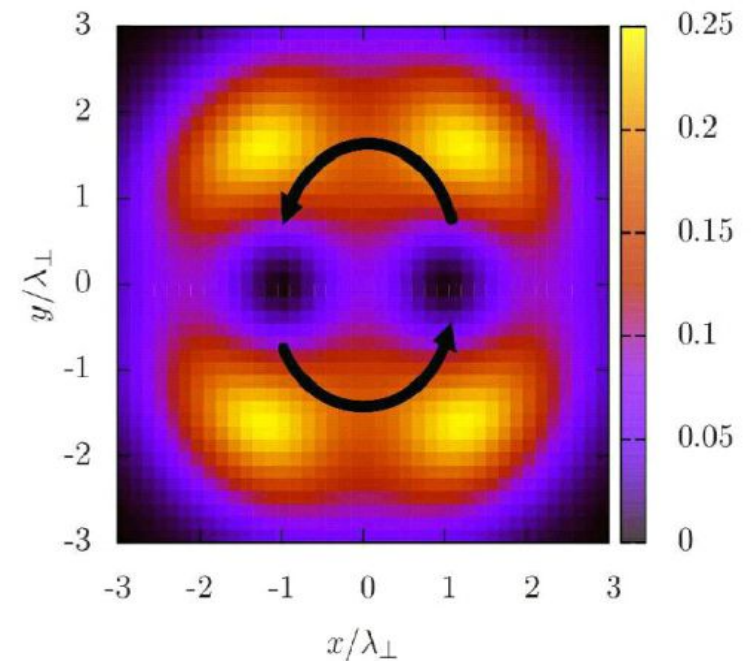


T. Grass, B. Juliá-Díaz, M. Lewenstein, Phys. Rev. A **86** 053629 (2012)

Quasiholes in FQH states

Holes in fractional quantum Hall systems:

- Candidates for anyonic quasi-particle excitations
- Fractional quantum statistics
- Atomic samples:
 - *Creation by a laser beam*
 - *Control of the position*
 - *Ramsey interferometry*



B. Paredes, P. Fedichev, J. I. Cirac, and P. Zoller. PRL **87**, 010402 (2001)

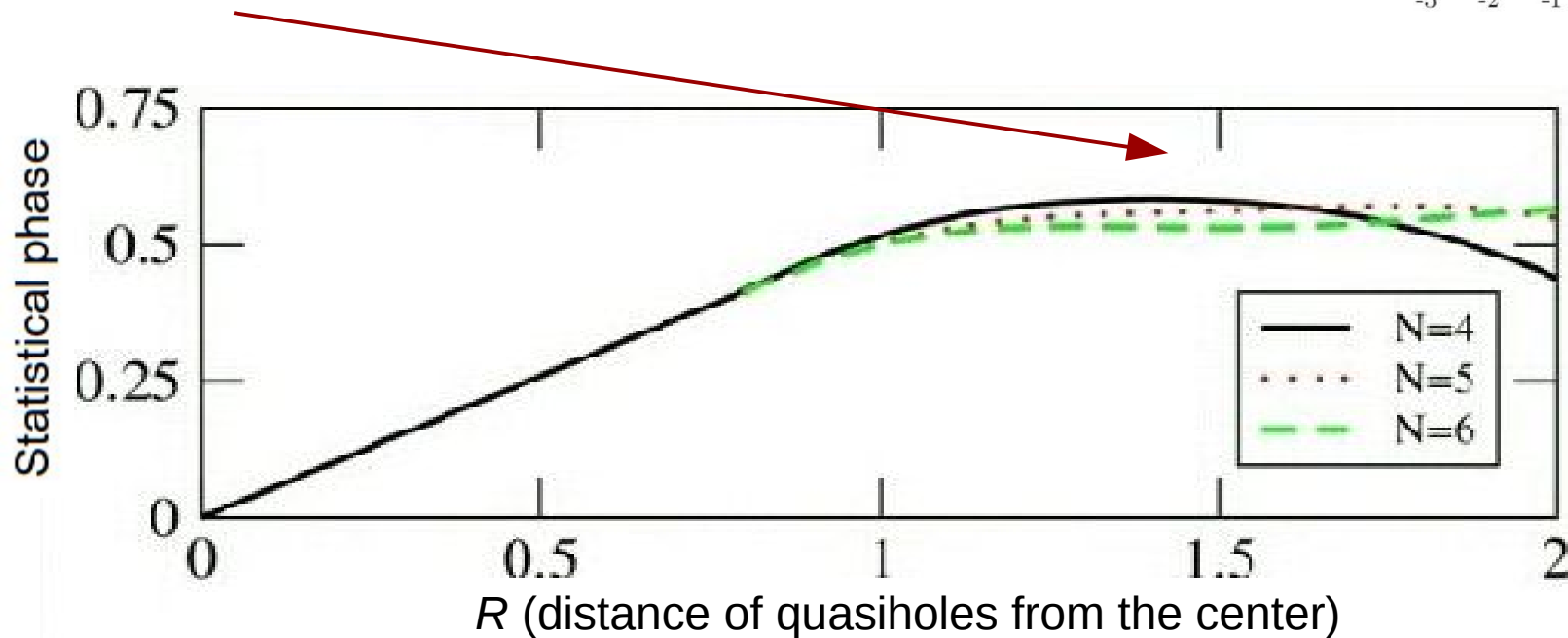
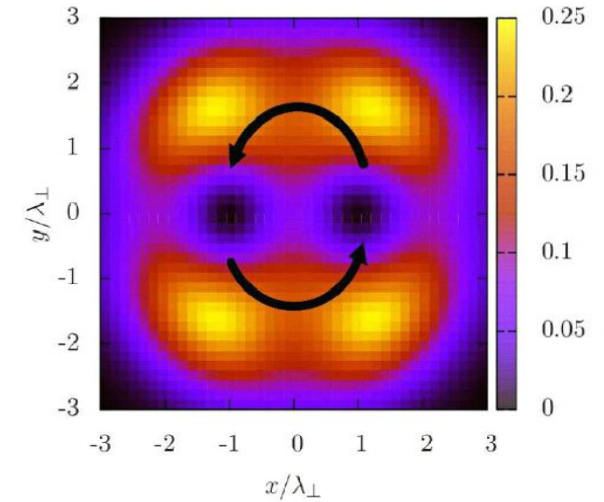
B. Juliá-Díaz, T. Grass, N. Barberán, M. Lewenstein, NJP **14**, 055003 (2012)

T. Grass, B. Juliá-Díaz, M. Lewenstein, PRA **89**, 013623 (2014)

Anyon braiding

Interchange two quasiholes and observe the phase difference: $\Psi \rightarrow e^{i\varphi} \Psi$

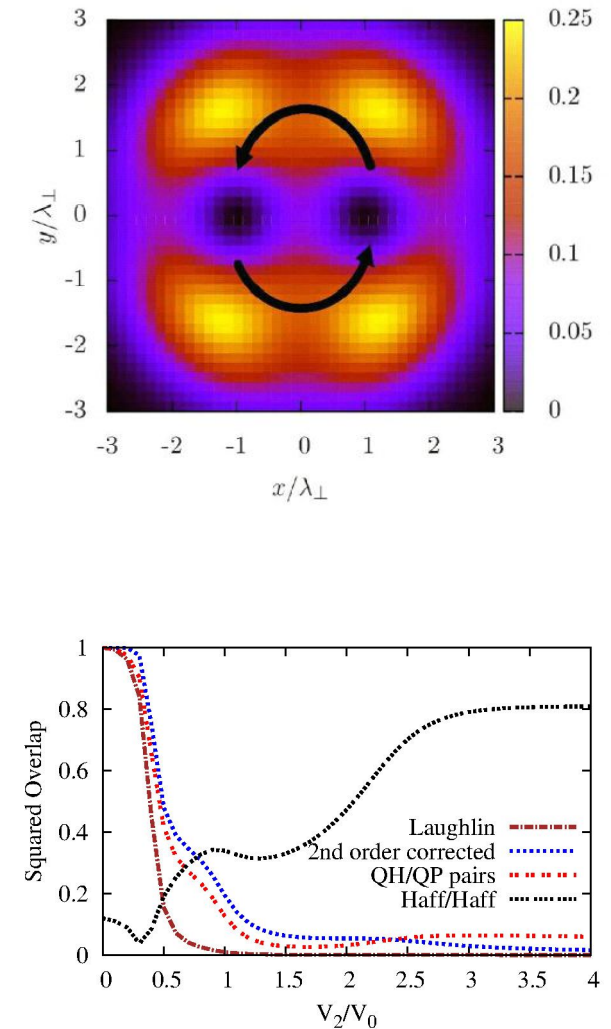
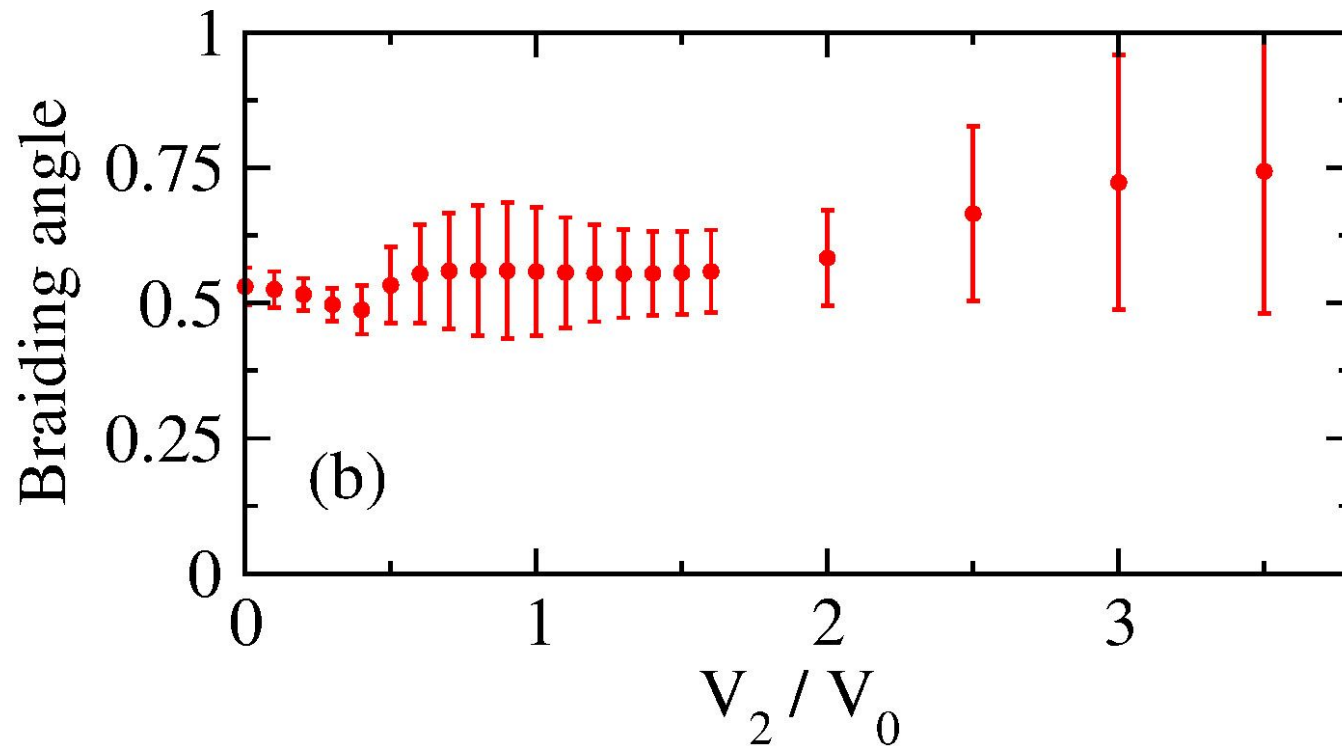
Fractional statistical phase becomes robust even in small Laughlin systems



B. Juliá-Díaz, T. Grass, N. Barberán, M. Lewenstein, NJP **14**, 055003 (2012)

Anyon braiding

To test the robustness of braiding phase, we switch on an additional, tunable d-wave interaction:



T. Grass, B. Juliá-Díaz, M. Lewenstein, PRA **89**, 013623 (2014)

Outline

1. Quantum Hall Physics – in general:

- Single-particle physics: Landau levels
- Many-body effects: Trial states

2. Fractional quantum Hall physics of spin-1/2 bosons:

- Abelian vs. Non-Abelian phases
- Numerical results: ambiguous
- Prospects of a quantum simulation

3. *Integer* quantum Hall physics of spin-1/2 bosons:

- Interactions are crucial
- Edge spectrum as a fingerprint

4. Anyon braiding in small systems

The people

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