

A precise determination of top quark electroweak couplings at the ILC

Ignacio García García
IFIC (UV-CSIC)



INDEX

- Introduction
- Experimental environment and data samples
- Event selection
- Observables and form factors
- Measurement of observables
 - *The Forward–Backward asymmetry: A_{FB}*
 - *Slope of the helicity distribution*
- Results
- Conclusions

Theory

- The top quark is the **heaviest elementary particle** and it is the most **strongly coupled** to the mechanism of **electroweak symmetry breaking**.
- In contrast to the situation at hadron colliders, the dominant pair production process $e^+e^- \rightarrow t\bar{t}$ involves only $t\bar{t}Z^0$ and $t\bar{t}\gamma$ **primary vertices**

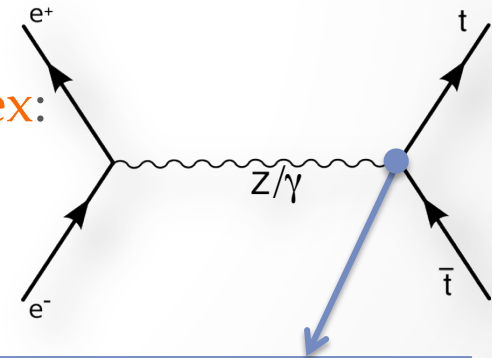
- A way to describe the **current** at the $t\bar{t}X$ **vertex**:

- $X = Z^0, \gamma$

- $V = \text{Vector coupling}$

- $A = \text{Axial coupling}$

arxiv.org/abs/hep-ph/0601112

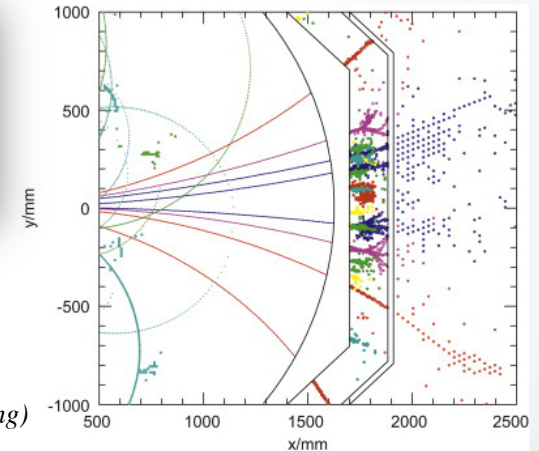
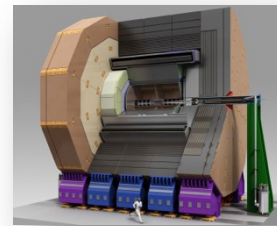


$$\Gamma_{\mu}^{ttX}(k^2, q, \bar{q}) = ie \left\{ \gamma_{\mu} \left(\tilde{F}_{1V}^X(k^2) + \gamma_5 \tilde{F}_{1A}^X(k^2) \right) + \frac{(q - \bar{q})_{\mu}}{2m_t} \left(\tilde{F}_{2V}^X(k^2) + \gamma_5 \tilde{F}_{2A}^X(k^2) \right) \right\}$$

International Linear Collider (ILC)

- The c.o.m. energy: $\sqrt{s} = 500 \text{ GeV}$ (default design)
- Luminosity: $\mathcal{L} = 500 \text{ fb}^{-1} = 5 \times 10^5 \text{ pb}^{-1}$ (estimated for 4 years of running)
- Beams are **polarised**:
 $P(e^-) \approx \pm 80\%$, $P(e^+) \approx \pm 30\%$.

ILD detector is optimised for Particle Flow Algorithm (PFLOW), i.e. measure particles in jet in the best suited sub-detectors



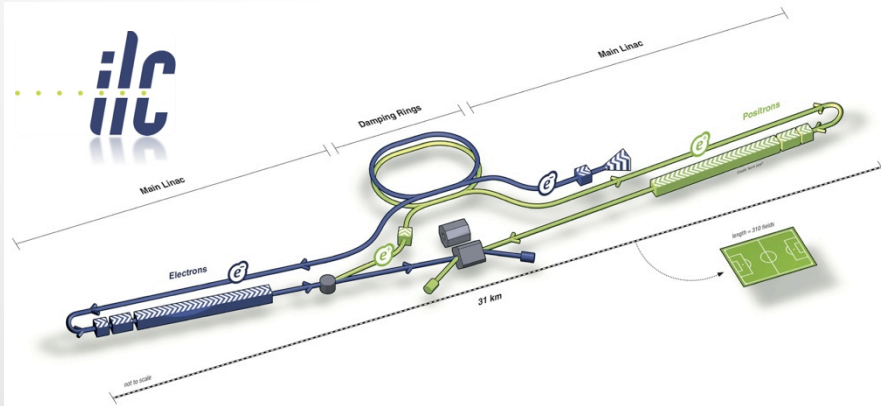
$$\sigma(d_0) = \left[5 \oplus \frac{1}{p_T} \frac{10}{\sqrt{\sin \theta}} \right] \mu\text{m}$$

$$\sigma(p_T) / p_T \approx 2 \cdot 10^{-5} p_T \oplus a$$

$$a = 1 \cdot 10^{-4} \text{ (multiple scattering)}$$

So the expected energy resolution is:

$$\sigma_E / E \sim 3\%$$



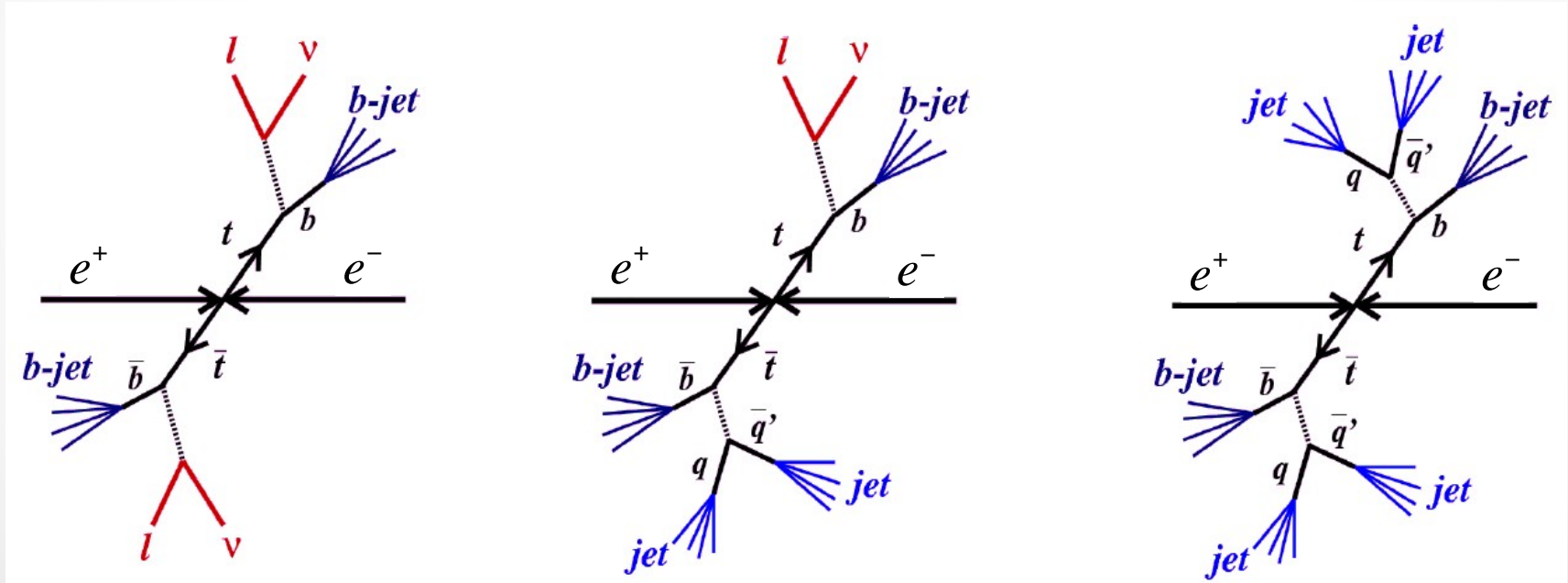
$t\bar{t}$ decay modes

$e^+e^- \rightarrow t\bar{t}$ gives three different final states:

Fully leptonic (10.3%)
2 jets + 2 leptons + 2 neutrinos

Semi-leptonic (43.5%)
4 jets + lepton + neutrino

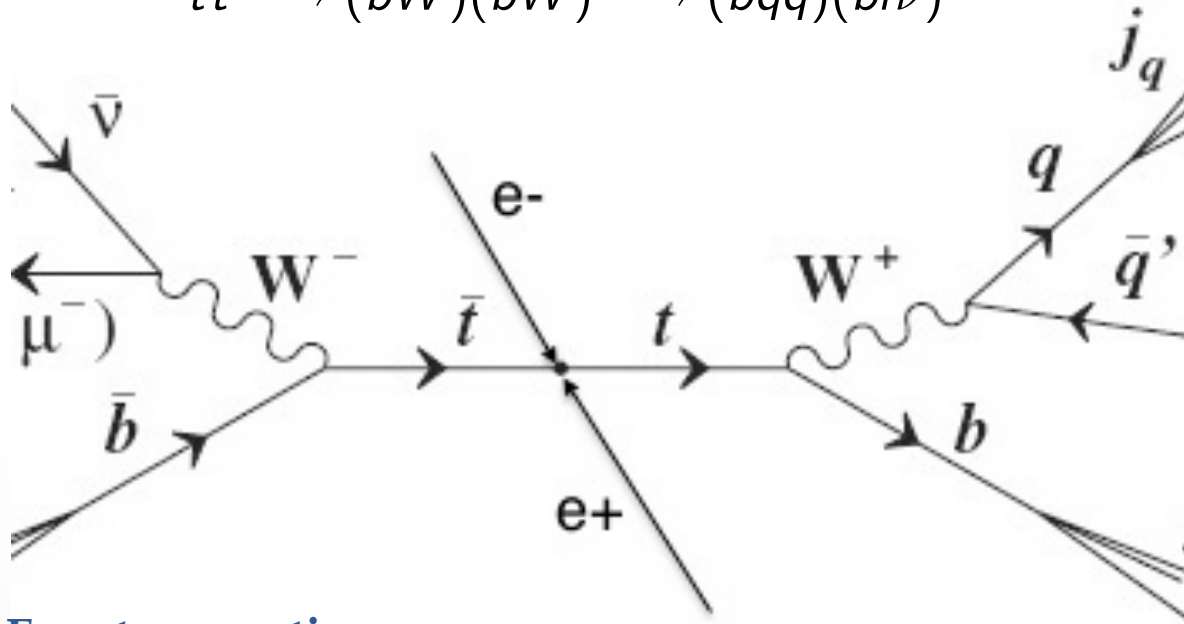
Fully hadronic (46.2%)
6 jets at final state



- ▣ This analysis concentrates mainly on events which have a **semi-leptonic final state**

<http://www-flc.desy.de/lcnotes/LC-REP-2013-007>

$$t\bar{t} \longrightarrow (bW)(bW) \longrightarrow (bqq)(bl\nu)$$



Event generation

- 1) **WHIZARD**: event generation (*samples for the DBD*)
- 2) **PYTHIA**: Parton shower and hadronisation

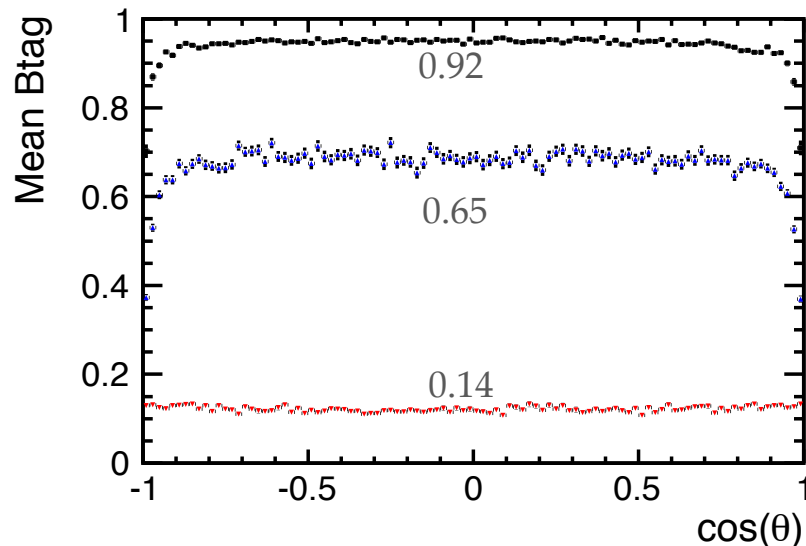
Event selection

- ▣ **Lepton identification** criteria:

- ▣ *Lepton is isolated from a jet* $x_T = p_{T,lepton}/M_{jet} > 0.25$ and $z = E_{lepton}/E_{jet} > 0.6$

Taking into account the τ leptons. → **Eff ~ 70%**

- ▣ ***b*-likeness or *b*-tag** is determined analysing secondary vertices → jet mass, decay length and particle multiplicity. A *b-tag* value is assigned to each jet.



$0 < b\text{-tag} < 1$

Tagged as *b*-jets

The 2 jets with higher *b*-tag value

Non *b*-jet

Event selection

- The signal is reconstructed by **choosing** the **combination of b quark jet and W boson** that minimises the following equation:

$$d^2 = \left(\frac{m_{cand.} - m_t}{\sigma_{m_t}} \right)^2 + \left(\frac{E_{cand.} - E_{beam}}{\sigma_{E_{cand.}}} \right)^2 + \left(\frac{p_b^* - 68}{\sigma_{p_b^*}} \right)^2 + \left(\frac{\cos\theta_{bW} - 0.23}{\sigma_{\cos\theta_{bW}}} \right)^2$$

- Some **cuts**:
 - **Hadronic mass** of the final state: $180 < m_{had.} < 420$ GeV
 - Reconstructed **W mass**: $50 < m_W < 250$ GeV
 - Reconstructed **top mass**: $120 < m_t < 270$ GeV
 - **Isolated lepton**: *the best candidate*
 - **b -tag values**: $b\text{-tag}_1 > 0.8$ & $b\text{-tag}_2 > 0.3$
- The **entire selection** retains:
 - **51.9%** for the configuration $P, P' = -1, +1$ (Left-handed electrons)
 - **55.0%** for $P, P' = +1, -1$ (Right-handed electrons)

Observables

- ▣ Total cross section (σ)
- ▣ The Forward-Backward Asymmetry (A_{FB})
- ▣ The slope of the distribution of the helicity angle (λ_{hel})

But actually there are **6 independent observables** = **3 observables** x **2 polarisations**

So once 6 observables are measured, we can obtain the following **CP conserving 5 couplings** of the top quark

$$\left. \begin{array}{l} \sigma(+), A_{FB}(+), \lambda_{hel}(+) \quad (+ = e_R^-) \\ \sigma(-), A_{FB}(-), \lambda_{hel}(-) \quad (- = e_L^-) \end{array} \right\} \Rightarrow \left\{ \begin{array}{l} F_{1V}^\gamma \quad * \quad F_{2V}^\gamma \\ F_{1V}^Z \quad F_{1A}^Z \quad F_{2V}^Z \end{array} \right\}$$

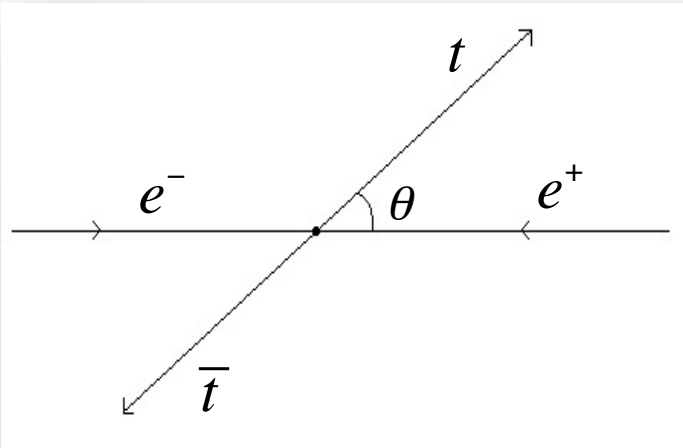
* $F_{1A}^\gamma = 0$ always because of the gauge invariance

Forward-Backward asymmetry: A_{FB}

▣ The Forward-Backward Asymmetry

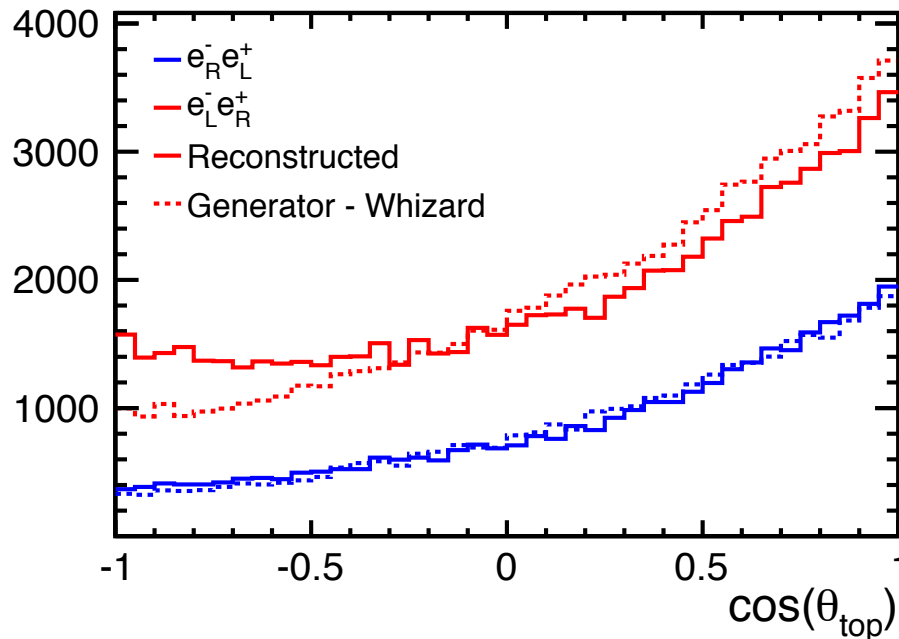
$$A_{FB} = \frac{N_{top}(\cos\theta > 0) - N_{top}(\cos\theta < 0)}{N_{top}(\cos\theta > 0) + N_{top}(\cos\theta < 0)}$$

$$-1 < A_{FB} < 1$$



- ▣ The **sign** of the **top** is the one of the **lepton**
- ▣ For \bar{t} we change θ to $\theta + \pi$

Results for A_{FB}



← We can see a clear **migration effect** for left-handed electrons

- ▣ This migration comes from the **wrong combination** of the **W** and the **b-jet** to **reconstruct the top** quark
- ▣ It occurs in about **30%** of the times.
- ▣ This gives a **wrong direction** of the reconstructed **top** and produces the **migration effect**.

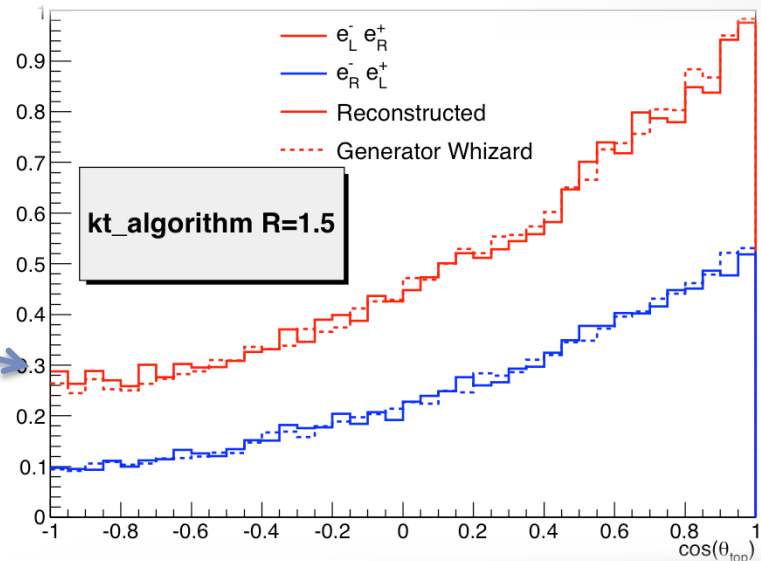
How to cure migration? χ^2 strategy

$$\chi^2 = \left(\frac{\gamma_t - 1.435}{\sigma_{\gamma_t}} \right)^2 + \left(\frac{E_b^* - 68}{\sigma_{E_b^*}} \right)^2 + \left(\frac{\cos\theta_{bW} - 0.26}{\sigma_{\cos\theta_{bW}}} \right)^2$$

- If we cut on χ^2 we reduce the number of wrong combinations of W and b-jet
- $\chi^2 < 15 \rightarrow$ Reconstruction efficiency : **29.6%**

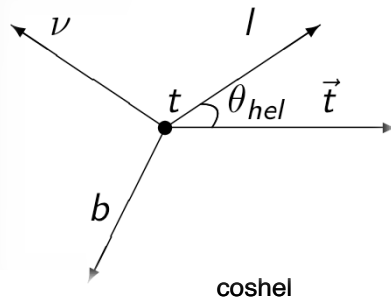
$\mathcal{P}, \mathcal{P}'$	$(A_{FB}^t)_{gen.}$	A_{FB}^t	$(\delta_{A_{FB}}/A_{FB})_{stat.} [\%]$
-1, +1	0.360	0.344	1.7 (for $\mathcal{P}, \mathcal{P}' = -0.8, +0.3$)
+1, -1	0.433	0.428	1.3 (for $\mathcal{P}, \mathcal{P}' = +0.8, -0.3$)

The χ^2 cut removes the migration effect
for left-handed electrons



Helicity angle (θ_{hel})

- In the rest frame of the top, θ_{hel} is the angle between the initial direction of the top and the lepton

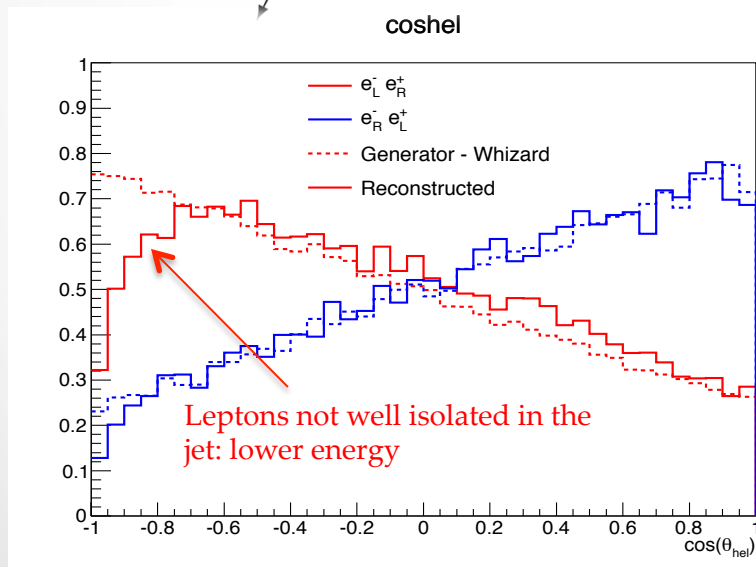


- The slope (λ_t) of the distribution gives the fraction of t_L and t_R in the sample.

$$\frac{1}{\Gamma} \frac{d\Gamma}{d\cos\theta_{hel}} = \frac{1 + \lambda_t \cos\theta_{hel}}{2} = \frac{1}{2} + (2F_R - 1) \frac{\cos\theta_{hel}}{2}$$

$$\lambda_t = 1 \text{ for } t_R \quad \lambda_t = -1 \text{ for } t_L$$

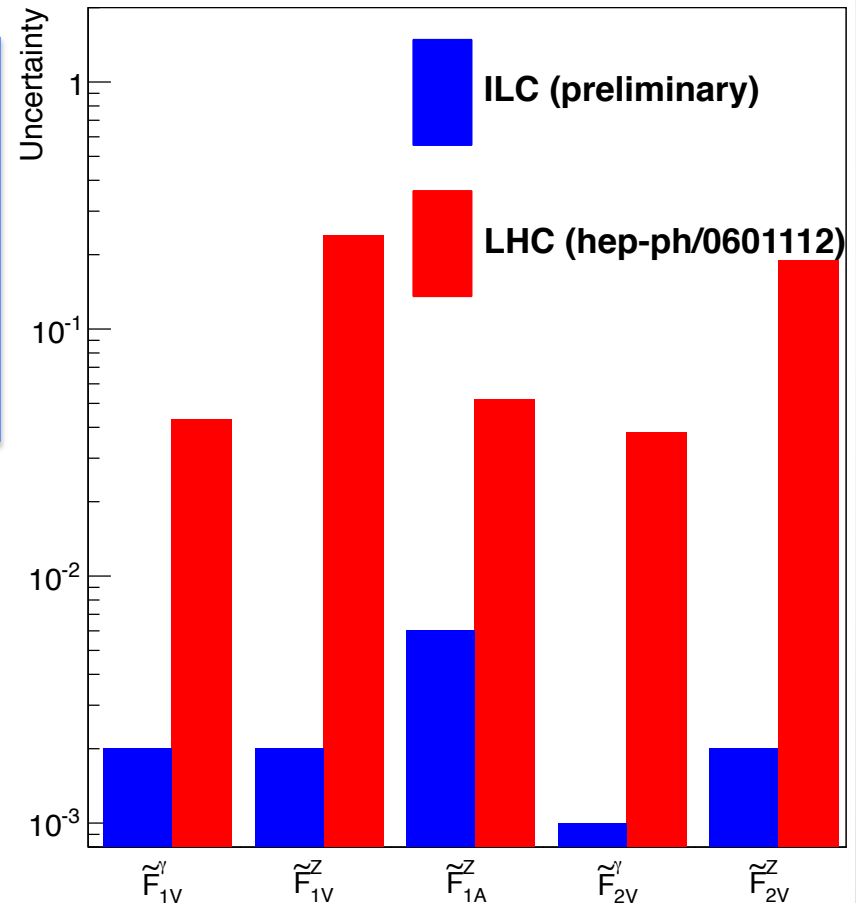
$\mathcal{P}, \mathcal{P}'$	$(\lambda_t)_{gen.}$	$(\lambda_t)_{rec.}$	$(\delta\lambda_t)_{stat.}$ for $\mathcal{P}, \mathcal{P}' = \mp 0.8, \pm 0.3$	$(\delta\lambda_t)_{syst.}$
-1, +1	-0.519	-0.489	0.016	0.011
+1, -1	0.544	0.547	0.016	0.010



Summary of the results

Coupling	SM value	LHC [1] $\mathcal{L} = 300 \text{ fb}^{-1}$	e^+e^- [6] $\mathcal{L} = 300 \text{ fb}^{-1}$ $\mathcal{P}, \mathcal{P}' = -0.8, 0$	e^+e^- [ILC DBD] $\mathcal{L} = 500 \text{ fb}^{-1}$ $\mathcal{P}, \mathcal{P}' = \pm 0.8, \mp 0.3$
$\Delta \tilde{F}_{1V}^\gamma$	0.66	+0.043 -0.041	- -	+0.002 -0.002
$\Delta \tilde{F}_{1V}^Z$	0.23	+0.240 -0.620	+0.004 -0.004	+0.002 -0.002
$\Delta \tilde{F}_{1A}^Z$	-0.59	+0.052 -0.060	+0.009 -0.013	+0.006 -0.006
$\Delta \tilde{F}_{2V}^\gamma$	0.015	+0.038 -0.035	+0.004 -0.004	+0.001 -0.001
$\Delta \tilde{F}_{2V}^Z$	0.018	+0.270 -0.190	+0.004 -0.004	+0.002 -0.002

- $F_{1}^{(\gamma/Z)}$ form factors can be extracted simultaneously considering $\Delta F_{2V}^{(\gamma/Z)} = 0$
- $F_{2V}^{(\gamma/Z)}$ are extracted **fixing** all $F_{1}^{(\gamma/Z)}$ to their SM values



Conclusions

- ▣ **Polarisation** is a powerful tool for analysis because it allows to **double** the number of **observables**
- ▣ **In LC with polarised beams** we can measure $\bar{t}tZ$ and $\bar{t}t\gamma$ vertices with **accuracies** one or two orders of magnitude **better than LHC**
- ▣ **Current aim:** We are looking for **new observables** sensitive to the **CPV form factors $F_{2A}^{\gamma/Z}$** . For instance spin correlations between the lepton and the top.

Thanks for
your attention
(you are free for leaving)

Particle Flow

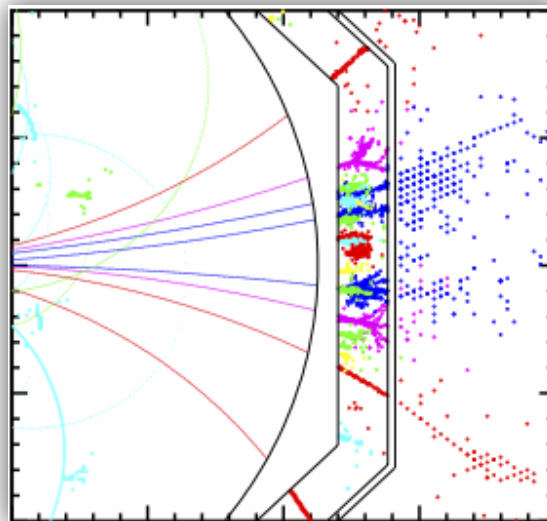
Particle Flow (a powerful tool to measure the energy of the jets)

- Measurement of the **charged particle momentum** in the tracker → charged component of the jet
- Measurement of the **momentum of the neutral component of the jet** = total energy measured in the calorimetry – energy of the charged particles in the calorimeter.
- Total energy of the jet = charged component + neutral component**

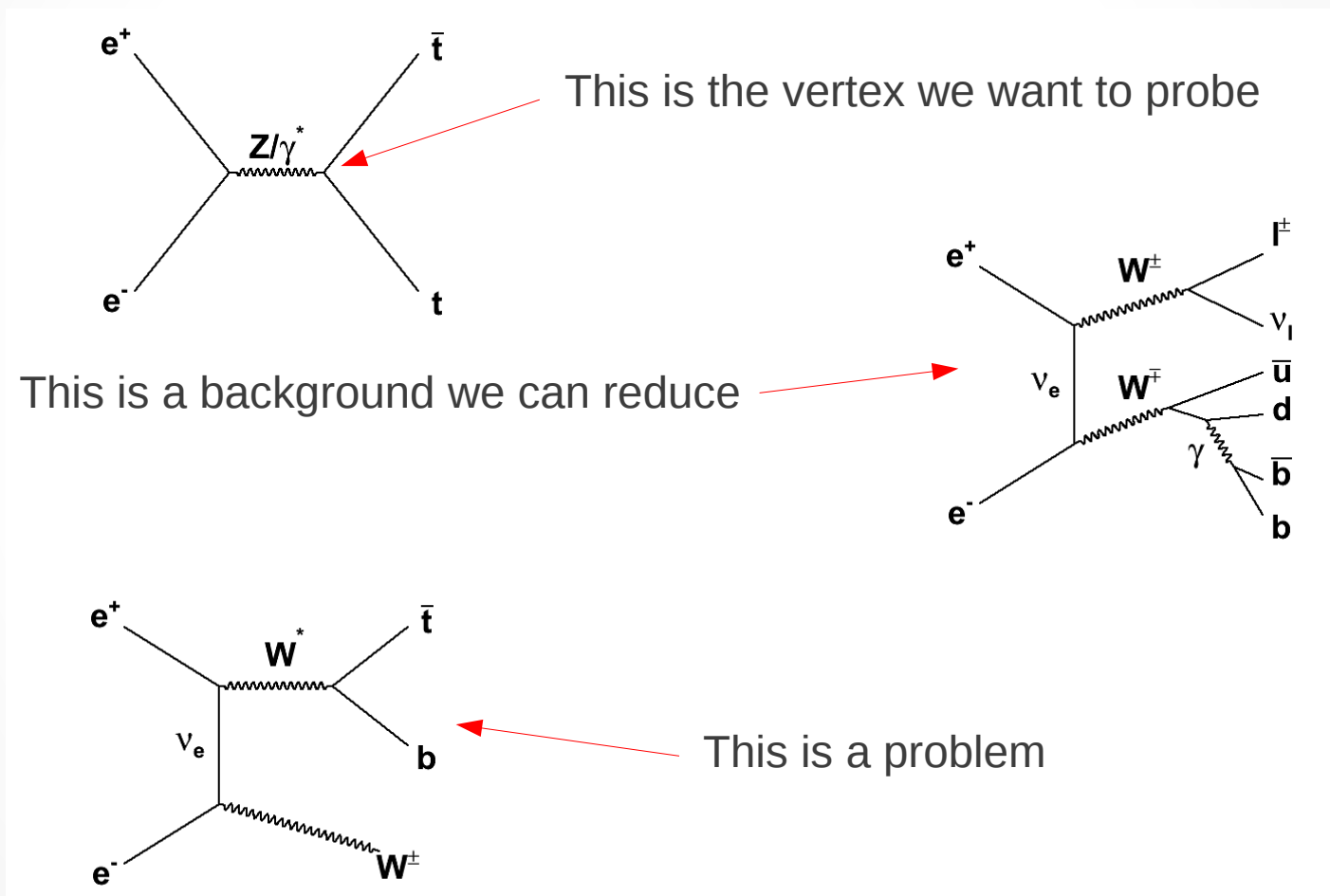
$$\sigma_E / E \approx 3\% \quad (E \text{ en GeV})$$

Great granularity of the calorimeters

Calormeter (Silicon-Tungsten)

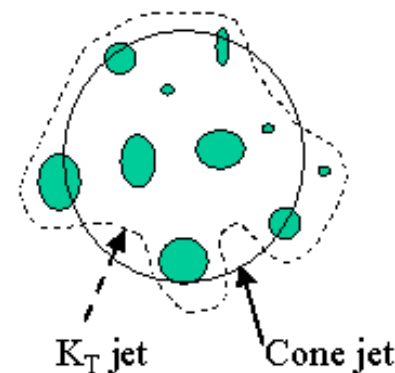


Single top



kt algorithm FastJet

<http://arxiv.org/pdf/1111.6097v1.pdf>



All clusters with $r < D$ are merged
 Clusters with $r > D$ can be merged if $\Delta E_T > 0$

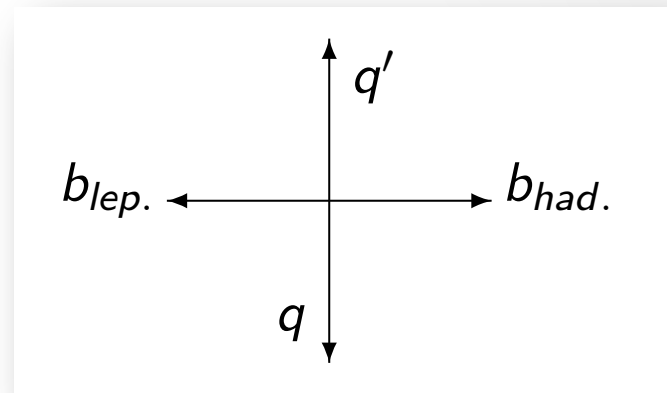
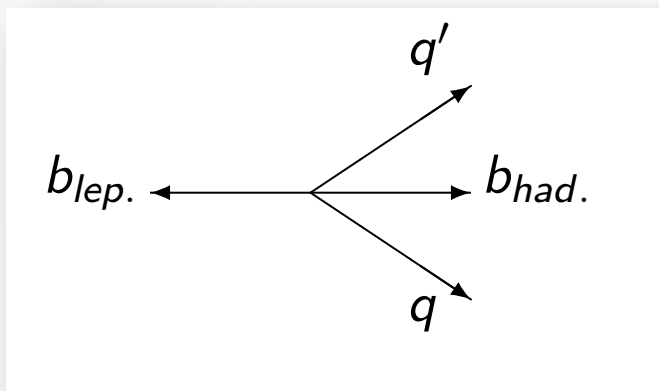
1. For each pair of particles i, j work out the k_t distance

$$d_{ij} = \min(k_{ti}^2, k_{tj}^2) \Delta R_{ij}^2 / R^2$$

with $\Delta R_{ij}^2 = (y_i - y_j)^2 + (\phi_i - \phi_j)^2$, where k_{ti} , y_i and ϕ_i are the transverse momentum, rapidity and azimuth of particle i and R is a jet-radius parameter usually taken of order 1; for each parton i also work out the beam distance $d_{iB} = k_{ti}^2$.

2. Find the minimum d_{\min} of all the d_{ij}, d_{iB} . If d_{\min} is a d_{ij} merge particles i and j into a single particle, summing their four-momenta (this is E -scheme recombination); if it is a d_{iB} then declare particle i to be a final jet and remove it from the list.
3. Repeat from step 1 until no particles are left.

Where does this migration comes from?



- ▣ **Right-handed** electron beam:
 - ▣ The W is emitted into the flight direction of the top together with a soft b
- ▣ In the case is the W is easily combine to good b to reconstruct the top

- ▣ **Left-handed** electron beam:
 - ▣ The W is emitted almost at rest together with a hard b
- ▣ In the case it is **harder to combine the W and the good b** to reconstruct the top

Observables SM values

- ▣ Total cross section (σ)
- ▣ The Forward-Backward Asymmetry (A_{FB})
- ▣ The slope of the distribution of the helicity angle (λ_{hel})

But actually there are **6 independent observables**

$\sigma(+)$	$A_{FB}(+)$	$\lambda_{hel}(+)$	(+ = e_R^-)
$\sigma(-)$	$A_{FB}(-)$	$\lambda_{hel}(-)$	(- = e_L^-)

- ▣ The expected values in the Standard Model are:

Observables	$e_L^- e_R^+$	$e_R^- e_L^+$
$\sigma(fb)$	1564	724
A_{FB}	0.38	0.47
F_R	0.25	0.76

where F_R is the fraction of right-handed tops

← $\lambda_{hel} = 2F_R - 1$