

# The zeta-regularized spectral determinant and vacuum energy of graph Schrödinger operators

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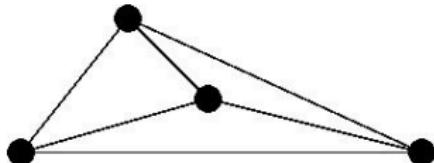
# Outline

1 Graph model

2 Dirichlet wire

3 General graph

# Quantum graph model



$V$  vertices  
 $B$  bonds

*Metric graph:* bond  $b$  corresponds to interval  $[0, L_b]$ .  
*Eigenproblem* on  $[0, L_b]$ ,

$$\left( i \frac{d}{dx_b} + A_b \right)^2 \psi_b(x_b) + V_b(x_b) \psi_b(x_b) = E \psi_b(x_b) . \quad (1)$$

Matching conditions at vertices so Schrödinger op. self-adjoint on graph. Hilbert space  $\mathcal{H} := \bigoplus_{b=1}^B L^2([0, L_b])$ .

**Graph spectrum:**  $0 < E_0 \leqslant E_1 \leqslant \dots$

## Spectral determinant

$$S(\gamma) = \prod_{j=0}^{\infty} (\gamma + E_j)$$

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- (99) Pascaud & Montambaux.
- (00) Akkermans, Comtet, Desbois, Montambaux & Texier.
- (06) Friedlander: Dirichlet to Neumann map,  $\delta$ -coupling conditions.
- (10) Texier:  $\delta$ -coupling conditions with potential, conjecture for general matching conditions [3].
- (11) JH & Kirsten: Laplacian with general vertex matching conditions [2].

## Vacuum energy

$$E_c = \frac{1}{2} \sum_{j=0}^{\infty} ' \sqrt{E_j}$$

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### Models of vacuum energy:

- (01) Leboeuf, Monastra, & Bohigas: Riemannium.
- (07) Fulling, Kaplan & Wilson: star graphs, repulsive Casimir effect for  $B > 3$ .
- (09) Berkolaiko, JH & Wilson: general graphs, effect of chaotic classical dynamics.

## Spectral zeta function

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## Spectral properties

$$S(\gamma) := \exp\left(-\zeta'(0, \gamma)\right) = \prod_{j=0}^{\infty} (\gamma + E_j)$$

$$E_c := \frac{1}{2}\zeta(-1/2, 0) = \frac{1}{2} \sum_{j=1}^{\infty} \sqrt{E_j}$$



## Dirichlet wire

Single interval  $[0, L]$  with  $\psi(0) = \psi(L) = 0$ .

$$\left( -\frac{d^2}{dx^2} + V(x) \right) \psi(x) = k^2 \psi(x) \quad (2)$$

$f_k(x)$  soln. with  $f_k(0) = 1$ ,  $f_k(L) = 0$ .  $\bar{f}_k(\bar{x})$  linearly independent  
soln. where  $\bar{x} = L - x$  &  $\bar{f}_k(0) = 1$ ,  $\bar{f}_k(L) = 0$ .

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Alternatively  $u_k(x)$ ,  $\bar{u}_k(\bar{x})$  solns. with  $u_k(0) = 0$  and  $u'_k(0) = 1$ ,  $\bar{u}_k(x) = f_k(L - x)/f'_k(L)$ .

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### Secular eqn.

$k^2$  is eigenval. of Schrödinger op. iff

$$\bar{u}_k(L) = \frac{-1}{f'_k(L)} = 0 .$$

## General method

Secular eqn.  $F(z) = 0$  related to  $\zeta(s, \gamma)$  via argument principle,

$$\zeta(s, \gamma) = \int_C \frac{dz}{2i\pi} (z^2 + \gamma)^{-s} \frac{d}{dz} \log(F(z)) . \quad (3)$$

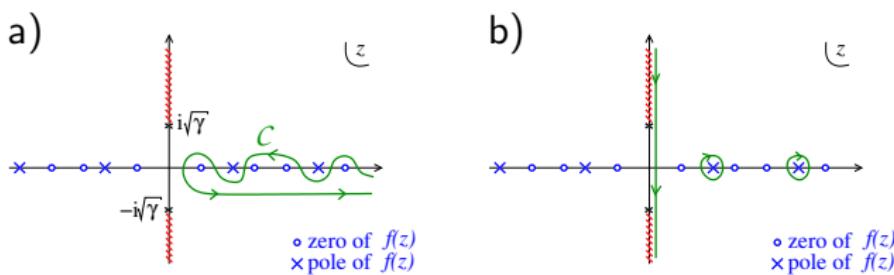


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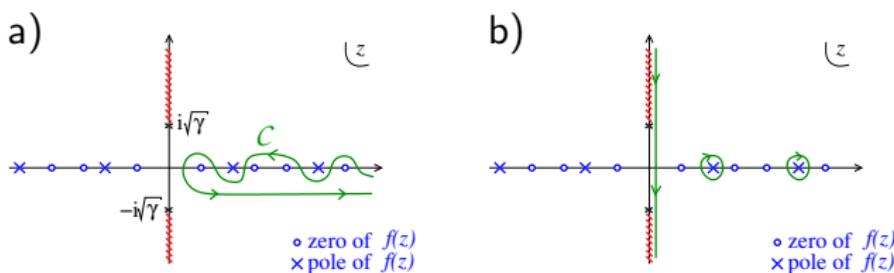


Figure: Contours (a) before, & (b) after, contour transformation.

**Dirichlet wire:**  $F(z) = u_z(L)$ , on imaginary axis  $z = it$ ,

$$\zeta_{\text{Dir}}(s, \gamma) = \frac{\sin \pi s}{\pi} \int_{\sqrt{\gamma}}^{\infty} dt (t^2 - \gamma)^{-s} \frac{d}{dt} \log(u_{it}(L)) .$$

Representation of  $\zeta_{\text{Dir}}(s, \gamma)$  valid  $1/2 < \text{Re } s < 1$ .

$$\log u_{it}(L) \underset{t \rightarrow \infty}{\sim} tL - \log 2t + O(t^0) \quad (4)$$

Subtracting and adding two terms in asymptotic expansion,

$$\begin{aligned} \zeta_{\text{Dir}}(s, \gamma) &= \frac{\sin \pi s}{\pi} \int_{\sqrt{\gamma}}^{\infty} dt (t^2 - \gamma)^{-s} \frac{d}{dt} [\log(u_{it}(L)) - tL + \log(2t)] \\ &\quad + L \frac{\Gamma(s - 1/2)}{2\sqrt{\pi}\Gamma(s)} \gamma^{\frac{1}{2}-s} - \frac{1}{2} \gamma^{-s} \quad -1/2 < \text{Re } s < 1. \end{aligned}$$

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## Spectral det of Dirichlet wire

$$S_{\text{Dir}}(\gamma) = \exp -\zeta'_{\text{Dir}}(0, \gamma) = 2u_{i\sqrt{\gamma}}(L) = \frac{-2}{f'_{i\sqrt{\gamma}}(L)}$$

## General graph

- $\psi, \psi'$  vectors of values of  $\psi$  & its covariant derivative at the  $2B$  ends of the  $B$  bonds.
- Matching conditions specified by matrices  $\mathbb{A}$  &  $\mathbb{B}$ ,

$$\mathbb{A}\psi + \mathbb{B}\psi' = \mathbf{0} . \quad (5)$$

- $\mathbb{A}, \mathbb{B}$  define a self-adjoint Schrödinger op. iff  $\text{rank}(\mathbb{A}, \mathbb{B}) = 2B$  and  $\mathbb{A}\mathbb{B}^\dagger = \mathbb{B}\mathbb{A}^\dagger$  (Kostrykin and Schrader).
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- *Local matching conditions:*  $\mathbb{A}, \mathbb{B}$  independent of  $\{L_b\}$  &  $E$ .

e.g. Neumann like matching at a vertex  $v$ :  $\psi$  continuous and  $\sum_{b \sim v} \psi'_b(0) = 0$ .

$$\mathbb{A} = \begin{pmatrix} 1 & -1 & 0 & \dots & 0 \\ 0 & 1 & -1 & \ddots & \vdots \\ \vdots & \ddots & \ddots & \ddots & 0 \\ 0 & \dots & 0 & 1 & -1 \\ 0 & \dots & 0 & 0 & 0 \end{pmatrix} , \quad \mathbb{B} = \begin{pmatrix} 0 & 0 & \dots & 0 \\ \vdots & \vdots & & \vdots \\ 0 & 0 & \dots & 0 \\ 1 & 1 & \dots & 1 \end{pmatrix} .$$

# Secular equation

$f_{b,k}(x_b), f_{\bar{b},k}(x_{\bar{b}})$  pair of lin. indep. solns of Schrödinger eqn on  $[0, L_b]$  s.t.  $f_{b,k}(0) = 1$  and  $f_{b,k}(L_b) = 0$ ,  $x_{\bar{b}} = L_b - x_b$ .

Wavefunction on bond  $b$ ,

$$\psi_{b,k}(x_b) = c_b f_{b,k}(x_b) e^{iA_b x_b} + c_{\bar{b}} f_{\bar{b},k}(x_{\bar{b}}) e^{iA_{\bar{b}} x_{\bar{b}}} . \quad (6)$$

$\psi = (c_1, \dots, c_B, c_{\bar{1}}, \dots, c_{\bar{B}})^T$  and  $\hat{\psi} = M(k)\psi$  where,

$$[M(k)]_{ab} = \delta_{a,b} f'_{b,k}(0) - \delta_{a,\bar{b}} f'_{\bar{b},k}(L_b) e^{iA_b L_b} . \quad (7)$$

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Substitute in matching conditions  $\mathbb{A}\psi + \mathbb{B}\psi' = \mathbf{0}$ .

## Secular equation

$$\det(\mathbb{A} + \mathbb{B}M(k)) = 0$$

$F(z) := \det(\mathbb{A} + \mathbb{B}M(z))$ , poles at eigenvalues of graph with Dirichlet bc's, as  $f'_{b,k}(L_b) = -1/u_{b,k}(L)$ ,

$$\zeta(s, \gamma) = \zeta_{\text{Dir}}(s, \gamma) + \zeta_{\text{Im}}(s, \gamma) .$$

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## Asymptotics:

$$F(it) \underset{t \rightarrow \infty}{\sim} \det \left( \mathbb{A} + \mathbb{B}(-tI + D(t)) \right) = \sum_{j=0}^{\infty} c_j t^{2B-j} \quad (8)$$

$D(t) = \sum_{j=1}^{\infty} t^{-j} \text{diag} \left\{ s_{1,j}(0), \dots, s_{B,j}(0), s_{\bar{1},j}(0), \dots, s_{\bar{B},j}(0) \right\}$ ,  
 $s_{b,1}(x_b) = -V_b(x_b)/2$ , subsequent fns. determined by recursion relation.

Let  $c_N$  be 1<sup>st</sup> non-zero coeff,  $N = 0$  and  $c_0 = \det \mathbb{B}$  if  $\det \mathbb{B} \neq 0$ .

$\zeta_{\text{Im}}(s, \gamma)$  obtained using previous contour integration method.

## Zeta function

$$\begin{aligned}\zeta(s, \gamma) = & \frac{\sin \pi s}{\pi} \int_{\sqrt{\gamma}}^{\infty} (t^2 - \gamma) - s \frac{d}{dt} \log(F(it)t^{N-2B}/c_N) dt \\ & + \frac{(2B-N)\sin \pi s}{2\pi s} \gamma^{-s} + \zeta_{\text{Dir}}(s, \gamma) \quad -1/2 < \text{Re } s < 1\end{aligned}$$

$$\begin{aligned}\zeta'_{\text{Im}}(0, \gamma) &= \log c_N - \log(F(i\sqrt{\gamma})\gamma^{\frac{N-2B}{2}}) - \frac{2B-N}{2} \log \gamma \\ &= \log c_N - \log \det(\mathbb{A} + \mathbb{B}M(i\sqrt{\gamma}))\end{aligned}$$

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## Spectral determinant

$$\begin{aligned}S(\gamma) &= S_{\text{Dir}}(\gamma) \exp(\zeta'_{\text{Im}}(0, \gamma)) \\ &= \left( \prod_{b=1}^B \frac{-2}{f'_{b, i\sqrt{\gamma}}(L_b)} \right) \frac{\det(\mathbb{A} + \mathbb{B}M(i\sqrt{\gamma}))}{c_N}\end{aligned}$$

# Vacuum energy

$$\log F(it) \underset{t \rightarrow \infty}{\sim} \log(c_N t^{2B-N}) + \frac{c_{N+J}}{c_N t^J} + O(t^{-(J+1)}) \quad (9)$$

$c_N$  and  $c_{N+J}$  first 2 non-vanishing coeffs. in expansion of  $\det(\mathbb{A} + \mathbb{B}(-tI + D(t)))$ .

$$\begin{aligned} \zeta_{\text{Im}}(s, \gamma) &= \frac{(2B - N) \sin \pi s}{2\pi s} \gamma^{-s} - \frac{J c_{N+J} \Gamma(s + J/2)}{2 c_N \Gamma(s) \Gamma(J/2 + 1)} \gamma^{-(2s+J)} \\ &+ \frac{\sin \pi s}{\pi} \int_{\sqrt{\gamma}}^{\infty} (t^2 - \gamma)^{-s} \frac{d}{dt} \left( \log \left( \frac{F(it)t^{N-2B}}{c_N} \right) - \frac{c_{N+J}}{c_N t^J} \right) dt \end{aligned}$$

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Generically  $J = 1$  and  $E_c = \zeta(-1/2, 0)/2$  divergent.

## Casimir force on bond $\beta$

$$\begin{aligned}\mathcal{F}_c^\beta &= \frac{\partial}{\partial L_\beta} E_c \\ &= \mathcal{F}_{c,\text{Dir}}^\beta + \int_0^\infty \frac{\partial}{\partial L_\beta} \log F(it) dt \\ F(it) &= \det(\mathbb{A} + \mathbb{B}M(it))\end{aligned}$$

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- Obtained repn of zeta fns of general graph Schrödinger ops.
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## Outlook

- Bootstrap results to thin networks.
- Implications for vac energy of chaotic systems.
- Extend to Dirac op.

## References

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